Lattice QCD in the Frontier of Electron-Ion Colliders

August 2024, Liverpool, UK









electron-ion colliders

shine light on atoms & proton



future electron-ion colliders

electron-ion collider in China



future circular collider @ CERN





Electron-Ion Collider @ Brookhaven National Laboratory science phase: 2033 –



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- Capability to achieve high polarization ($\sim 70\%$) of both electron and proton beams.
- Variable center-of-mass energies $\sqrt{s} = 20 100$ GeV, upgradable up to $\sqrt{s} = 140$ GeV, for e + p collisions. • Luminosity up to $10^{33} - 10^{34}$ cm⁻²s⁻¹ for e + p collisions.
- Capability for colliding ion beams from deuterons to heavy nuclei, such as gold, lead or uranium.
- Flexibility to accommodate more than one interaction regions.



Major research and innovation infrastructure investment announced **UK Research** and Innovation **UK-US collaboration**

Another project will receive £58.8 million from UKRI in a partnership with the US Department of Energy (DOE), to develop new detector and accelerator infrastructure to address fundamental questions on the nature of matter.

The technology will be built by:

- two STFC national laboratories, Daresbury Laboratory in Cheshire and the Rutherford Appleton Laboratory in Oxfordshire
- the universities of Birmingham, Brunel, Glasgow, Lancaster, Liverpool, Oxford and York
- the Cockcroft Institute for Accelerator Science and Technology in Cheshire

It will be installed at the Electron-Ion Collider (EIC), a major new particle accelerator facility at the Brookhaven National Laboratory in New York in the US.





Dark matter 27%

Visible

matter

5%

free Hydrogen & Helium: 4%

stars: 0.5% neutrinos: 0.3% heavy elements: 0.03%

68% Dark energy



Structure within the Atom

Quark Size < 10⁻¹⁹ m

C

Nucleus Size ≈ 10⁻¹⁴ m.

e⁻

Atom Size = 10⁻¹⁰ m

If the protons and neutrons in this picture were 10 cm across, then the quarks and electrons would be less than 0.1 mm in size and the entire atom would be about 10 km across.

d

Electron Size < 10⁻¹⁸ m

e⁻

Neutron and Proton Size ≈ 10⁻¹⁵ m

partonic origin of proton mass

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proton in the eyes of the light: partonic image

an effective description observed from an infinite-momentum frame

 x_i : fraction of the momentum of proton carried by a parton

In c.m. frame electron sees a Lorentz-contracted proton; parton's virtual life Lortentz-dialted

• $Q^2 \to \infty$, electron crosses proton in $t \to 0$: sees partons 'frozen' on approx. mass-shell

deep inelastic scattering: $e + P \longrightarrow e + X$

factorization ~ perturbative \bigotimes non-perturbative $\sigma(y) \sim c(y, x, \mu) \bigotimes f_{LC}(x, \mu)$

parton distribution function (PDF)

distribution, at scale μ , of longitudinal momentum fractions of a parton inside hadron moving with infinite momentum

science at EIC needs help from lattice QCD

Summary of the National Academy of Science report

"The scientific challenges that would unfold with EIC require a robust theory program, not simply to design and interpret experiments, but also to develop the broad implications in an understanding of the quantum world, both through analytic theory as well as through lattice QCD simulations on large-scale computers."

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nonperturbative partonic image of hadron

regularize QCD after taking the lightcone, $P_z \rightarrow \infty / z^2 \rightarrow 0$, limit

$$f(x,\mu) \sim \left\langle H(P_z) | \hat{O}(z^-,\mu) | H(P_z) \right\rangle$$

timelike separated bilocal operator

lattice QCD in EIC era —

how to 'see' a parton on the lattice?

partonic structures from lattice QCD

hadron at rest

spacelike separated bilocal operator

renormalize: scale μ

 $M(z^2,\mu) \sim \left\langle H(0) \left| \hat{O}(z,\mu) \right| H(0) \right\rangle$

fast-moving hadron

 $\left\langle H(0) \,|\, \hat{O}(z,\mu) \,|\, H(0) \right\rangle$

$\left\langle H(P_z) | \hat{O}(z^-, \mu) | H(P_z) \right\rangle$

 $\tilde{c}(y, x, \mu, P_z) \otimes f_{LOCD}(x, \mu)$

momentum space

nonperturbative objects on the lightcone, $f_{LC}(x, \mu)$, and from lattice QCD, $f_{LOCD}(x, \mu)$, shares same infrared singularities, i.e. governed by same evolution equations

factorization of $M(y, \mu, P_{\gamma}) \sim \text{ perturbative } \otimes \text{ non-perturbative}$

 $\tilde{c}(y, x, \mu, z^2) \otimes f_{LOCD}(x, \mu)$

position space

factorization: perturbative \otimes non-perturbative

 $M(y, \mu, P_{z}) \sim \tilde{\sigma}(y, x, \mu, P_{z}) \otimes \tilde{f}(x, \mu)$

