

Lattice QCD calculation of hadron mass splitting

Chenfei Lu

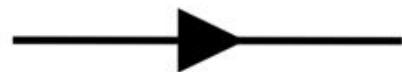
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Isospin Breaking by Perturbation

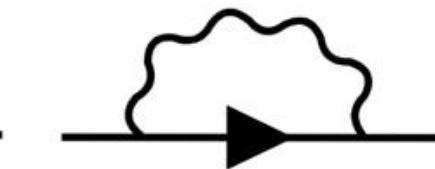
bare quark propagator:



+



+



QED

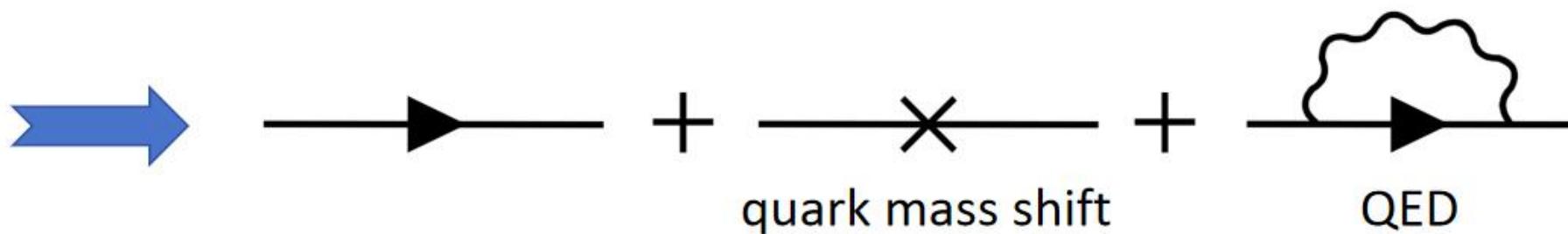
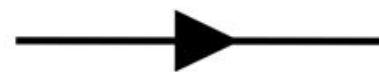


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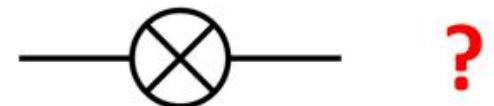
QED self-energy counter term needs renormalization.

Isospin Breaking by Perturbation

bare quark propagator:



QED self-energy counter term needs renormalization.



However, such counter term does not appear
in some cases of hadron mass splittings.

Example: π

		quark mass shift	QED			counter term
coefficients of diagrams						
m_{π^\pm}	1	$F_s(0) \cdot \Delta m_q$	$2e^2/9$	0	$-5e^2/9$	$-5e^2/9 \cdot CT$
m_{π^0}	1	$F_s(0) \cdot \Delta m_q$	$-5e^2/18$	$e^2/2$	$-5e^2/9$	$-5e^2/9 \cdot CT$
$m_{\pi^\pm} - m_{\pi^0}$	0	0	$e^2/2$	$-e^2/2$	0	0

No ultra-violet divergence nor counter term in $m_{\pi^\pm} - m_{\pi^0}$.

For baryons made of uds, the same thing happens in $m_{\Sigma^+} + m_{\Sigma^-} - 2m_{\Sigma^0}$.

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How to make it general?

Generalize to Any Hadron

- Within pure QCD

used on lattice

$$m_H = m_H^0 + \langle H | \bar{q} q | H \rangle (m_q^{phy} - m_q) + O((m_q^{phy} - m_q)^2)$$

u-d mass splitting and fine tuning.

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|
u-d mass splitting and fine tuning.
- We square it for mesons
$$m_\pi^2 = (m_\pi^0)^2 + \frac{2m_\pi^0 \langle \pi | \bar{q}q | \pi \rangle (m_q^{phy} - m_q)}{|} + O\left((m_q^{phy} - m_q)^2\right)$$

|
relates mildly to pion mass.

