

# Precision event generation for top-quark production



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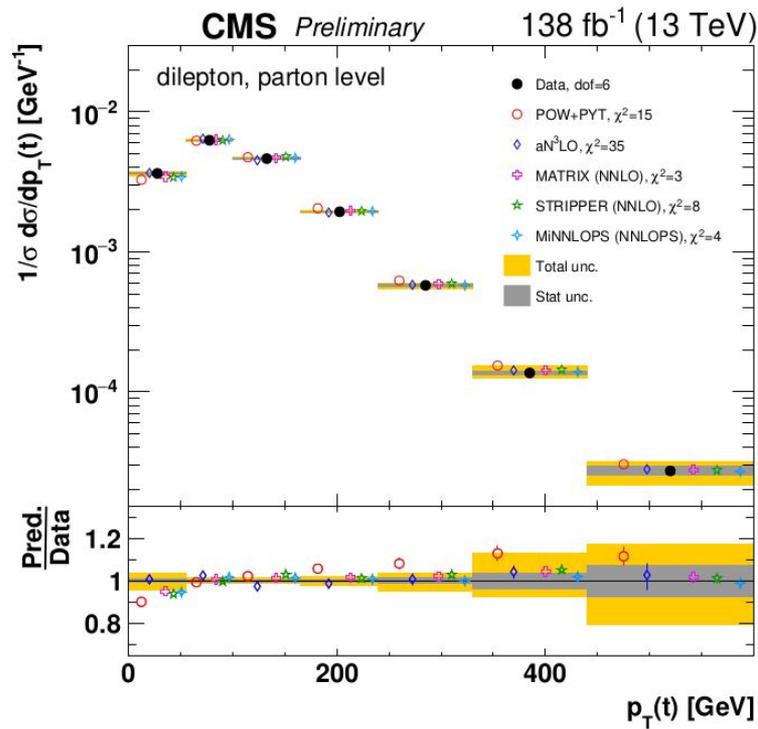
# Outline & introduction

- review of some recent results
- focus mostly on  $t\bar{t}$ , few comments on  $t\bar{t}+X$

- Fully-differential partonic results: NNLO QCD (prod. and decay), NLO EW corrections known (+ analytic resummation)
- MC event generators (from data to parton-level / background for BSM / measure parameters):
  - stable tops: NLO+PS established, **NNLO+PS available**, approximated decays
  - **full final state** (leptons+jets) with QCD corrections: NLO+PS methods available (+ **ongoing developments**)
  - full assessment of TH uncertainties missing ( $\rightarrow$  NLL showers will play crucial role)

# $t\bar{t}$ measurements: status

[CMS-PAS-TOP-20-006]

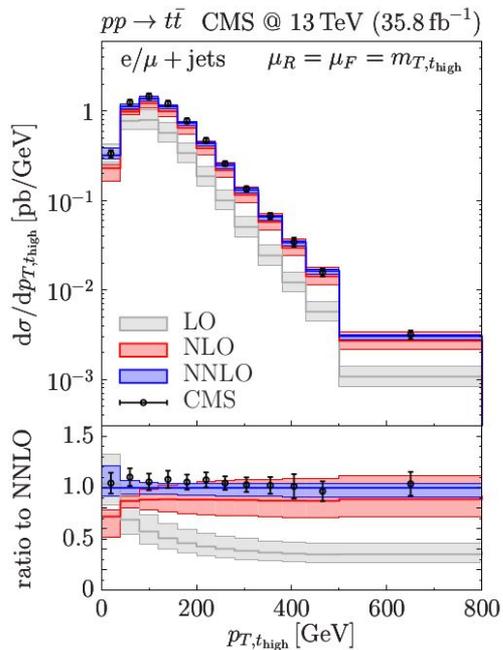


- TH in this plot: NLO+PS, NNLO, NNLO+PS (and aN3LO)

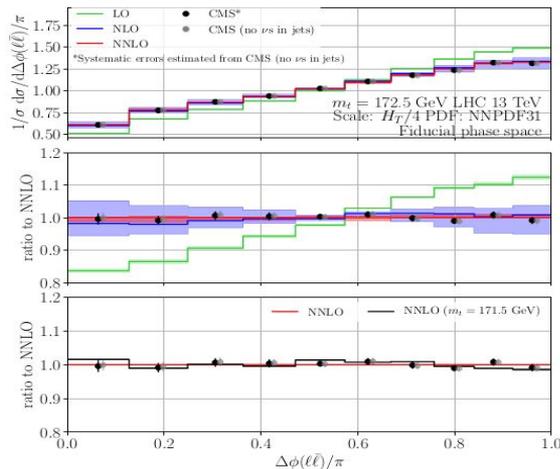
- scale choice:  $\mu=H_T/4=(m_T(t)+m_T(\bar{t}))/4$   
studied in [1606.0350]

- NNLO corrections  $\rightarrow$  fix long-standing issue with  $p_T(t)$

# $t\bar{t}$ (differential) parton-level computations

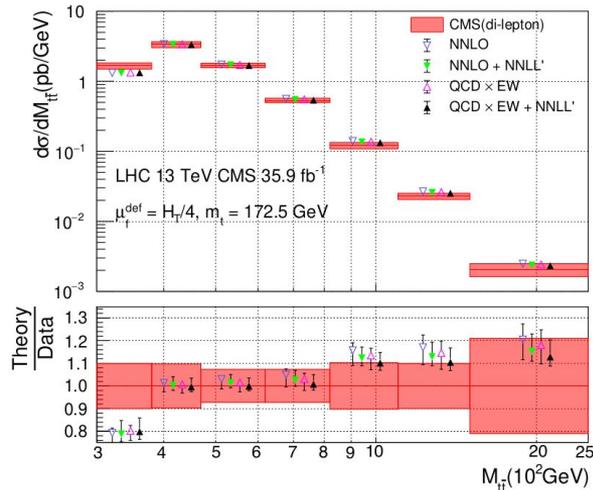


[Catani, Devoto, et al. '19]



[Czakon, Mitov, Poncelet '20]

$$d\sigma = d\sigma_{t\bar{t}} \times \frac{d\Gamma_t}{\Gamma_t} \times \frac{d\Gamma_{\bar{t}}}{\Gamma_{\bar{t}}}$$



[Czakon et al. 1901.08281]

- NNLO known for prod & decay
- decay: NWA, spin correlations
- NNLO x NNLO x NNLO

- EW/QCD combination

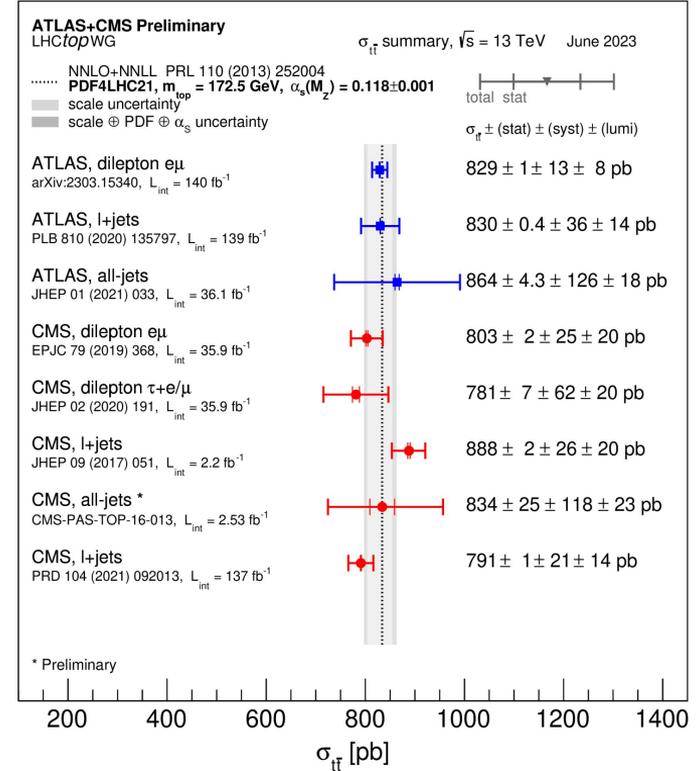
$$\Sigma_{\text{QCD} \times \text{EW}} \equiv \Sigma_{\text{QCD} + \text{EW}} + (K_{\text{QCD}}^{\text{NLO}} - 1) \times \Sigma_{\text{NLO EW}}$$

