

# Dark Matter: from simulations to detection

Nassim Bozorgnia

Institute for Particle Physics Phenomenology  
Durham University

# Dark Matter halo

*What is the distribution of Dark Matter (DM) in halo of our Galaxy?*



# Dark Matter halo

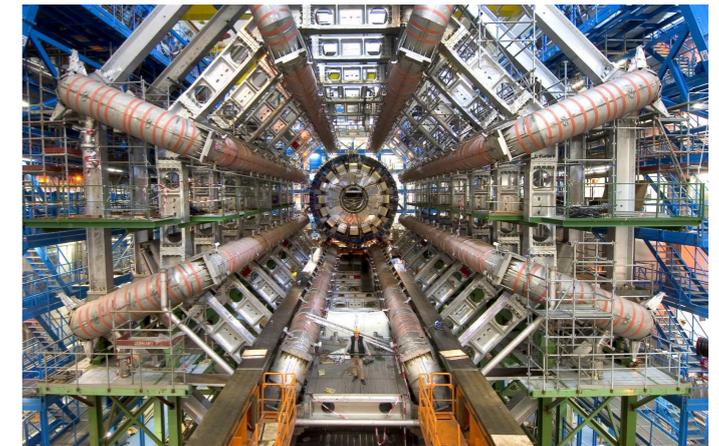
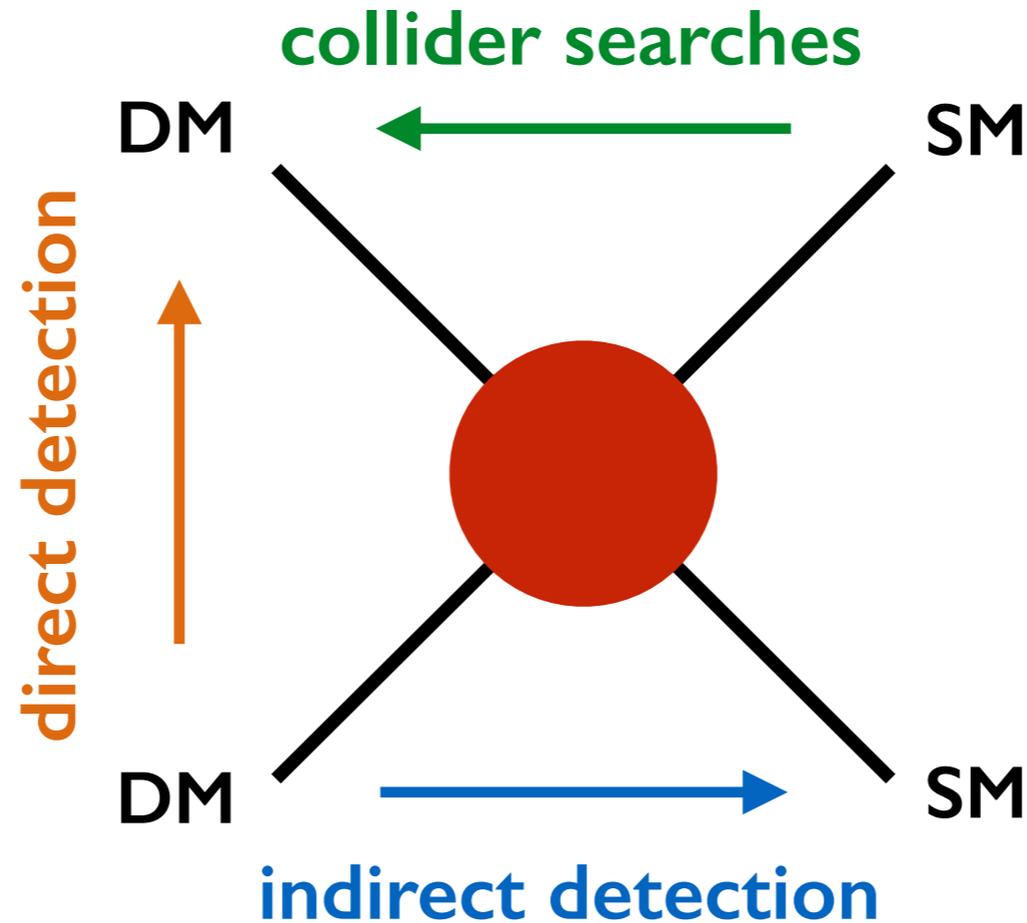
*What is the distribution of Dark Matter (DM) in halo of our Galaxy?*