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$$m_\pi^2 = (m_\pi^0)^2 + \frac{2m_\pi^0 \langle \pi | \bar{q}q | \pi \rangle (m_q^{phy} - m_q)}{\text{lattice calculation}} + O((m_q^{phy} - m_q)^2)$$

relates mildly to pion mass.

prediction

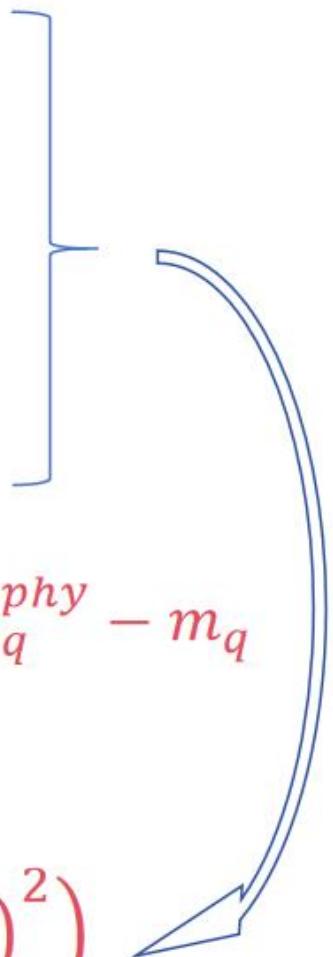
$$m_H = m_H^0 + \frac{m_\Omega^{phy} \langle H | \bar{q}q | H \rangle}{2m_\pi^0 \langle \pi | \bar{q}q | \pi \rangle} \frac{m_\pi^2 - (m_\pi^0)^2}{m_\Omega^{phy}} + O((m_q^{phy} - m_q)^2)$$

Z_s cancels

input

cancel $m_q^{phy} - m_q$

used on lattice



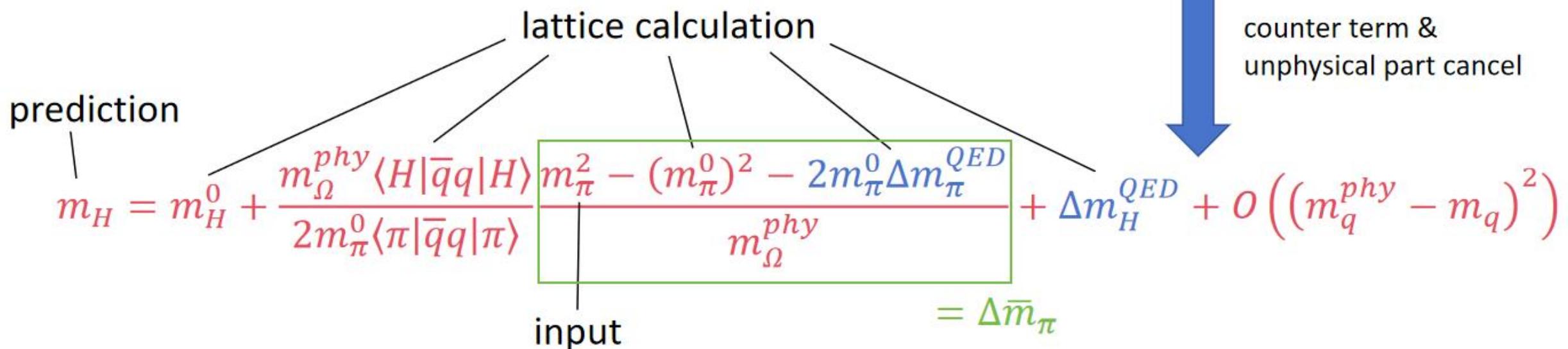
Introduce QED Effect

UV divergence is cut off by lattice spacing

$$m_H = m_H^0 + \langle H | \bar{q} q | H \rangle (m_q^{phy} - m_q + CT) + \Delta m_H^{QED} + O((m_q^{phy} - m_q)^2)$$

QED self-energy counter term of quark

$$m_\pi^2 = (m_\pi^0)^2 + 2m_\pi^0 \langle \pi | \bar{q} q | \pi \rangle (m_q^{phy} - m_q + CT) + 2m_\pi^0 \Delta m_\pi^{QED} + O((m_q^{phy} - m_q)^2)$$