# Parton-level results: summary

- QCD corrections: clear picture, good perturbative convergence
- different scale choices possible:  $H_T$ -like seems to do a good job for distributions
  - insights on hardness vs. kinematics in  $t\bar{t}$ :  
[Caola,Dreyer et al. 2101.06068]
- Full NLO EW available [Denner,Pellen '16], possible to validate different approximations.
- Combination of EW and QCD corrections: done for stable top quarks (possible to supplement it with resummation(s))
- Residual uncertainty (size, PDF vs  $\mu$ ) depends on observables



Foundations on which MC developments are based



[ATL-PHYS-PUB-2023-014]

# MC event generators: status

- Several NLO+PS generators available: **POWHEG BOX**, **MG5\_aMC@NLO**, **Sherpa**, **Herwig** (through `Matchbox`).
- Approximations made to include decays and off-shell effects (e.g. `ttb_NLO_dec: prod@NLO x decay@NLO`)
- Exact MC simulation for decay, offshell effects and interferences:  
**bb41 NLO+PS** generator [Ježo,Lindert et al. '16]
  - **NEW: NLOPS matching for off-shell  $tt+Wt$  production with semileptonic decays** [Ježo,Lindert,Pozzorini '23]
- Multijet merging up to 2 jets @ NLO+PS [MEPS@NLO, FxFx, '12-'13]
- **$t\bar{t}$  NNLO+PS results available** [Mazzitelli,Monni,Nason,ER,Wiesemann,Zanderighi '20-'21]
- **NNLO+PS available for  $b\text{-bar}$ , ongoing work on  $t\bar{t}+H$**  [Mazzitelli,Wiesemann,Ratti,Zanderighi '23/Mazzitelli,Wiesemann WIP]

# NNLO+PS: goals and accuracy

- ▶ **NNLO accuracy** for observables inclusive on radiation.  $[d\sigma/dy_F]$
- ▶ **NLO(LO) accuracy** for  $F + 1(2)$  jet observables (in the hard region).  $[d\sigma/dp_{T,j_1}]$ 
  - appropriate scale choice for each kinematics regime
- ▶ **Sudakov resummation** from the Parton Shower (PS)  $[\sigma(p_{T,j} < p_{T,veto})]$
- ▶ preserve the PS accuracy (leading log - LL)
  - possibly, no merging scale required.

- ▶ methods: **reweighted MINLO'** (“**NNLOPS**”) [Hamilton, et al. '12,'13,...], **UNNLOPS** [Höche, Li, Prestel '14,...] / [Plätzer '12] **Geneva** [Alioli, Bauer, et al. '13,'15,'16,...], **MINNLO<sub>PS</sub>** [Monni, Nason, ER, Wieseemann, Zanderighi '19,...]

[Notation: From this point,  $X = \sum_k \left(\frac{\alpha_S}{2\pi}\right)^k [X]^{(k)}$ ]

## top pair-production @ NNLO+PS: MiNNLO (I)

- ▶ from  $p_T$  resummation, differential cross section for  $F+X$  production can be written as:

$$\frac{d\sigma}{dp_T d\Phi_F} = \frac{d}{dp_T} \left\{ \mathcal{L}(\Phi_F, p_T) \exp(-\tilde{S}(p_T)) \right\} + R_{\text{finite}}(p_T)$$

$$\mathcal{L}(\Phi_F, p_T) \ni \{H^{(1)}, H^{(2)}, C^{(1)}, C^{(2)}, (G^{(1)} \cdot G^{(1)})\} \quad R_{\text{finite}}(p_T) = \frac{d\sigma_{\text{FJ}}}{d\Phi_F dp_T} - \frac{d\sigma^{\text{sing}}}{d\Phi_F dp_T}$$

- ▶ recast it, to match the POWHEG  $\bar{B}^{(\text{FJ})}(\Phi_{\text{FJ}})$

$$\frac{d\sigma}{d\Phi_F dp_T} = \exp[-\tilde{S}(p_T)] \left\{ D(p_T) + \frac{R_{\text{finite}}(p_T)}{\exp[-\tilde{S}(p_T)]} \right\}$$

$$D(p_T) \equiv -\frac{d\tilde{S}(p_T)}{dp_T} \mathcal{L}(p_T) + \frac{d\mathcal{L}(p_T)}{dp_T} \quad \tilde{S}(p_T) = \int_{p_T}^Q \frac{dq^2}{q^2} \left[ A_f(\alpha_S(q)) \log \frac{Q^2}{q^2} + B_f(\alpha_S(q)) \right]$$

- ▶ expand the **above integrand** in power of  $\alpha_S(p_T)$ , keep the terms that are needed to get NLO<sup>(F)</sup> & NNLO<sup>(F)</sup> accuracy, when integrating over  $p_T$

# top pair-production @ NNLO+PS: MiNNLO (II)

$$\frac{d\bar{B}(\Phi_{\text{FJ}})}{d\Phi_{\text{FJ}}} = \exp[-\tilde{S}(p_{\text{T}})] \left\{ \frac{\alpha_{\text{S}}(p_{\text{T}})}{2\pi} \left[ \frac{d\sigma_{\text{FJ}}}{d\Phi_{\text{FJ}}} \right]^{(1)} \left( 1 + \frac{\alpha_{\text{S}}(p_{\text{T}})}{2\pi} [\tilde{S}(p_{\text{T}})]^{(1)} \right) + \left( \frac{\alpha_{\text{S}}(p_{\text{T}})}{2\pi} \right)^2 \left[ \frac{d\sigma_{\text{FJ}}}{d\Phi_{\text{FJ}}} \right]^{(2)} + [D(p_{\text{T}})]^{(\geq 3)} F_{\ell}^{\text{corr}}(\Phi_{\text{FJ}}) \right\}$$

$$- [D(p_{\text{T}})]^{(\geq 3)} = \underbrace{-\frac{d\tilde{S}(p_{\text{T}})}{dp_{\text{T}}} \mathcal{L}(p_{\text{T}})}_D + \frac{d\mathcal{L}(p_{\text{T}})}{dp_{\text{T}}} - \frac{\alpha_{\text{S}}(p_{\text{T}})}{2\pi} [D(p_{\text{T}})]^{(1)} - \left( \frac{\alpha_{\text{S}}(p_{\text{T}})}{2\pi} \right)^2 [D(p_{\text{T}})]^{(2)}$$

-  $F_{\ell}^{\text{corr}}(\Phi_{\text{FJ}})$ : projection  $\rightarrow$  recover  $[D(p_{\text{T}})]^{(\geq 3)}$  when integrating over  $\Phi_{\text{FJ}}$  at fixed  $(\Phi_{\text{F}}, p_{\text{T}})$

. The second radiation is generated by the usual POWHEG mechanism.

$$d\sigma = \bar{B}(\Phi_{\text{FJ}}) d\Phi_{\text{FJ}} \left\{ \Delta_{\text{pwg}}(\Lambda_{\text{pwg}}) + d\Phi_{\text{rad}} \Delta_{\text{pwg}}(p_{\text{T,rad}}) \frac{R(\Phi_{\text{FJ}}, \Phi_{\text{rad}})}{B(\Phi_{\text{FJ}})} \right\}$$

. if emissions are strongly ordered, same emission probabilities as in  $k_t$ -ordered shower  
 $\rightarrow$  LL shower accuracy preserved

# top pair-production @ NNLO+PS: MiNNLO for $t\bar{t}$

- ▶ Starting point: resummation formula for  $t\bar{t}$  transverse momentum.

