

Non-singlet axial current improvement for massless and massive sea quarks

Justus Kuhlmann

Patrick Fritzsch, Jochen Heitger



Relevance for further improvement and physics

- ▶ exp. Wilson-clover fermion framework
- ▶ massive $\hat{=}$ at $N_f = 3$ symmetric point

Relevance for further improvement and physics

- ▶ exp. Wilson-clover fermion framework
- ▶ massive $\hat{=}$ at $N_f = 3$ symmetric point

- ▶ needed for improv. quark current mass
- ▶ decay constants & matrix elements

Relevance for further improvement and physics

- ▶ exp. Wilson-clover fermion framework
- ▶ massive $\hat{=}$ at $N_f = 3$ symmetric point

- ▶ needed for improv. quark current mass
- ▶ decay constants & matrix elements

- ▶ improvement and renormalisation:
 - ▶ c_V, c_T, Z_A
 - ▶ no $c_A \Rightarrow$ no improvement of other channels

Determination of c_A

- ▶ Schrödinger functional boundary conditions
- ▶ similar to quenched [[hep-lat/9609035](#)], $N_f = 2$ [[hep-lat/0503003](#)] and std. Wilson-Clover $N_f = 3$ [[1502.04999](#), [hep-lat/0703006](#)]
- ▶ derive from PCAC mass

$$m_{\text{PCAC}} = \frac{\partial_0 f_A}{2f_P} + c_A a \frac{\partial_0^2 f_P}{2f_P} = r + c_A as$$

$$m_{\text{PCAC}}^{(0)} = m_{\text{PCAC}}^{(1)} \quad \Leftrightarrow \quad c_A = -\frac{r^{(1)} - r^{(0)}}{s^{(1)} - s^{(0)}}$$

- ▶ states (0) and (1) are the PS ground and first excited state in our setup
- ▶ PCAC relation holds for both

The wavefunction method

- ▶ construct pseudoscalar states
 - ▶ H-like basis wavefunctions: $\omega_1 = e^{-r/a_0}$, $\omega_2 = r e^{-r/a_0}$, $\omega_3 = e^{-r/(2a_0)}$
 - ▶ also include $\omega_4 = \text{cons.}$, $\omega_5 = -r^2 e^{-r/a_0}$ with $r = |\vec{y} - \vec{x}|$

The wavefunction method

- ▶ construct pseudoscalar states
 - ▶ H-like basis wavefunctions: $\omega_1 = e^{-r/a_0}$, $\omega_2 = r e^{-r/a_0}$, $\omega_3 = e^{-r/(2a_0)}$
 - ▶ also include $\omega_4 = \text{cons.}$, $\omega_5 = -r^2 e^{-r/a_0}$ with $r = |\vec{y} - \vec{x}|$
- ▶ diagonalise boundary-to-boundary corr. func. $(F_1)_{i,j} = -\langle O(\omega_i)O'(\omega_j) \rangle$

The wavefunction method

- ▶ construct pseudoscalar states
 - ▶ H-like basis wavefunctions: $\omega_1 = e^{-r/a_0}$, $\omega_2 = r e^{-r/a_0}$, $\omega_3 = e^{-r/(2a_0)}$
 - ▶ also include $\omega_4 = \text{cons.}$, $\omega_5 = -r^2 e^{-r/a_0}$ with $r = |\vec{y} - \vec{x}|$
- ▶ diagonalise boundary-to-boundary corr. func. $(F_1)_{i,j} = -\langle O(\omega_i)O'(\omega_j) \rangle$
- ▶ employ eigenvectors of $(F_1)_{i,j}$ to project $f_A(x_0)$ and $f_P(x_0)$ onto the eigenstates

The wavefunction method

- ▶ construct pseudoscalar states
 - ▶ H-like basis wavefunctions: $\omega_1 = e^{-r/a_0}$, $\omega_2 = r e^{-r/a_0}$, $\omega_3 = e^{-r/(2a_0)}$
 - ▶ also include $\omega_4 = \text{cons.}$, $\omega_5 = -r^2 e^{-r/a_0}$ with $r = |\vec{y} - \vec{x}|$
- ▶ diagonalise boundary-to-boundary corr. func. $(F_1)_{i,j} = -\langle O(\omega_i)O'(\omega_j) \rangle$
- ▶ employ eigenvectors of $(F_1)_{i,j}$ to project $f_A(x_0)$ and $f_P(x_0)$ onto the eigenstates
- ▶ evaluate $c_A(x_0)$ with projected correlation functions
- ▶ later: choice of x_0 and wavefunction basis is part of the improvement condition