Lattice Calculation

- For lattice with 3 flavors of quarks, we need 3 kinds of mesons

$$\begin{aligned} m_H &= m_H^0 + \sum_q \langle H | \bar{q}q | H \rangle (m_q^{phy} - m_q + CT) + \Delta m_H^{QED} + O\left((m_q^{phy} - m_q)^2\right) \\ &= m_H^0 + \alpha \Delta \bar{m}_{\pi^\pm} + \beta \Delta \bar{m}_{K^\pm} + \gamma \Delta \bar{m}_{K^0} + \Delta m_H^{QED} + O\left((m_q^{phy} - m_q)^2\right) \end{aligned}$$

α, β, γ : combinations of ratios of $F_s(0)$  need to compute properties of π & K .

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α, β, γ : combinations of ratios of $F_s(0)$  need to compute properties of π & K .

- Both input and output need lattice spacing to convert between lattice and physical unit

$$a^{QCD} m_H^{phy} = m_H^0$$

The lattice spacing changes while introducing QED. We need another hadron (eg. Ω)

$$a^{full} m_\Omega^{phy} = m_\Omega = m_\Omega^0 + \alpha \Delta \bar{m}_{\pi^\pm} + \beta \Delta \bar{m}_{K^\pm} + \gamma \Delta \bar{m}_{K^0} + \dots$$

as input to solve the new spacing a^{full} .