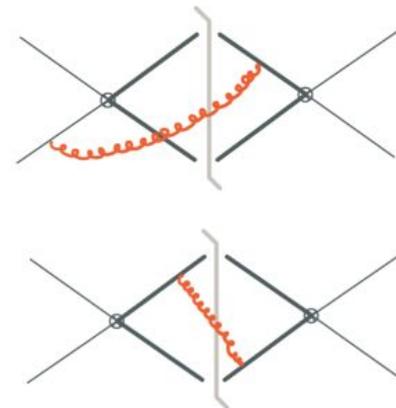
[Catani, Grazzini, Torre '14]

Very schematically:

$$d\sigma_{\text{res}}^F \sim \frac{d}{dp_T} \left\{ e^{-S} \text{Tr}(\mathbf{H}\mathbf{\Delta}) (C \otimes f) (C \otimes f) \right\}$$

$$S = - \int \frac{dq^2}{q^2} \left[ \frac{\alpha_s(q)}{2\pi} (A^{(1)} \log(M/q) + B^{(1)}) + \frac{\alpha_s^2(q)}{(2\pi)^2} (A^{(2)} \log(M/q) + B^{(2)}) + \dots \right]$$

$$\text{Tr}(\mathbf{H}\mathbf{\Delta}) = \langle M | \mathbf{\Delta} | M \rangle, \quad \mathbf{\Delta} = \mathbf{V}^\dagger \mathbf{D} \mathbf{V}, \quad \mathbf{V} = \exp \left\{ - \int \frac{dq^2}{q^2} \left[ \frac{\alpha_s(q)}{2\pi} \mathbf{\Gamma}_t^{(1)} + \frac{\alpha_s^2(q)}{(2\pi)^2} \mathbf{\Gamma}_t^{(2)} \right] \right\}$$



- ▶ With some approximations (respecting our goal), terms due to soft interference can be rearranged so that the “resummation” can be eventually recasted as:

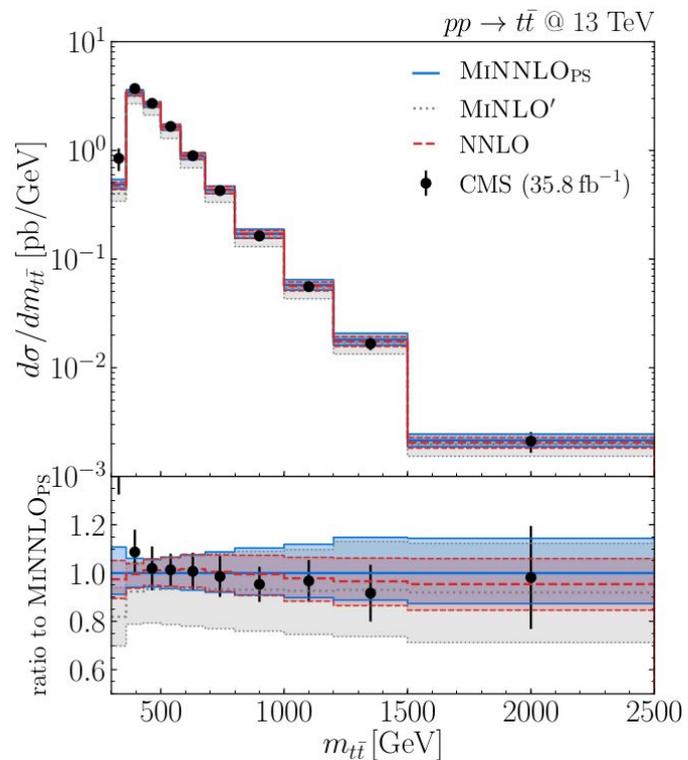
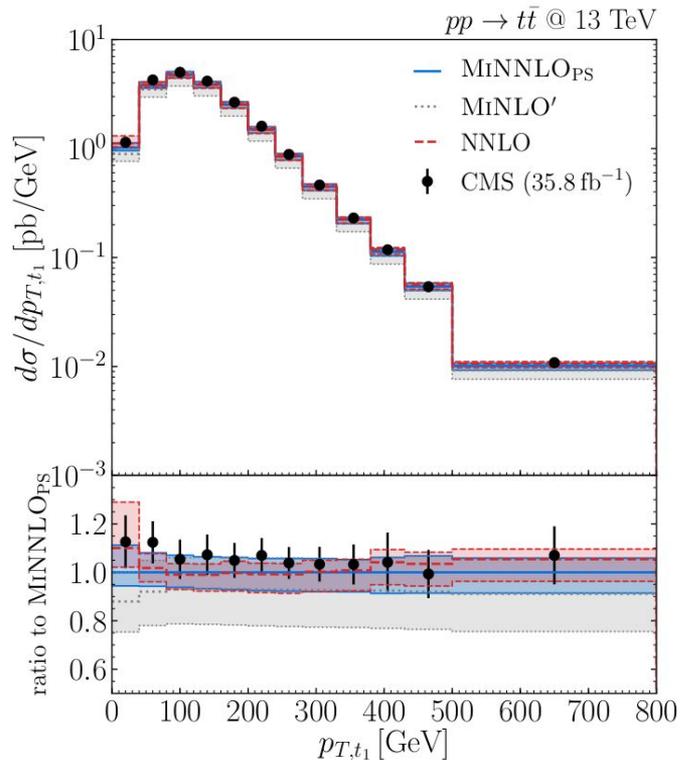
$$d\sigma_{\text{res}}^F \sim \frac{d}{dp_T} \left\{ \sum_{i \in \text{colours}} e^{-\bar{S}_i} c_i \underbrace{\overline{H(C \otimes f)(C \otimes f)}}_{\equiv \overline{\mathcal{L}}_i} \right\} + \mathcal{O}(\alpha_S^5)$$

inputs from [Catani, Devoto, Grazzini, Kallweit, Mazzitelli + Sargsyan '19]

paper: [Catani, Devoto, Grazzini, Mazzitelli '23]

- ▶ Each term has the “same structure” as in the color-singlet case!

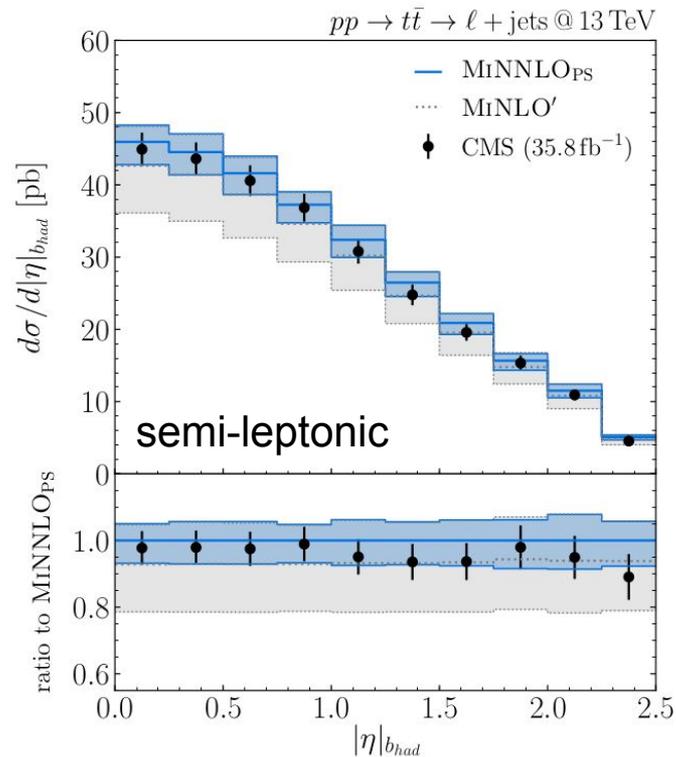
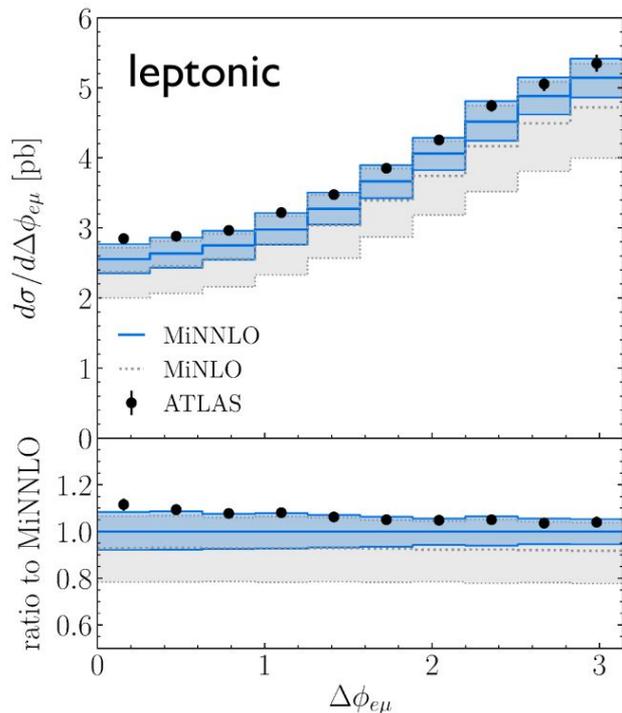
# top pair-production @ NNLO+PS: results I