Ensembles

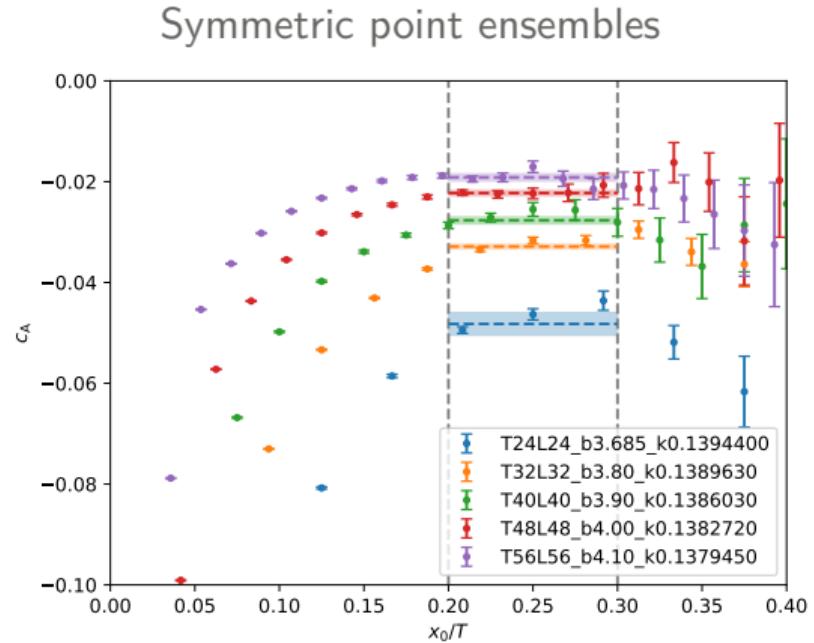
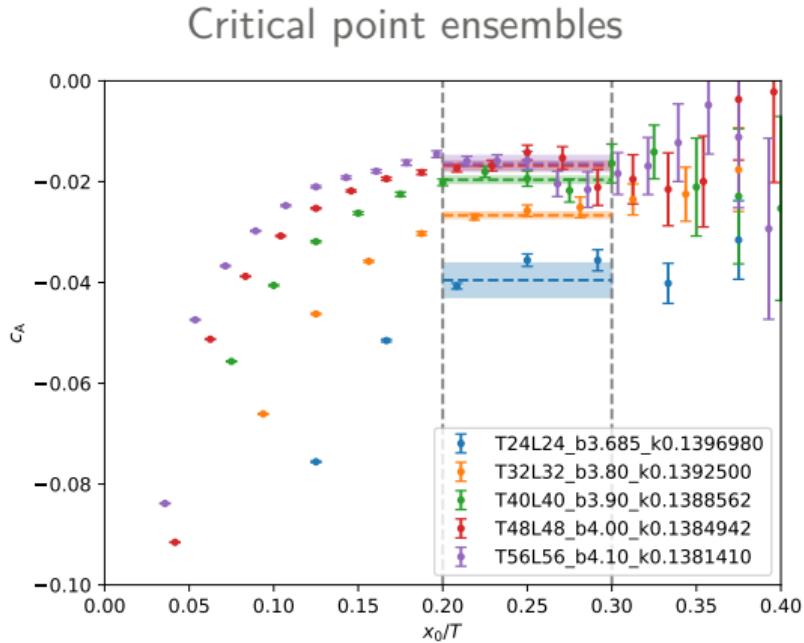
$T = L \approx 3 \text{ fm}$ Schrödinger-Functional ensembles, exp. Wilson-Clover fermions

| L/a | β | $\kappa_1 \approx \kappa_{\text{cr}}$ | κ_2 | $\kappa_3 \approx \kappa_{\text{sym}}$ | $\approx a$ [fm] |
|-------|---------|---------------------------------------|------------|--|------------------|
| 24 | 3.685 | 0.1396980 | 0.1395500 | 0.1394400 | 0.120 |
| 32 | 3.80 | 0.1392500 | — | 0.1389630 | 0.095 |
| 40 | 3.90 | 0.1388562 | 0.1386148 | 0.1386030 | 0.080 |
| 48 | 4.00 | 0.1384942 | 0.1384880 | 0.1382720 | 0.064 |
| 56 | 4.10 | 0.1381410 | 0.1380000 | 0.1379450 | 0.055 |
| 96 | 4.37 | — | — | — | 0.035 |