include a^{full}

Numerical Analysis

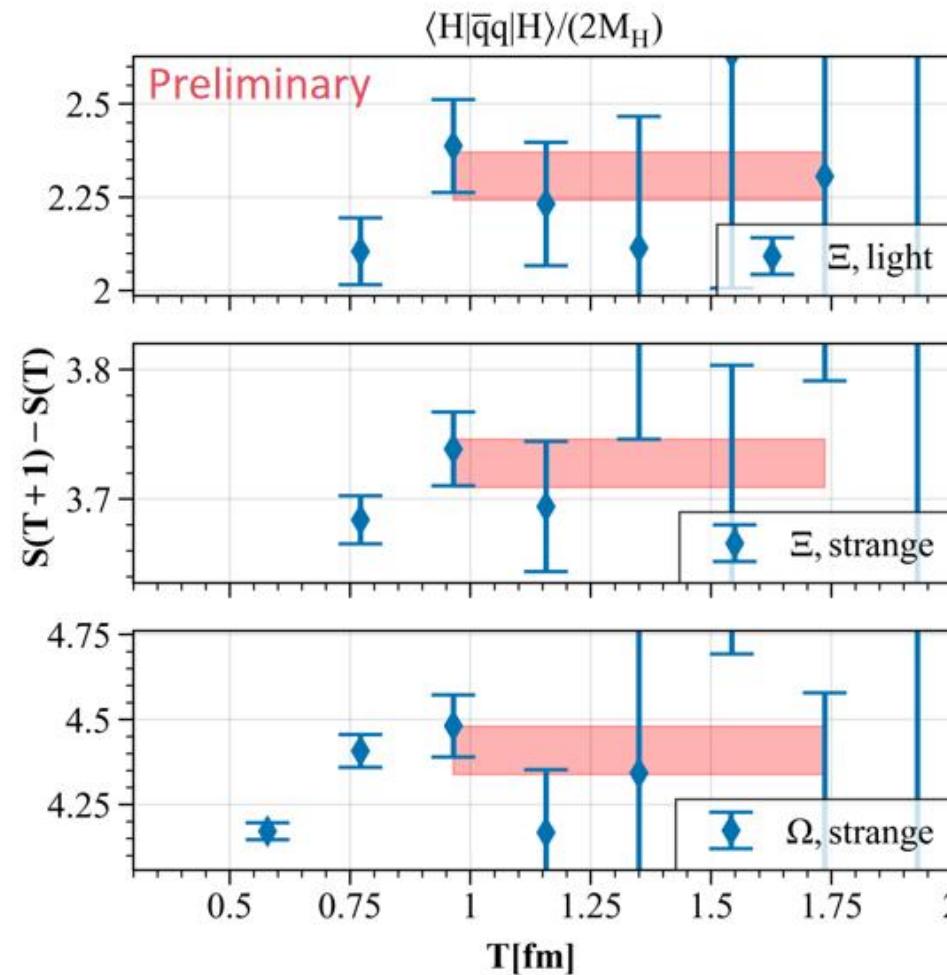
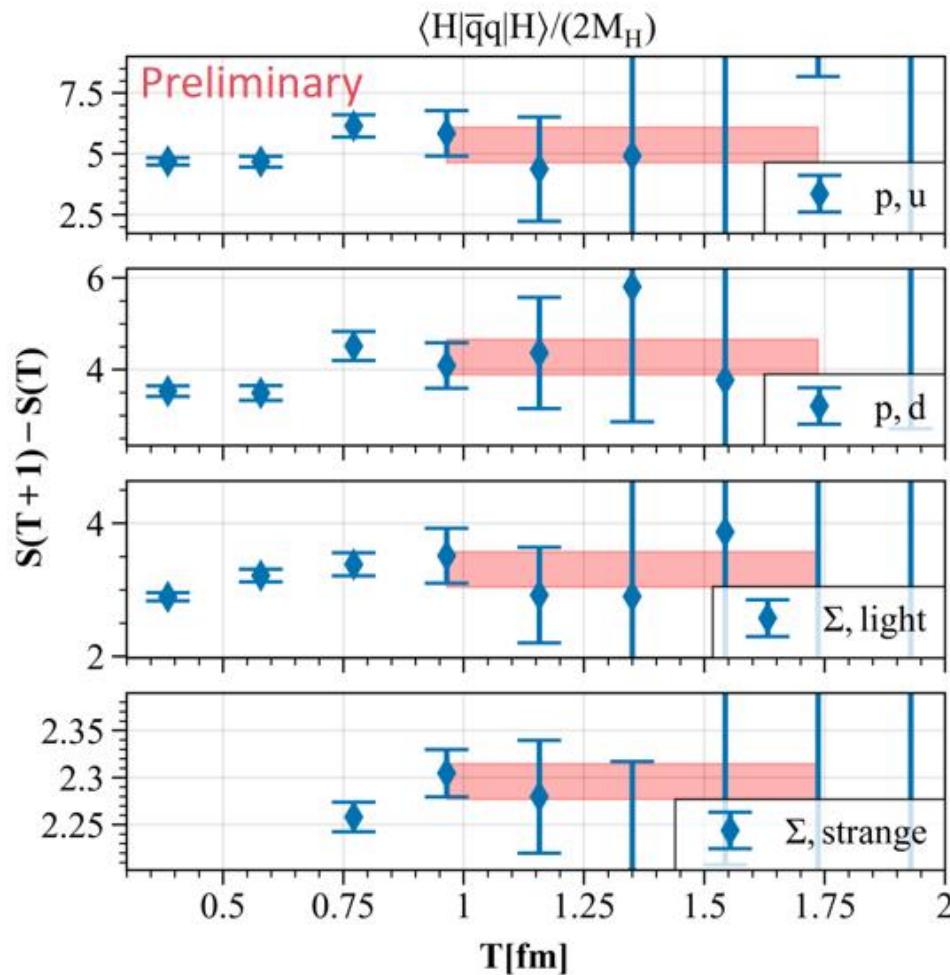
We compute $\left\{ \begin{array}{l} m^0, \quad \text{by 2pt} \\ Z_s F_s(0), \quad \text{by 3pt of several particles} \\ \Delta m^{QED}, \quad \text{by 4pt} \end{array} \right.$ $\left. \begin{array}{l} \pi \\ K \\ \Omega \\ N \\ \Sigma \\ \Xi \\ \vdots \end{array} \right\}$ lattice spacing prediction