- nice agreement with NNLO (and with data - both ATLAS and CMS).  $\mu_{\text{core}} = H_T/4$

-  $m(t\bar{t})$ , 1st bin: finite-width effects + non-relativistic effects + QED

# top pair-production @ NNLO+PS: results II



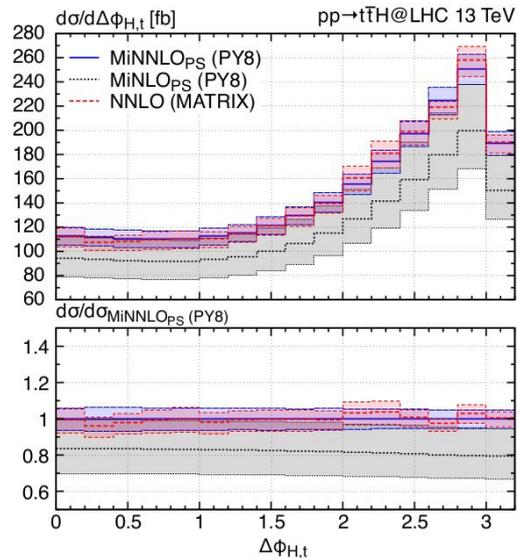
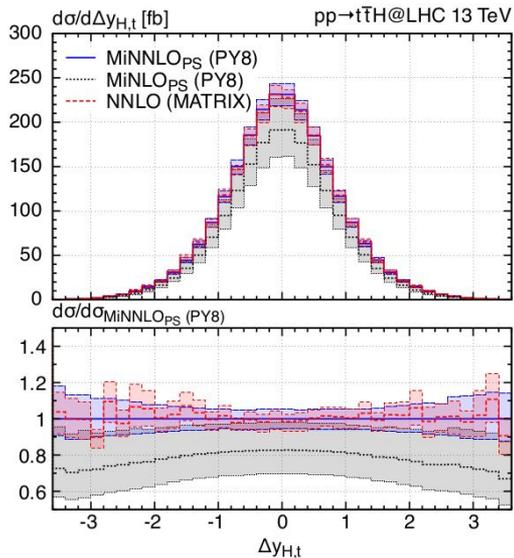
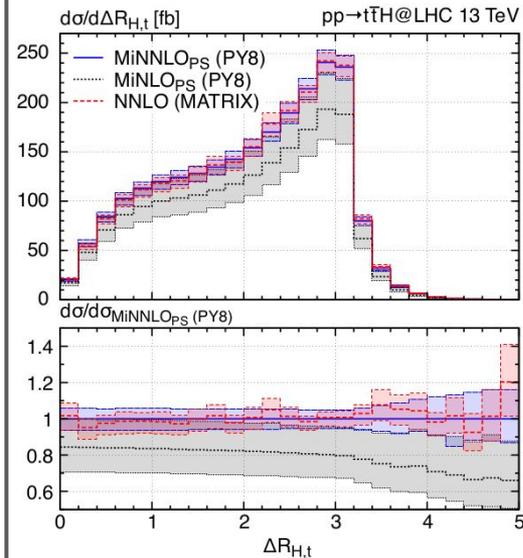
- implemented top-quark decays @ tree level + approximated off-shell effects
- NB: if analysis probes off-shell/non-resonant regions → **bb41 NLO+PS** should be method of choice

# ttH production @ NNLO+PS

[Mazzitelli, Wiesemann WIP]

PRELIMINARY

[slide from Wiesemann's talk at SM@LHC]



Marius Wiesemann (MPP Munich)

Heavy Flavour production at the LHC with NNLO+PS

July 14, 2023

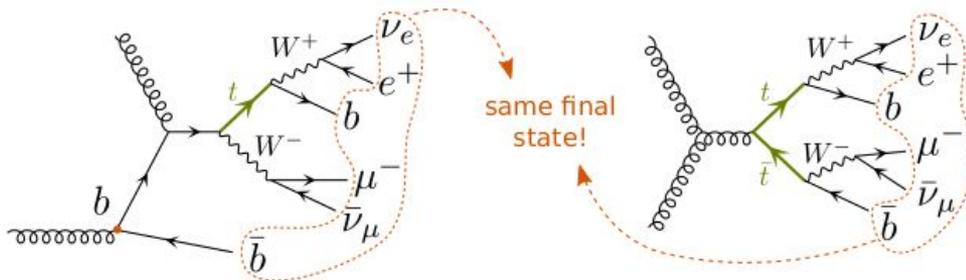
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- soft function computation extended (arbitrary kinematics):
- 2-loops using “soft Higgs” approximation

[Devoto, Mazzitelli WIP]

[Catani et al. 2210.07846 → talk by S. Devoto]

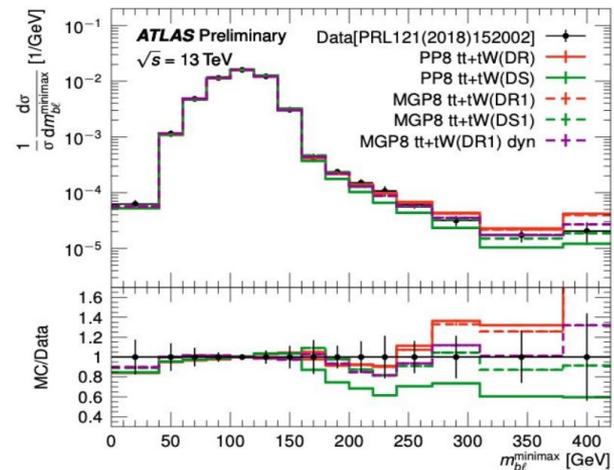
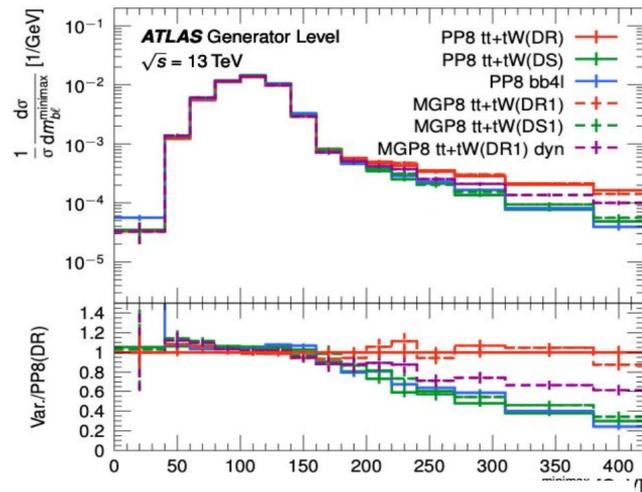
# $t\bar{t}$ / $Wt$ : EXP/TH status



[diagrams from slides by T. Ježo]