Uncertainties in the DM distribution → *prevents a precise determination of the properties of the DM particle.*



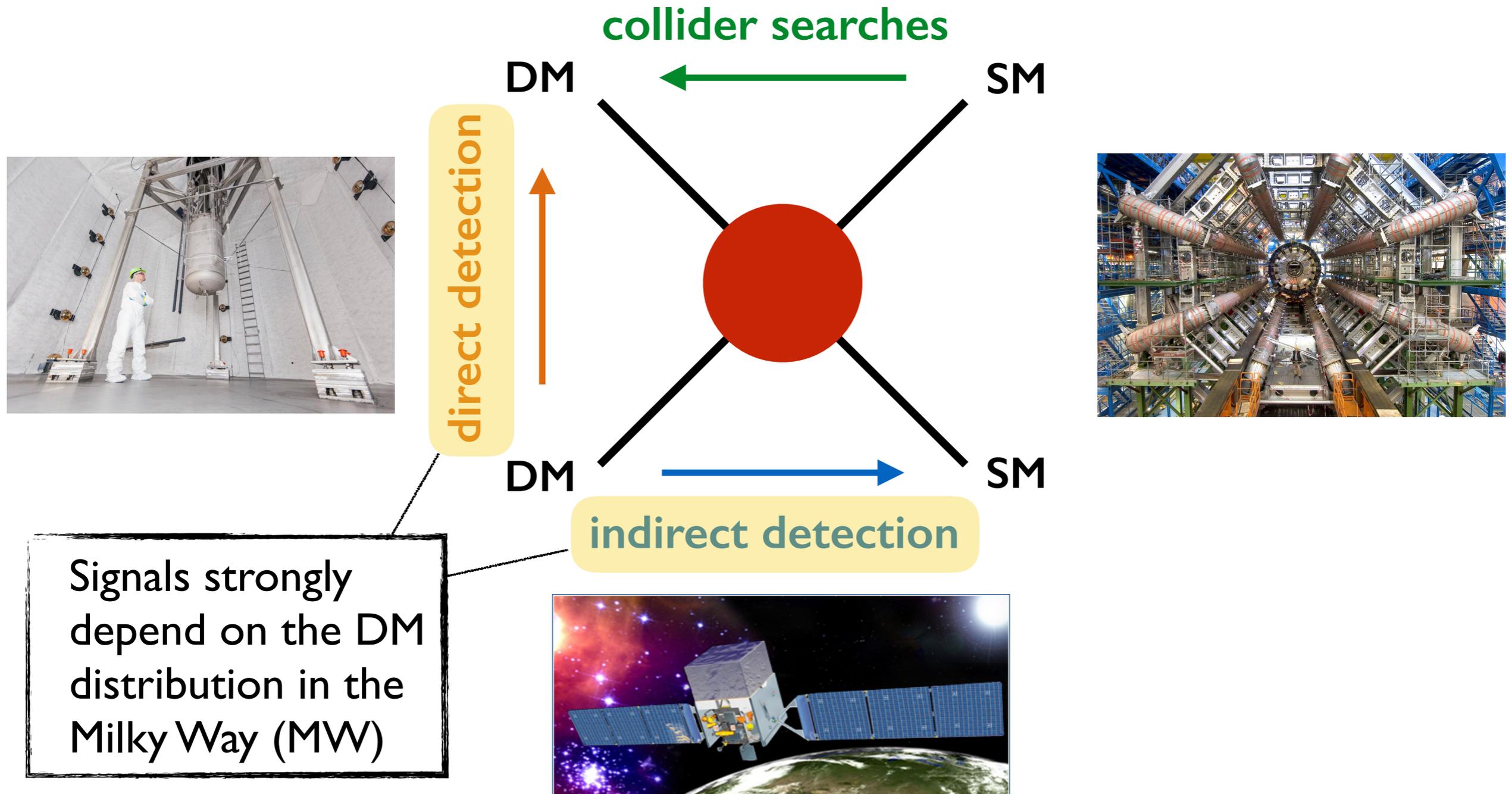