Name	N	L	T	a	m_π
24D	94	24	64	0.19fm	143MeV

Generated by RBC/UKQCD, 2+1 flavor

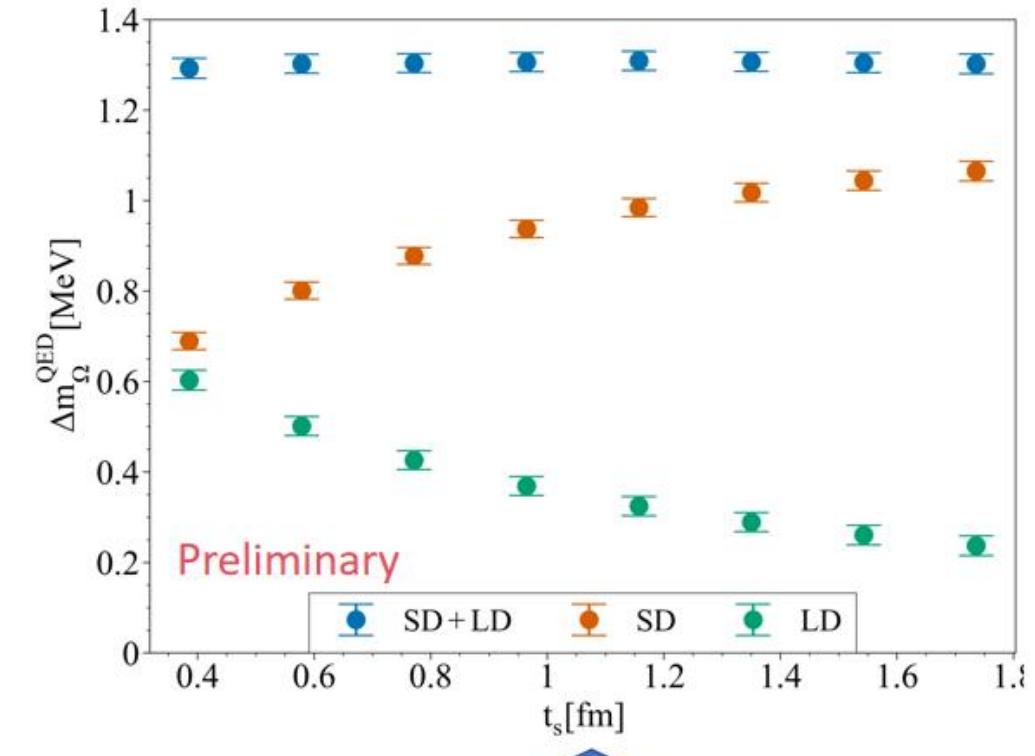
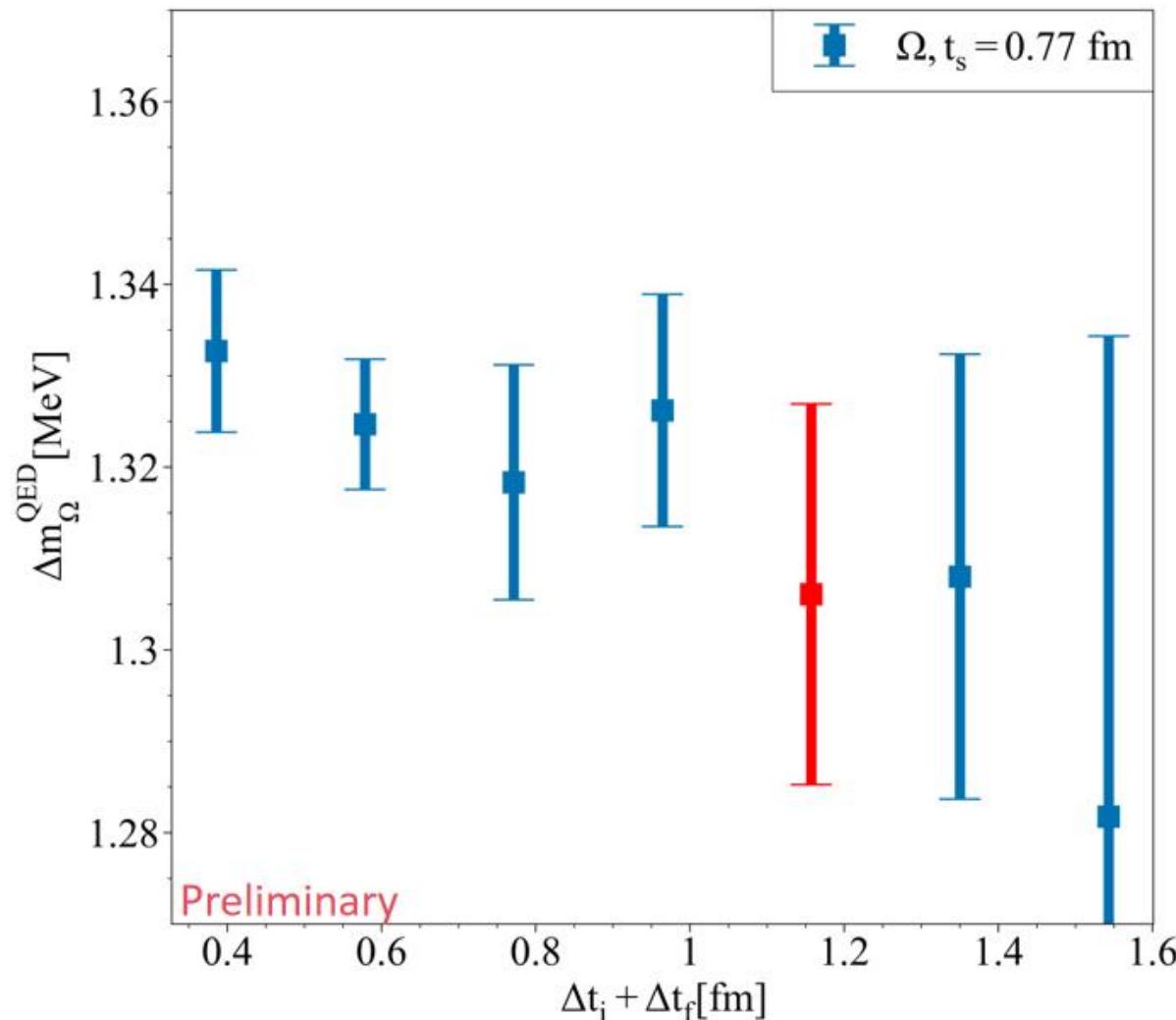
$$\begin{cases} m^0, & \text{by 2pt} \\ Z_s F_s(0), & \text{by 3pt} \\ \Delta m^{QED}, & \text{by 4pt} \end{cases}$$