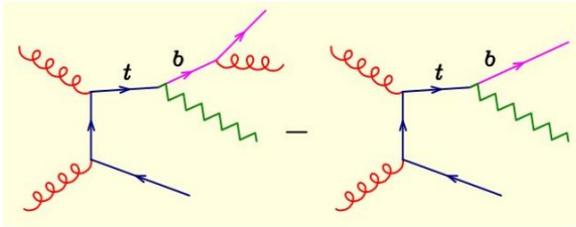
- Interference  $t\bar{t}/Wt$ : **long-standing issue**
  - Approximations made (DR/DS) to treat processes separately, with unstable resonances in final state
  - DR/DS differences might become an issue in BSM search regions
- **bb41 NLOPS** generator in **POWHEG-BOX-RES**: interferences/off-shell effects, top decay at NLO
  - resonant-aware methods also in **MG5\_aMC@NLO**

[ATL-PHYS-PUB-2021-042]



$$m_{b\ell}^{\text{minimax}} = \min\{\max(m_{b_1, \ell_1}, m_{b_2, \ell_2}), \max(m_{b_2, \ell_1}, m_{b_1, \ell_2})\} \quad 14$$

# Off-shell $t\bar{t} + tW$ : status



$$d\sigma = d\Phi_{\text{rad}} \bar{B}(\Phi_B) \frac{R(\Phi_B, \Phi_{\text{rad}})}{B(\Phi_B)} \times \exp \left[ - \int \frac{R(\Phi_B, \Phi_{\text{rad}})}{B(\Phi_B)} d\Phi_{\text{rad}} \right]$$

$\Phi_B \rightarrow (\Phi_B, \Phi_{\text{rad}})$  mapping does not preserve virtuality  
 $\Rightarrow R/B$  can become large also far from collinear singularity



- NW limit spoiled
- BW shape distorted

## “Resonant-aware” NLO+PS methods

[Ježo, Nason ‘15 / Frederix et al. ‘16]

- project full MEs onto “resonance histories”:  $B = \sum_{f_b} B_{f_b}$ , where  $B_{f_b} \equiv \frac{P^{f_b}(\Phi_B)}{\sum_{f'_b} P^{f'_b}(\Phi_B)} B(\Phi_B)$   
 $P^{f_b}(\Phi_B)$  (products of) Breit-Wigner functions  $\Leftrightarrow$  resonance history  $f_b$
- each term (Born- and real-like) is attributed to an unique resonance history
  - virtuality-preserving mappings
  - POWHEG radiation can be assigned a resonance  $\rightarrow$  (up to) 1 emission per resonance

# Off-shell $t\bar{t}$ + $tW$ : status $\rightarrow$ recent progress

- **bb41 NLO+PS** [[1607.04538](#)]: generator available for leptonic decays  
(used also for TH studies, e.g. [[Ferrario Ravasio, Ježo et al. '18](#)])
- **NEW**: [[Ježo, Lindert, Pozzorini 2307.15653](#)]: improvements + semi-leptonic decays
  1. width effects
  2. improved resonance histories

# Off-shell $t\bar{t}$ + $tW$ : recent progress

[Ježo,Lindert,Pozzorini 2307.15653]

## 1. width effects:

NWA: 
$$d\sigma_{\text{prod}\times\text{dec}} = d\sigma \frac{d\Gamma}{\Gamma}$$

naive: 
$$d\sigma_{\text{prod}\times\text{dec}}^{\text{NLO}} = d\sigma_0 \frac{d\Gamma_0}{\Gamma_{\text{NLO}}} + d\sigma_1 \frac{d\Gamma_0}{\Gamma_{\text{NLO}}} + d\sigma_0 \frac{d\Gamma_1}{\Gamma_{\text{NLO}}} \quad \int_{\text{dec}} d\sigma_{\text{prod}\times\text{dec}}^{\text{NLO}} = d\sigma_0 + d\sigma_1 - d\sigma_1 \frac{\Gamma_1}{\Gamma_{\text{NLO}}}$$

consistent [0907.3090]: 
$$\int_{\text{dec}} d\sigma_{\text{prod}\times\text{dec}}^{\text{NLO}_{\text{exp}}} = d\sigma_0 \left[ \left(1 - \frac{\Gamma_1}{\Gamma_0}\right) \int_{\text{dec}} \frac{d\Gamma_0}{\Gamma_0} + \int_{\text{dec}} \frac{d\Gamma_1}{\Gamma_0} \right] + d\sigma_1 \int_{\text{dec}} \frac{d\Gamma_0}{\Gamma_0} = d\sigma_0 + d\sigma_1$$

multi resonances: 
$$d\sigma_{\text{prod}\times\text{dec}}^{\text{NLO}_{\text{exp}}} = \left( \prod_{r \in \mathcal{R}} \frac{\Gamma_{r,\text{NLO}}}{\Gamma_{r,0}} \right) d\sigma_{\text{prod}\times\text{dec}}^{\text{NLO}} - \left( \sum_{r \in \mathcal{R}} \frac{\Gamma_{r,1}}{\Gamma_{r,0}} \right) d\sigma_{\text{prod}\times\text{dec}}^{\text{LO}}$$

offshell  $\rightarrow$  NLO+PS: 
$$d\sigma_{\text{off-shell}}^{\text{NLO}_{\text{exp}}} = \left( \prod_{r \in \mathcal{R}} \frac{\Gamma_{r,\text{NLO}}}{\Gamma_{r,0}} \right) \left[ d\sigma_{\text{off-shell}}^{\text{NLO}} - \left( \sum_{r \in \mathcal{R}} \frac{\Gamma_{r,1}}{\Gamma_{r,0}} \right) d\sigma_{\text{off-shell}}^{(0)} \right]$$

$$\delta\kappa_{\text{spurious}}^{t\bar{t}+X} \simeq -2 \frac{d\sigma_1}{d\sigma_0} \frac{\Gamma_{t,1}}{\Gamma_{t,0}}$$

17%

- last equation: used in new **bb41-d1/s1** generators

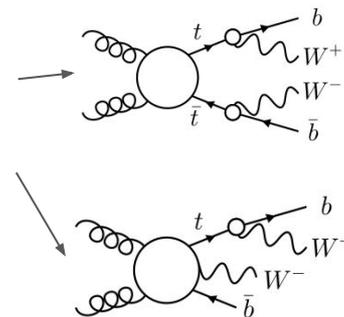
# Off-shell $t\bar{t}$ + $tW$ : recent progress

[Ježo,Lindert,Pozzorini 2307.15653]

## 2. improved resonance histories:

- original **bb41**:  $\{t\bar{t}, Z\}$
- **bb41-d1/s1**:  $\{t\bar{t}, tW^+, tW^-\}$   
 → expect better treatment of “decay” emissions

resonance history	production subprocess
$t\bar{t}$	$pp \rightarrow t\bar{t}$
$tW^-$	$pp \rightarrow tW^-\bar{b}$
$\bar{t}W^+$	$pp \rightarrow \bar{t}W^+b$
$Z$	$pp \rightarrow Z/H + b\bar{b}$



$$\rho_{t\bar{t}}^{(\text{hist})}(\Phi_B)|_{\text{ME}} = |\mathcal{A}_{t\bar{t}}|^2,$$

- ME-based projectors:  $\rho_{\bar{t}W^+}^{(\text{hist})}(\Phi_B)|_{\text{ME}} = |\mathcal{A}_{\bar{t}W^+}|^2$