# Dark Matter searches

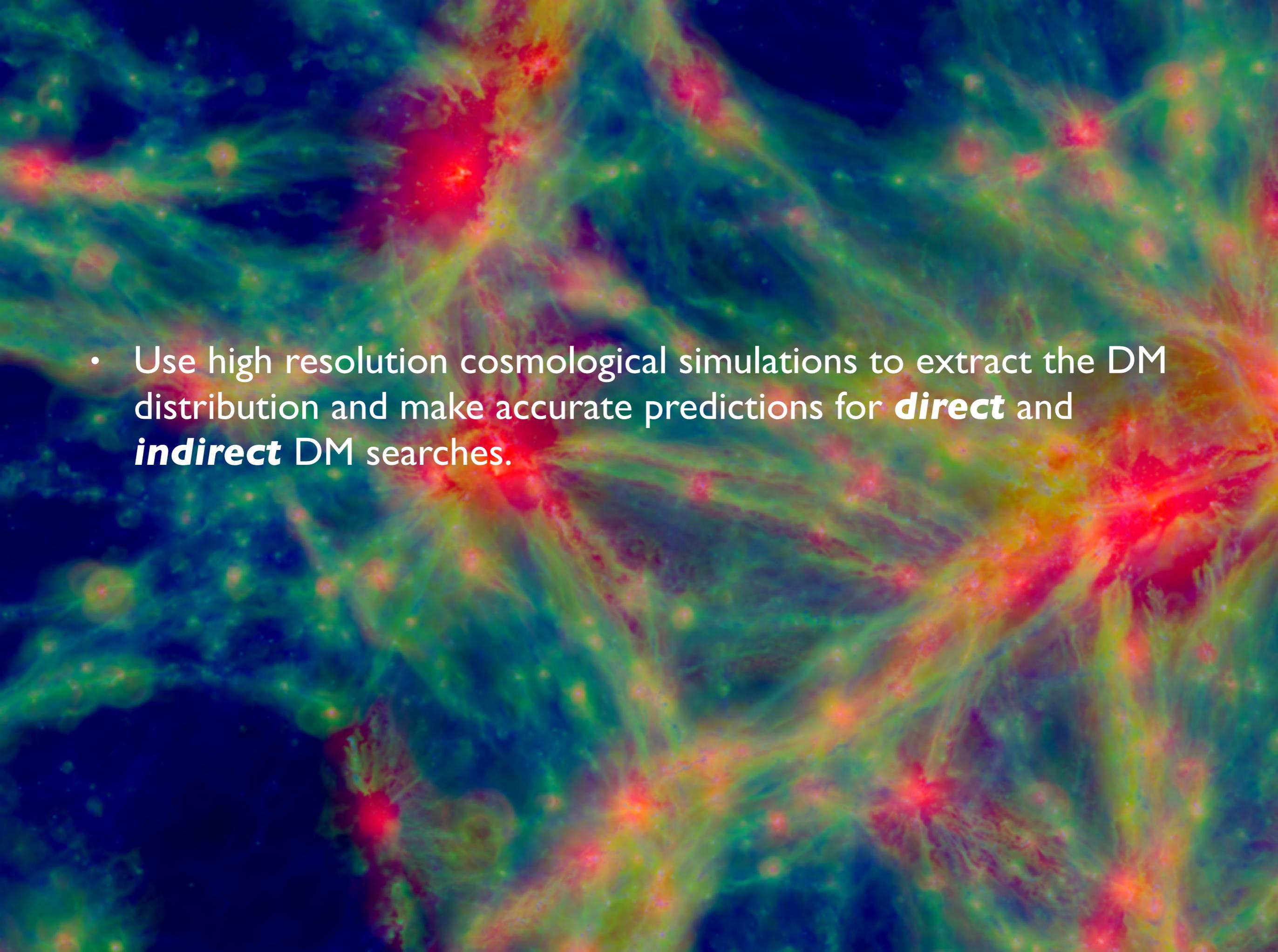
- Searching for **WIMPs** in three complementary ways:



# Dark Matter searches

- Searching for **WIMPs** in three complementary ways:



- 
- Use high resolution cosmological simulations to extract the DM distribution and make accurate predictions for **direct** and **indirect** DM searches.

# Prospects for direct DM searches

# Direct DM detection event rate

- The differential event rate:

$$\frac{dR}{dE_R} = \frac{\rho_\chi}{m_\chi m_N} \int_{v > v_{\min}} d^3v \frac{d\sigma_{\chi N}}{dE_R} v f_{\text{det}}(\mathbf{v}, t)$$

where  $v_{\min} = \sqrt{m_N E_R / (2\mu_{\chi N}^2)}$  is the minimum DM speed required to produce a recoil energy  $E_R$ .

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astrophysics

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- Astrophysical inputs:**
  - local DM density:** *normalization in event rate.*
  - local DM velocity distribution:** *enters the event rate through an integration.*

# Direct DM detection event rate

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astrophysics

- For standard spin-independent and spin-dependent interactions:

$$\frac{dR}{dE_R} = \frac{\sigma_0 F^2(E_R)}{2m_\chi \mu_{\chi N}^2} \rho_\chi \eta(v_{\min}, t)$$

particle physics      astrophysics

where

$$\eta(v_{\min}, t) \equiv \int_{v > v_{\min}} d^3v \frac{f_{\text{det}}(\mathbf{v}, \mathbf{t})}{v}$$

halo integral

# Local Dark Matter distribution

*What is the distribution of DM in the Sun's neighborhood?*



# Standard Halo Model

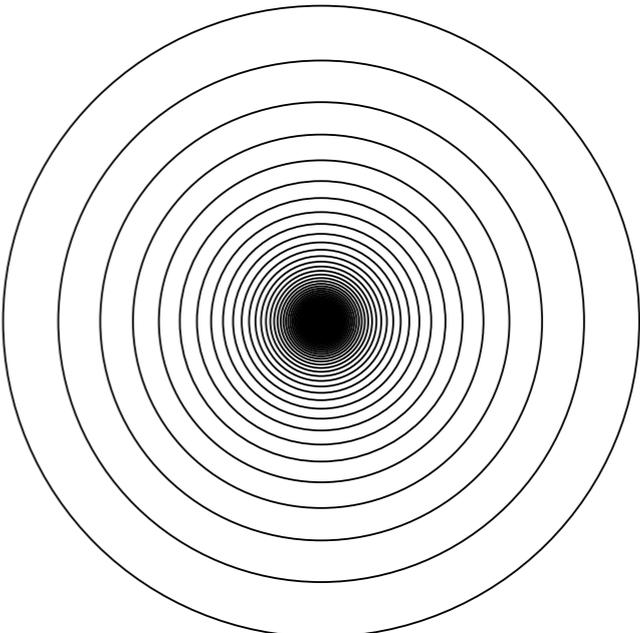
- The simplest model for the DM distribution in our Galaxy is the *Standard Halo model (SHM)*: isothermal sphere with an isotropic Maxwell-Boltzmann velocity distribution.