- interested in 2 LCPs: chiral and $N_f = 3$ sym. point
- matching sym. point of OpenLat [\[2201.03874\]](#)

Improvement of the axial-vector current

c_A estimators



Interpolation

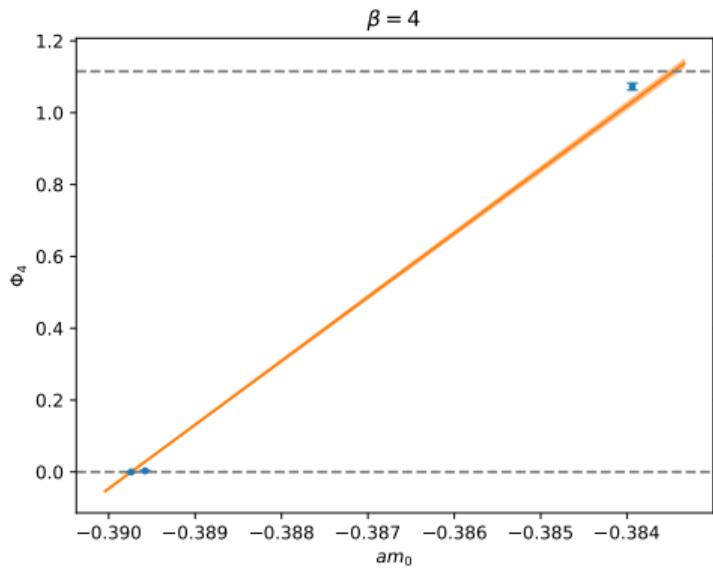
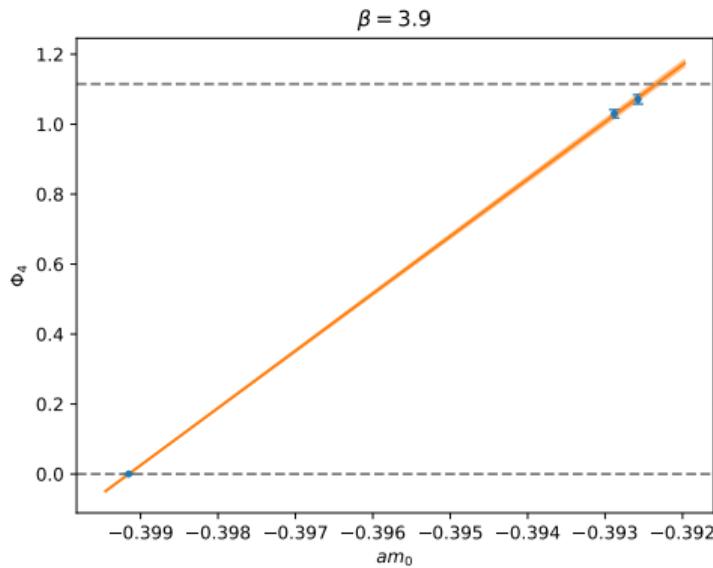
... to the symmetric and critical point

- ▶ ensembles not exactly tuned
- ▶ able to interpolate to the desired points due to 2 or 3 ensembles per β
- ▶ determine points of interest as in OpenLat ensembles [2201.03874]
- ▶ define:

$$\Phi_4^{\text{SF}} = \frac{3}{2} 8t_0 |m_{\text{eff}}| m_{\text{eff}} \quad \Rightarrow \quad \Phi_4^{\text{SF}} \Big|_{m_{0,\text{cr}}} = 0, \quad \Phi_4^{\text{SF}} \Big|_{m_{0,\text{sym}}} = 1.115$$