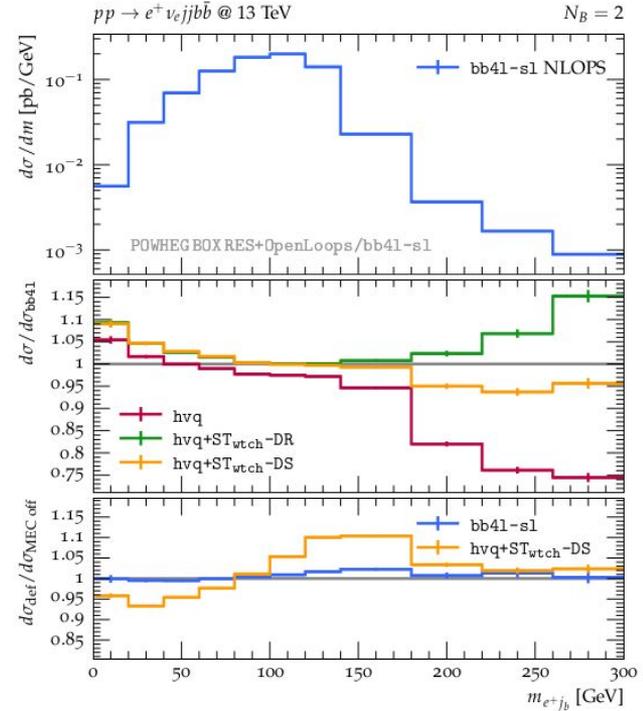
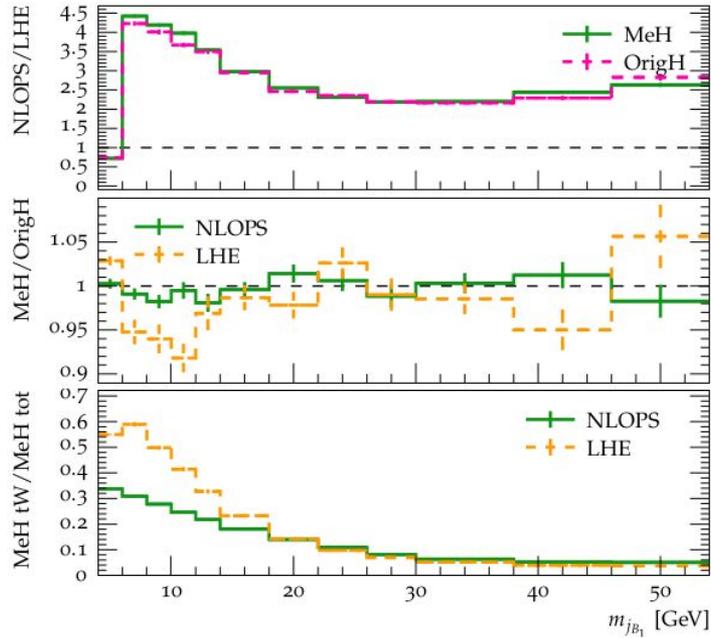
$$\rho_{tW^-}^{(\text{hist})}(\Phi_B)|_{\text{ME}} = |\mathcal{A}_{tW^-}|^2$$

- several variations, new mappings, thorough validation

$$\sigma_{t\bar{t}} = \lim_{\xi_t \rightarrow 0} \left( \xi_t^2 \sigma_{\text{bb41}} \Big|_{\Gamma_t \rightarrow \xi_t \Gamma_t} \right)$$

# Off-shell $t\bar{t}$ + $tW$ : results

[Ježo,Lindert,Pozzorini 2307.15653]



- LHS: resonance histories ambiguities (in off-shell regions\*)  $\rightarrow$  negligible after PS

- RHS: even when interference effects large,  $\text{hvq}+\text{wtch-DS} \approx \text{bb41}$

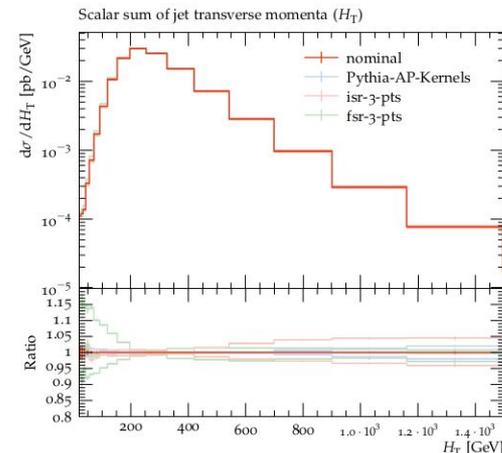
- naive width  $\rightarrow \left. \frac{\sigma_{\text{bb41-s1}}}{\sigma_{\text{hvq+ST}}} \right|_{\text{no } 1/\Gamma_t \text{ expansion}} = 1.074$

\*  $Q_{\text{off-shell}} = \max\{|Q_t - m_t|, |Q_{\bar{t}} - m_t|\} > 60 \text{ GeV}$  <sup>19</sup>

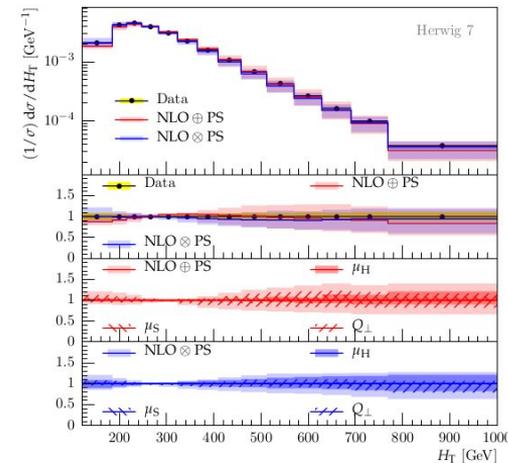
# TH uncertainties in MC generators

## Possible TH uncertainties:

- scale variation (hard matrix elements)
  - matching uncertainties (e.g.  $h_{\text{damp}}$  in POWHEG, hard veto scale in MC@NLO,...)
  - change matching scheme / shower [e.g. [Matchbox study, 1810.06493](#)]
  - other “shower-related” pQCD uncertainties can be probed
  - recoil scheme
  - .....
  - non-perturbative parameters & tuning
- 
- possible to include such variations (→ certainly not the ultimate solution)
  - within the current paradigms for matching & merging, rethinking needed for some of the above items, once matching to NLL parton showers will be achieved (Panscales, Alaric, FHP, ....)



[LH19 (2003.01700)]



[1810.06493]

# Conclusions

- MC event generators play a role in ~ all EXP analysis
  - data to parton-level, background for BSM, measure parameters
- processes with top quarks: NNLO+PS available (with MiNNLOPS method), approximate decays
- full off-shell final state (leptons+jets): NLO+PS “resonance-aware” methods available
  
- full assessment of TH uncertainties missing (→ NLL showers will play crucial role)
- $t\bar{t}+X$ : once NNLO known (exactly or approximately) → NNLO+PS possible
  
- Other approaches: *Geneva* for  $t\bar{t}$ ...[[2111.03632](#) (0-jettiness resummation)]

*Thank you for your attention!*

Backup slides

# top pair-production @ NNLO+PS: MiNNLO for $t\bar{t}$

$$\frac{d\sigma}{d\vec{q}_\perp d\Phi_F} \sim \sum_f |M_{f\bar{f} \rightarrow t\bar{t}}^{(0)}|^2 \int \frac{d^2\vec{b}}{(2\pi)^2} e^{i\vec{b} \cdot \vec{q}_\perp} e^{-R_f(b)} \text{Tr}(\mathbf{H}_f \Delta_{\text{soft}}) \sum_{i,j} (C_{fi} \otimes h^{[i]})(C_{\bar{f}j} \otimes h^{[j]})$$

$$\text{Tr}(\mathbf{H}_f \Delta_{\text{soft}}) = \frac{\langle M_{f\bar{f}} | \Delta | M_{f\bar{f}} \rangle}{|M_{f\bar{f}}^{(0)}|^2}, \quad \Delta = \mathbf{V}^\dagger \mathbf{D} \mathbf{V}$$

- LL+NNLO accuracy: azimuthally averaged distribution becomes

$$\mathbf{V} = \mathcal{P} \exp \left\{ - \int_{\frac{b_0^2}{b^2}}^{M_{t\bar{t}}^2} \frac{dq^2}{q^2} \Gamma_t(\Phi_{t\bar{t}}, \alpha_s(q)) \right\}$$