Drukier, Freese, Spergel, 1986

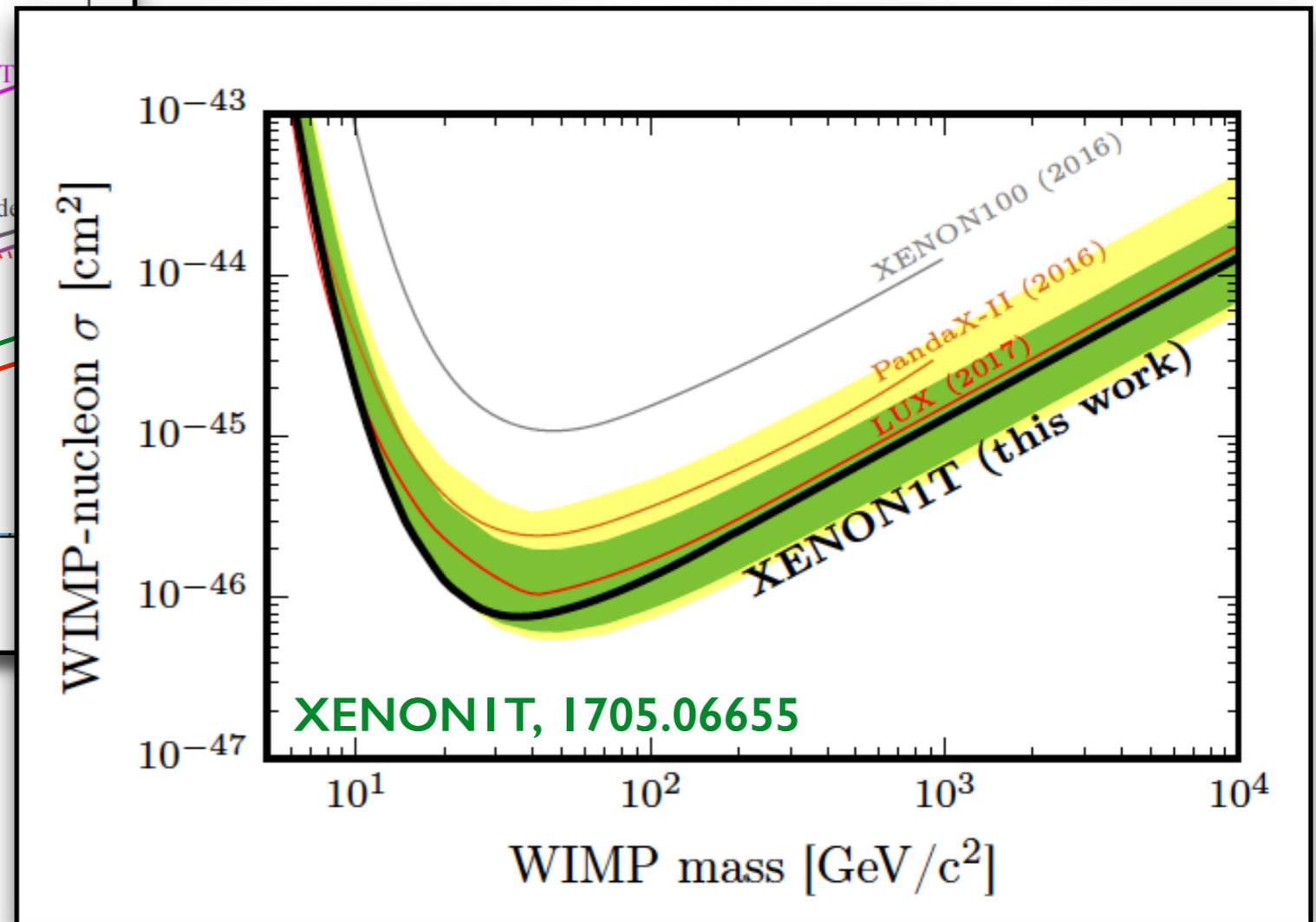