$$S(T) = \sum_{\tau=0}^T \langle O_H(T) J_s(\tau) O_H^+(0) \rangle = C_0 + \langle H | \bar{q} q | H \rangle \cdot T + (\text{excited states})$$



m^0 , by 2pt
 $Z_s F_s(0)$, by 3pt
 Δm^{QED} , by 4pt

$$\Delta m^{QED} = \frac{e^2}{2} \int d^4x \langle H | J^\mu(x) J^\nu(0) | H \rangle S_{\mu\nu}^\gamma(x)$$



Exponentially suppressed finite volume effects

Preliminary Results

$$a^{full} m_H^{phy} = m_H = m_H^0 + \alpha \Delta \bar{m}_{\pi^\pm} + \beta \Delta \bar{m}_{K^\pm} + \gamma \Delta \bar{m}_{K^\pm} + \Delta m_H^{QED} + O\left((m_q^{phy} - m_q)^2\right)$$

➤ For Ω

$$\alpha = -1.198(23),$$

$$\beta = \gamma = 1.264(20)$$

$$\frac{a^{full \ theory}}{a^{pure \ QCD}} = 1.0195(50)$$

➤ To do next...

- $N, \Sigma, \Xi, \Delta \dots$, for example, $\alpha_N, \beta_N, \gamma_N, \Delta m_N^{QED}, a^{full} \rightarrow \Delta m_N$
- Disconnected diagrams
- Add finer lattice to do continuum extrapolation

Conclusion

- We develop a new method to compute isospin breaking effect on hadron mass splitting on lattice.
- Using this method we can compute hadron mass splitting without ultraviolet divergence nor renormalization.
- We apply it on a real lattice to compute several hadrons.

Thank you for your listening.