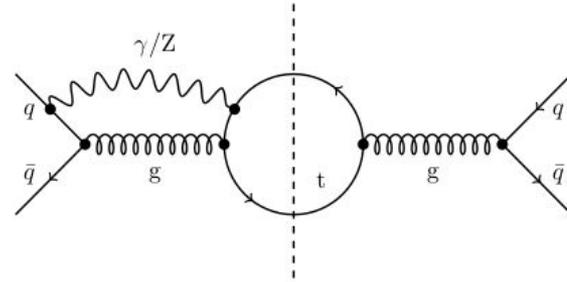
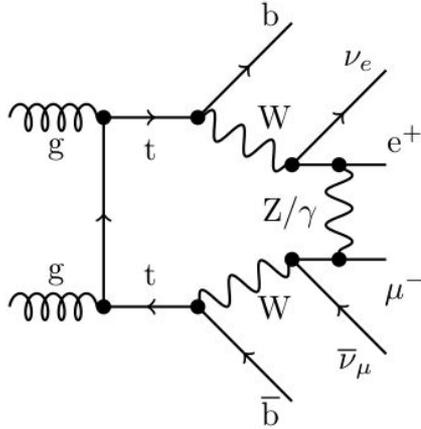
$$\left[ \frac{d\sigma}{d\vec{q}_\perp d\Phi_F} \right]_\phi \sim \frac{d}{dq_\perp} \left[ \sum_f e^{-S_f(q_\perp)} \langle M_{f\bar{f}}^{(0)} | (\mathbf{V}_{\text{NLL}})^\dagger \mathbf{V}_{\text{NLL}} | M_{f\bar{f}}^{(0)} \rangle [\text{Tr}(\mathbf{H}_f \mathbf{D}_{\text{soft}}) \sum_{i,j} (C_{fi} \otimes h^{[i]})(C_{\bar{f}j} \otimes h^{[j]})] \phi \right]$$

$$B(\alpha_s) = \frac{\alpha_s}{2\pi} B^{(1)} + \frac{\alpha_s^2}{(2\pi)^2} B^{(2)}$$

$$\sim \frac{\langle M_{f\bar{f}}^{(0)} | \mathbf{D}^{(1)} | M_{f\bar{f}}^{(0)} \rangle}{|M_{f\bar{f}}^{(0)}|^2} \left( (C_{fi}^{(1)} \otimes f_i)(f_{\bar{f}}) \right)$$

- diagonalization of  $\mathbf{V}_{\text{NLL}}$  → recast as sum of “colour-singlet-like” terms

# EW corrections: from total cross sections....

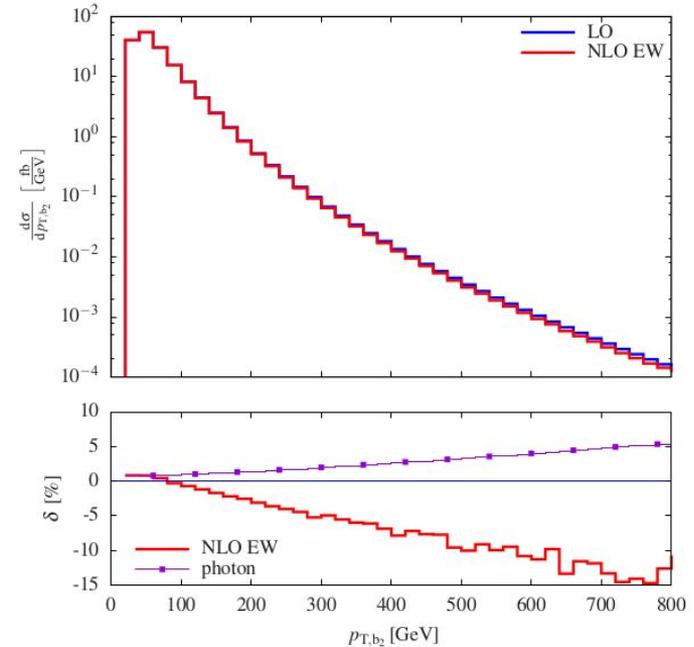
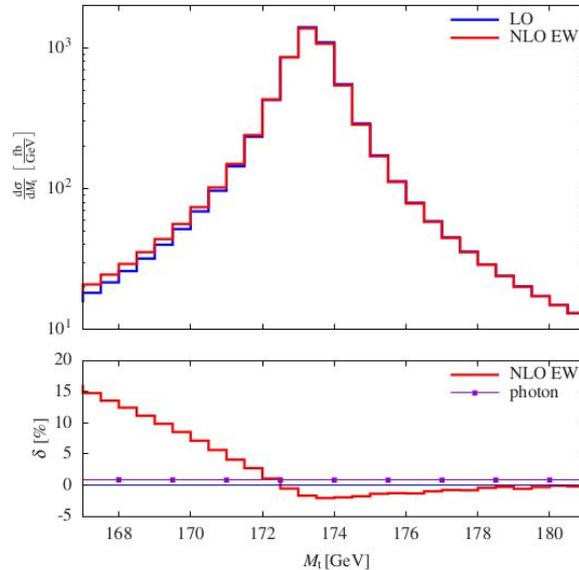


- fully exclusive results (2 lep channel)
- NLO EW =  $O(\alpha_s^2 \alpha^5)$
- total x-section: EW corrections < 1%
- no uncertainty on EW corrections  
(but irrelevant for total x-section)

Ch.	$\sigma_{\text{LO}}$ [fb]	$\sigma_{\text{NLO EW}}$ [fb]	$\delta$ [%]
gg	2824.2(2)	2834.2(3)	0.35
$q\bar{q}$	375.29(1)	377.18(6)	0.50
$gq(/\bar{q})$		0.259(4)	
$\gamma g$		27.930(1)	
pp	3199.5(2)	3211.7(3)	0.38

[Denner,Pellen '16]

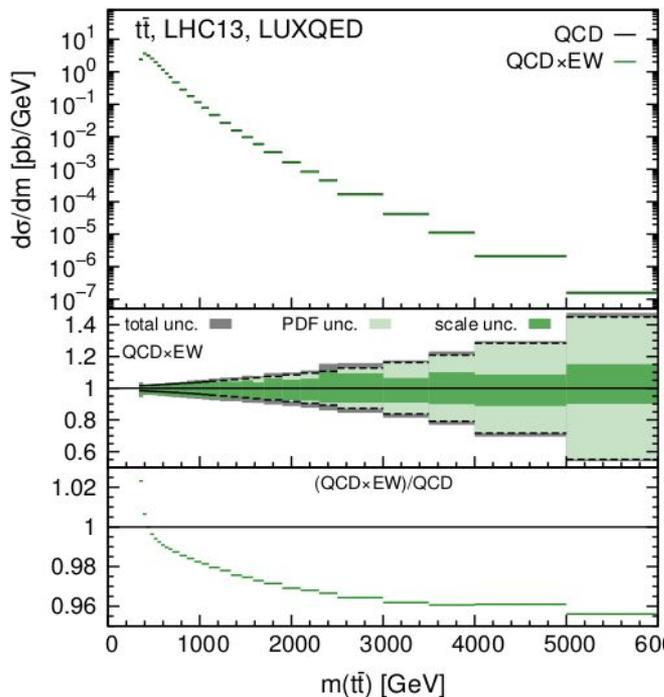
## ...to differential distributions



- corrections  $\sim 10\text{-}15\%$  (radiative tails (left) / Sudakov logs (right))
- here: no approximations. Possible to study “DP approximations” for tops or Ws.
- scales fixed at top mass (enter in  $\alpha_s$  & PDFs)
- TH uncertainties: needed? Uncharted territory: change of scheme (?),...