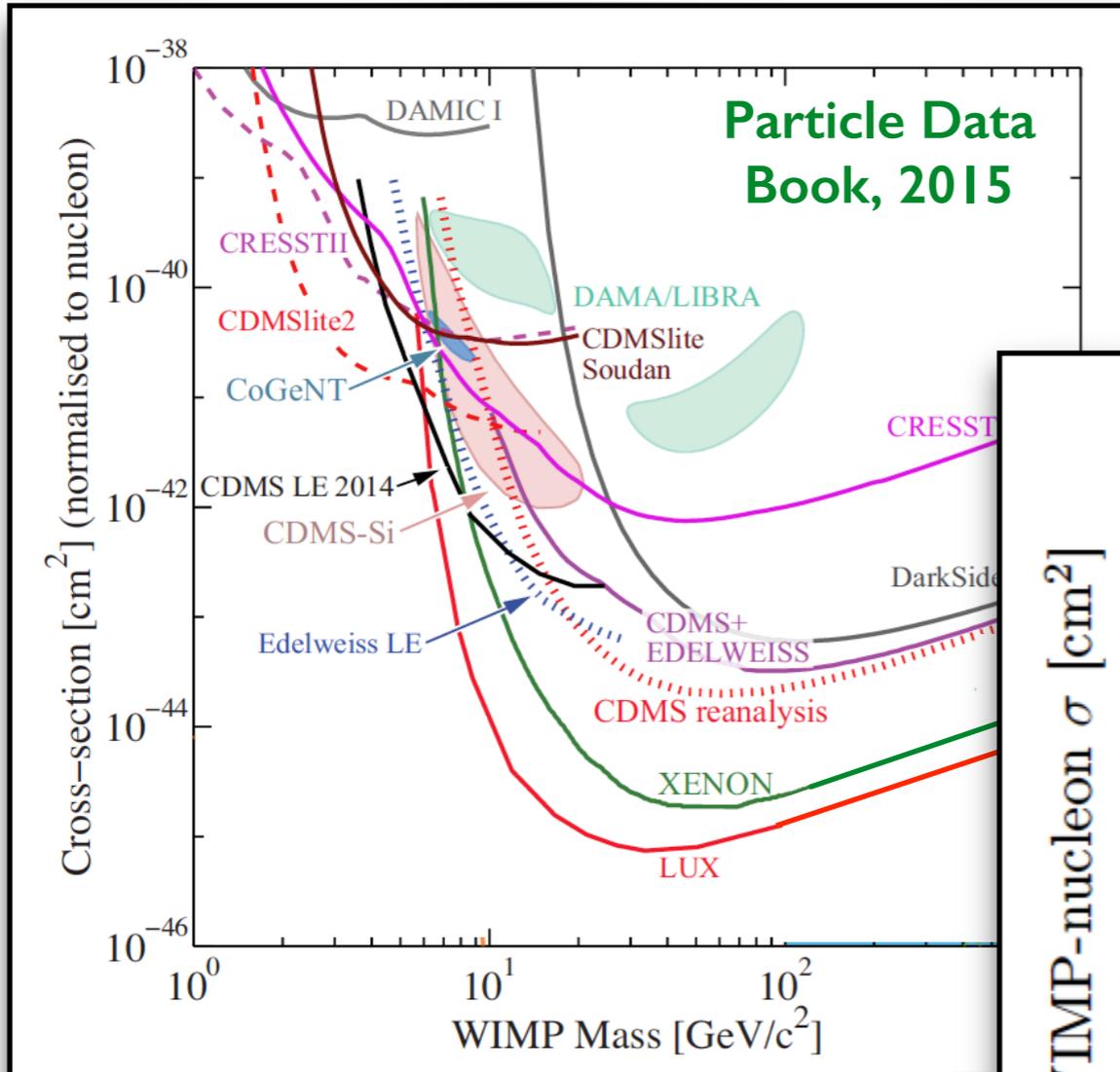
# Standard Halo Model