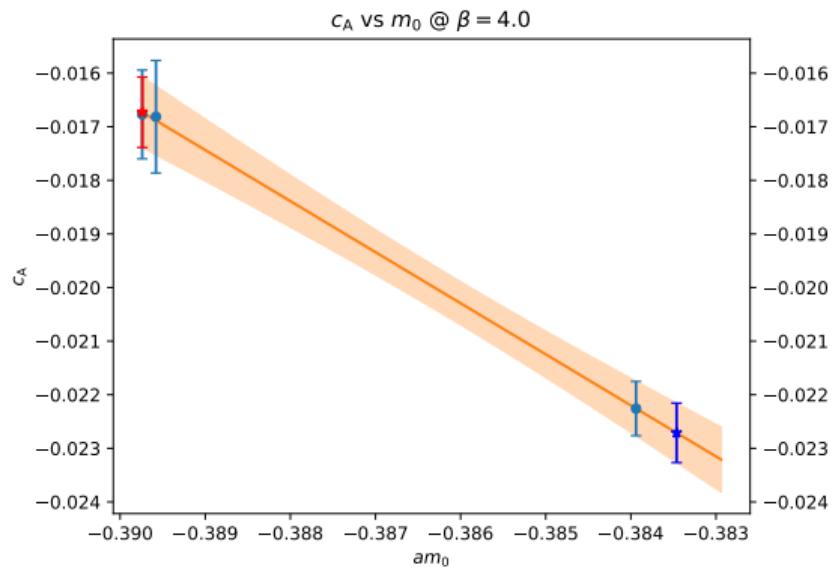
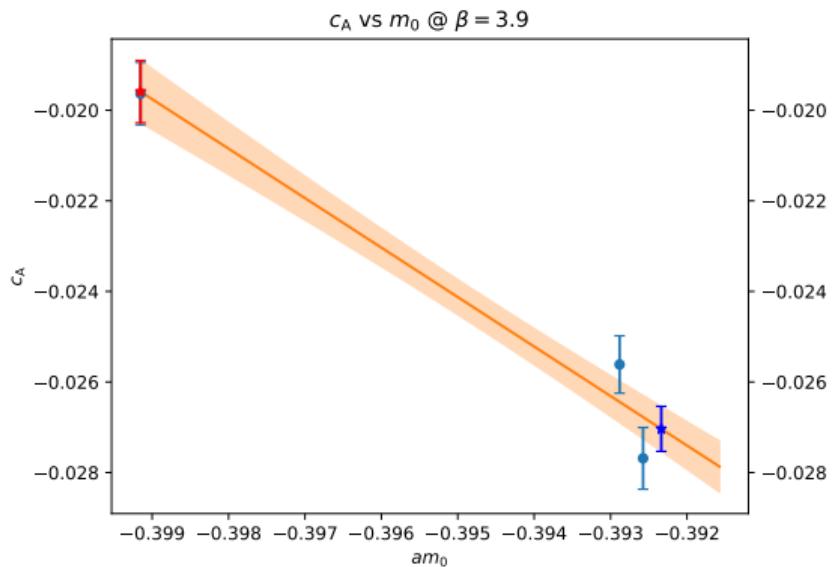
Improvement of the axial-vector current