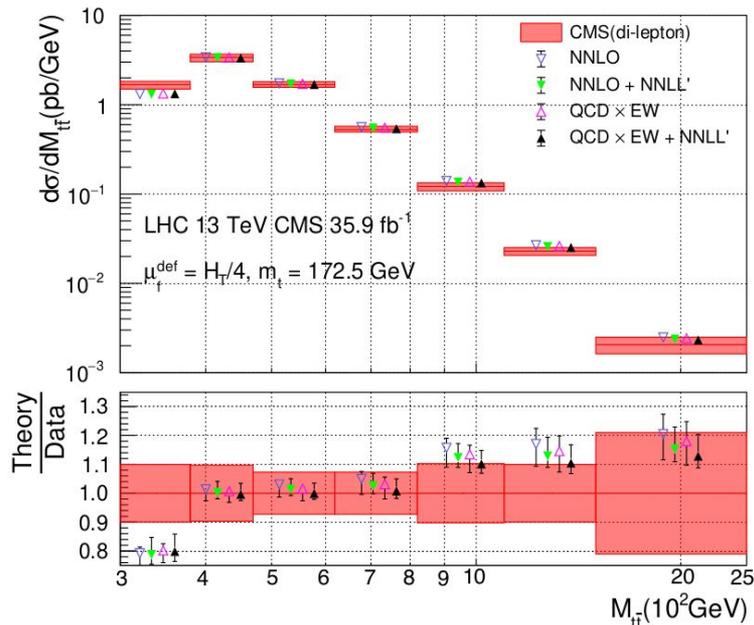
# EW and QCD combined results

[1705.04105]



preferred combination: multiplicative  $\Rightarrow$

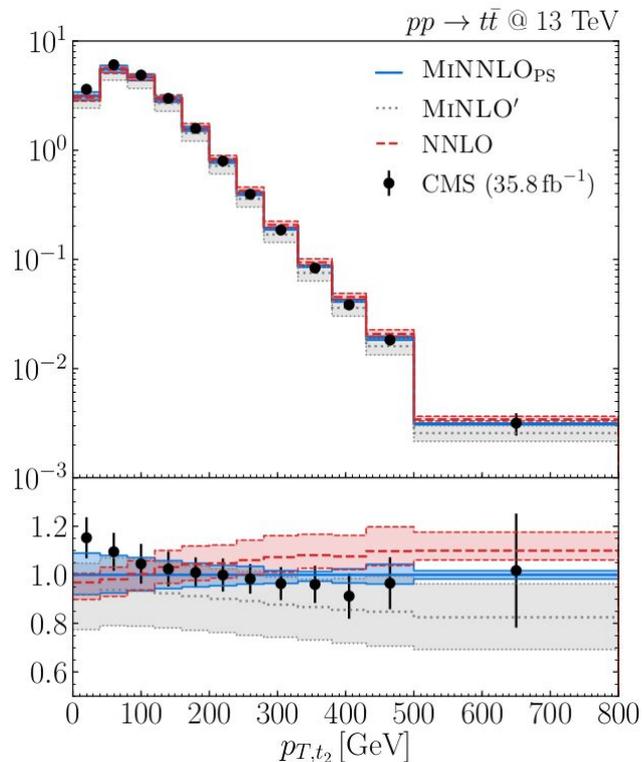
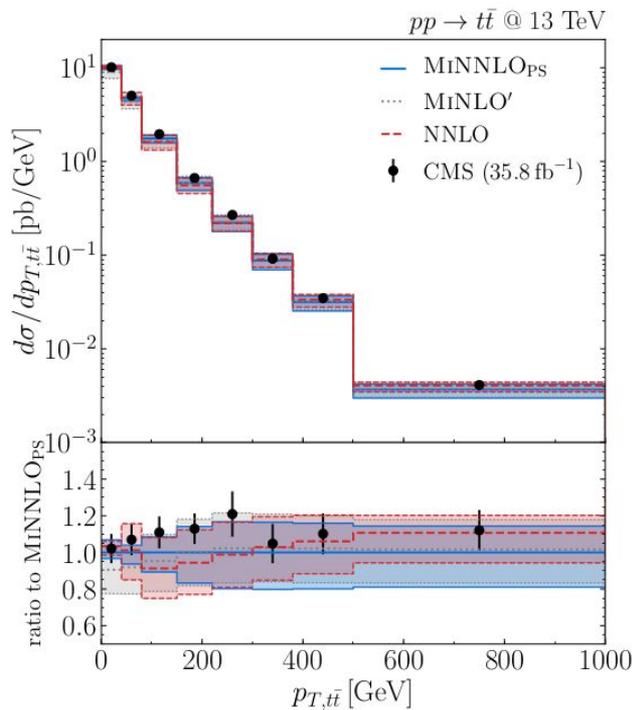
[1901.08281]



$$\Sigma_{\text{QCD}\times\text{EW}} \equiv \Sigma_{\text{QCD}+\text{EW}} + (K_{\text{QCD}}^{\text{NLO}} - 1) \times \Sigma_{\text{NLO EW}}$$

- stabilize scale uncertainty of  $\Sigma_{\text{NLO EW}} = O(\alpha_s^2 \alpha)$
- $\Sigma_{\text{mixed}} = O(\alpha_s^3 \alpha) \approx \Sigma_{\text{NLO QCD}} \Sigma_{\text{NLO EW}} / \Sigma_{\text{LO QCD}}$
- correct in regime “soft gluon” + “Sudakov log”

# top pair-production @ NNLO+PS: results III



- nice agreement with NNLO (and with data - both ATLAS and CMS).  $\mu_{\text{core}} = H_T/4$

# Resonant-aware NLO+PS: details I

1. complete matrix elements for  $W^+W^-b\bar{b}$ : need to project each partonic subprocess onto all possible “resonance histories”:

- each contribution should be dominated by a single resonance history:

$$B = \sum_{f_b} B_{f_b}, \quad \text{where} \quad B_{f_b} \equiv \frac{P^{f_b}(\Phi_B)}{\sum_{f'_b} P^{f'_b}(\Phi_B)} B(\Phi_B)$$

$P^{f_b}(\Phi_B)$  (products of) Breit-Wigner functions  $\Leftrightarrow$  resonance history  $f_b$

- for real contributions, split also according to compatible FKS regions

$\Rightarrow$  a term  $R_{\alpha_r}$  is dominant if the collinear partons of region  $\alpha_r$  have the smallest  $k_T$ , and the corresponding resonance history is the closest to its mass shell.

2. each term (Born-like and real) is attributed to an unique resonance history

- virtuality-preserving mappings between  $\Phi_B$  and  $(\Phi_B, \Phi_{\text{rad}})$  can be used
- POWHEG radiation(s) can now be assigned to a resonance
- (& other technical but crucial subtleties...)

# Resonant-aware NLO+PS: details II

“multiplicative POWHEG”: keep multiple emissions before showering



- by default POWHEG is additive: keeps only the hardest emission
- keep hard radiation and the emissions from all decaying resonances, then merge them into a single radiation phase space with several radiated partons, up to one for each resonance

$\Leftrightarrow \quad d\sigma = \bar{B}(\Phi_B) d\Phi_B \left[ \Delta(q_{\text{cut}}) + \sum_{\alpha} \Delta(k_T^{\alpha}) \frac{R_{\alpha}(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}))}{B(\Phi_B)} d\Phi_{\text{rad}} \right]$

$\Leftrightarrow \quad d\sigma = \bar{B}(\Phi_B) d\Phi_B \prod_{\alpha=\alpha_b, \alpha_{\bar{b}}, \alpha_{1SR}} \left[ \Delta_{\alpha}(q_{\text{cut}}) + \Delta_{\alpha}(k_T^{\alpha}) \frac{R_{\alpha}(\Phi_{\alpha}(\Phi_B, \Phi_{\text{rad}}^{\alpha}))}{B(\Phi_B)} d\Phi_{\text{rad}}^{\alpha} \right]$

► in the above case, the interface to parton shower becomes more complicated.

# TH uncertainties in MC generators: looking ahead

- plots from: [\[Hamilton et al. 2301.09645\]](#)
- thrust in  $e^+e^-$ : NLO+PS multiplicative matching + NLL shower
- dots: modified splitting function in hard region
- dashes:  $\mu_R$  scale variation (also in hard matrix elements)
- here matching fulfils NNDL accuracy (i.e. the same accuracy of a NLL resummation matched with NLO)

$$\Sigma(O < e^L) = h_1(\alpha_s L^2) + \sqrt{\alpha_s} h_2(\alpha_s L^2) + \alpha_s h_3(\alpha_s L^2) + \dots$$