- regularize QCD on a lattice, then $P_7 \rightarrow \infty / z^2 \rightarrow 0$
- opposite order of limits from light-cone quantization
- two limits don't commute
- difference is UV physics, can be taken care of through perturbative matching

 $M(y, \mu, z^2) \sim \tilde{\sigma}(y, x, \mu, z^2) \otimes \tilde{f}(x, \mu)$

parton physics

NNLO NLO LO $C(\mathcal{S},\mu) \sim \alpha_s^0(\mu) + \alpha_s(\mu) f\left(\ln[\mathcal{S}\mu]\right) + \alpha_s^2(\mu) f\left(\ln[\mathcal{S}\mu]\right) + \cdots$

lattice QCD

 $\mathcal{S} = 2xP_z, z^2$

isovector quark PDF of unpolarized proton at NNLO

physical quark masses

A. Hanlon et al., Phys. Rev. D107, 7, 074509 (2023)

valance quark PDF of pion at NNLO

Sullivan process @ EIC

physical quark masses, continuum-extrapolated

Y. Zhao et al., <u>Phys. Rev. Lett. 128, 14, 142003 (2022)</u>
X. Gao et al., <u>Phys. Rev. D06, 11, 114510 (2022)</u>

quark energy contributions to hadron masses

 $E_q(2 \text{ GeV})/m_p \approx 40\%$

$$-m_H \int_0^1 x f(x,\mu) \, dx$$

4.0 DNN-abAsymp JAM21nlo **J**AM21nlonll Model-4*p* 3.5 --- xFitter 3.0 0.4 0.3' (x) v bx 2.5 $q^{v}(x)$ 2.0 0.1 1.5 0.0 0.25 0.50 0.75 1.00 1.0 X 0.5 $\mu = 2 \text{ GeV}$ 0.0 ⊾ 0.0 0.2 0.4 0.6 1.0 0.8 X

pion

 $E_q(2 \text{ GeV})/m_\pi \approx 30\%$

partonic origin of proton spin

deeply virtual compton scattering: $e + P \rightarrow e + P' + \gamma$

r is Fourier conjugate of *t*

generalized parton distributions (GPD)

distribution of the longitudinal momentum fractions of partons in the transverse plane the hadron

proton GPD: unpolarized quarks inside ...

unpolarized proton

J. Miller et. al., Acta Phys. Polon. Supp. 167, 7-A6 (2023); Phys. Rev. D 106, 1, 114512 (2022)

from proton GPD to proton spin

contributions of quarks' total angular momentum to proton spin:

$$A_{2,0}(t) = \int_{-1}^{1} x H^{q}(x,\xi=0,t) dx \qquad B_{2,0}(t) = \int_{-1}^{1} dx$$

and it's distribution in the transverse plane $J^q(r_1)$

$J^{q} = \frac{1}{2} \left[A_{20}(0) + B_{20}(0) \right]$

$xE^q(x,\xi=0,t)dx$ -1

distributions of quarks' angular momenta

X. Gao et. al., Phys. Rev. D108, 1, 014507 (2023)

 $\hat{\chi}$

down

 $\mu = 2 \text{ GeV}$

2+1 dimensional image ...

quark's longitudinal momentum fraction

 r_{χ} [fm]

partonic image of proton: 3-dimensional momentum space

quarks moving transverse to proton's motion

semi-inclusive deep inelastic scattering

transverse momentum-dependent PDF (TMDPDF) b_T is Fourier conjugate of k_T

 $f(x, k_T, \mu, \eta)$

Quark Polarization

xp_

 k_{T}

0

Nucleon

Polarization

X

evolution of TMD functions across collision energies ...

Collins-Soper kernel

$$\gamma^{\overline{\mathrm{MS}}}(b_T,\mu)$$

 $= \frac{\partial \phi(x, b_T, \eta, \mu)}{\partial \ln \sqrt{\eta}}$

property of QCD vacuum — independent of hadronic state

nonperturbative Collins-Soper kernel

nonperturbative

 $\gamma^{\overline{\mathrm{MS}}}(b_{\perp},\mu)$

perturbative

nonperturbative Collins-Soper kernel from LQCD

unitary chiral quarks, physical mass

TMDPDF: polarized vs. unpolarized proton

unitary chiral quarks, physical mass

electromagnetic form factor at large momenta

$F(Q^2 \gg \Lambda_{QCD}) \sim pQCD \otimes DA$

distribution amplitude (DA)

large momenta π^+ from factor at EIC ...

2102.11788

$F_{\pi}(Q^2 \gg \Lambda_{OCD}) \sim pQCD \otimes pion DA$

NNLO pQCD Phys. Rev. Lett. 132, 20190 (2024)

unitary chiral quarks, physical mass

from QCD to EIC ...

Q. Shi et. al., <u>2404.04412</u>

beginning of a new journey...from QCD to EIC

small x & saturation physics

EIC

RHIC

1.0

Storage Ring

Detector