Finding the symmetric and chiral point



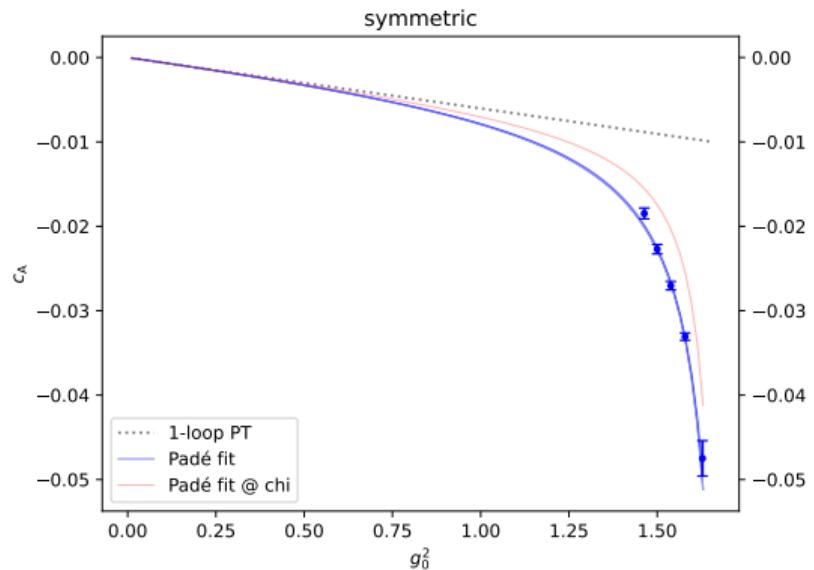
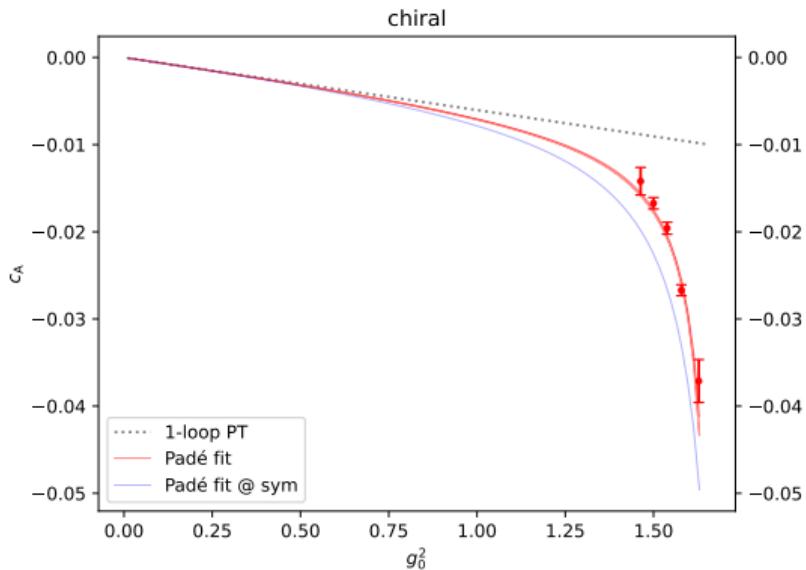
Improvement of the axial-vector current

Interpolations in c_A



Improvement of the axial-vector current

Interpolations in g_0^2



First scaling test of improvement

- ▶ Example: Calculate $f_{\pi K}$ with stabilised Wilson fermions
- ▶ symmetric point OpenLat ensembles
- ▶ improve with $c_A = 0$ vs $c_A(g_0^2)|_{\text{chi}}$ vs $c_A(g_0^2)|_{\text{sym}}$

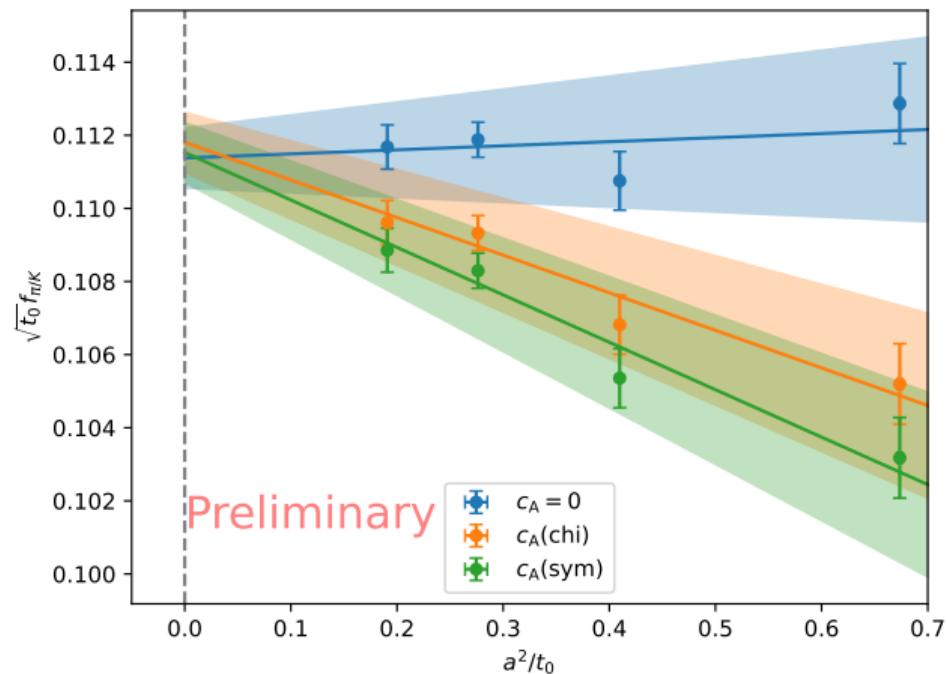
$$f_{\pi K}^{\text{RI}} = Z_A(1 + b_A a m_q + \bar{b}_A a \text{Tr}[M_q]) \frac{\sqrt{2} \mathcal{A}_{A_0 P}}{\sqrt{\mathcal{A}_{PP}} m_\pi}$$

with $C_{XX} \propto \mathcal{A}_{XX} e^{-mx_0}$

- ▶ renormalisation: Z_A preliminary, b_A from pert. theory, \bar{b}_A neglected

First scaling test of improvement

Results

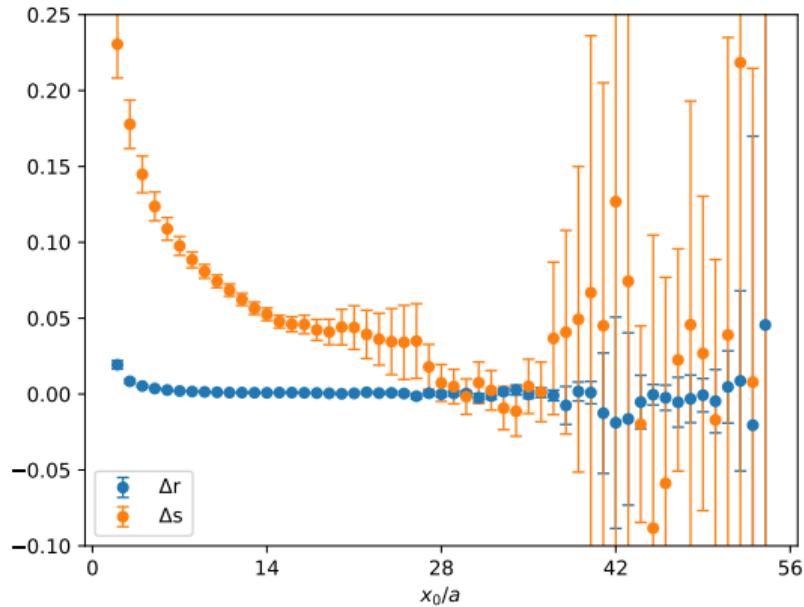
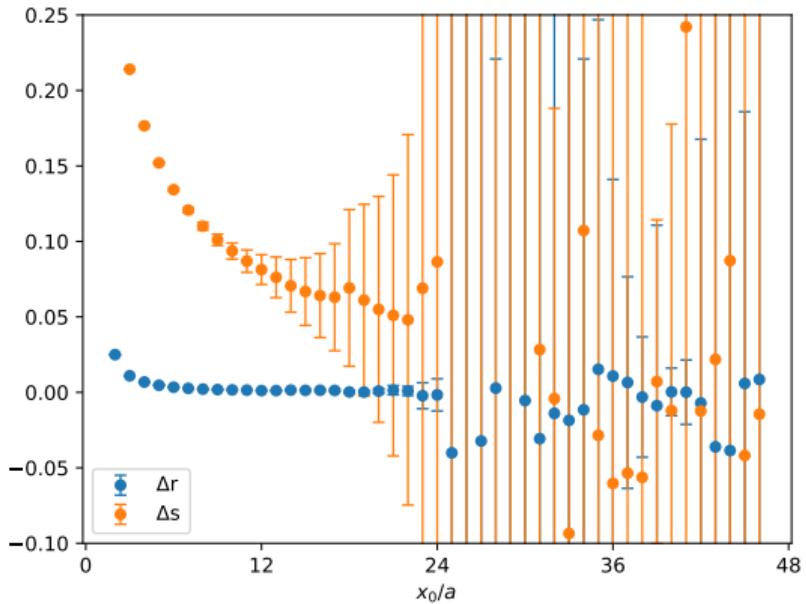


Outlook

- ▶ finish Z_A , Z_V , b_V and \bar{b}_V through SF
- ▶ further improvement and renormalisation currently in the works:
 - ▶ vector and tensor current improvement (c_V , c_T)
 - ▶ current quark mass renormalisation ($b_A - b_P$, b_m , Z)
 - ▶ determination of Z_A , Z_V , Z_S/Z_P through χ SF

Backup

Δr and Δs behaviour



Backup

Behaviour of f_P projected