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Drukier, Freese, Spergel, 1986

- Hydrostatic equilibrium: pressure balances gravitational potential
  - Density profile:  $\rho(r) \propto r^{-2}$
  - Local DM density:  $0.3 \text{ GeV/cm}^3$
  - Typical DM speed: 220 km/s
- 
- Actual DM distribution may *deviate substantially* from the SHM.

# Direct detection results

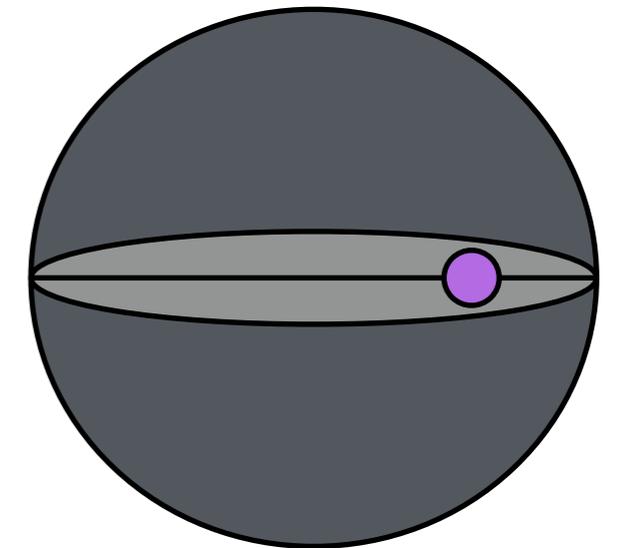


- Assumption in these kinds of plots: **SHM**

# Local Dark Matter density

## From observations:

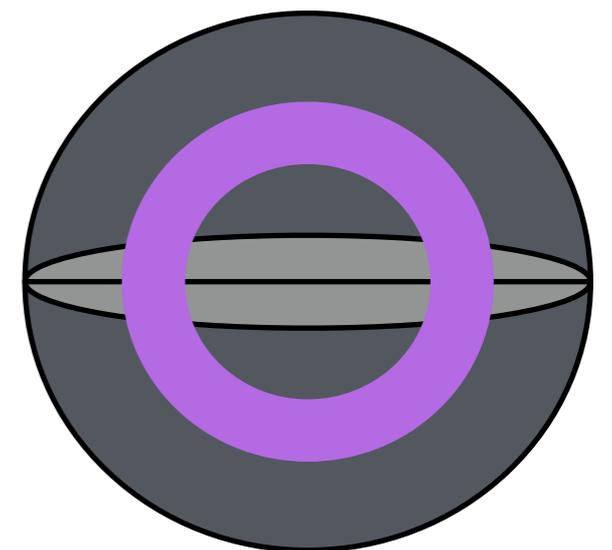
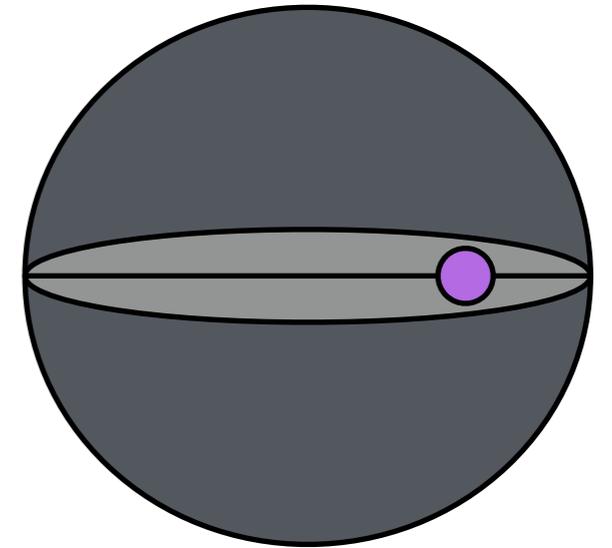
- **Local estimates:** use kinematical data from a nearby population of stars.
- Robust measurements, but need to account for the local contribution of baryons which has significant uncertainties. → *large error bars*



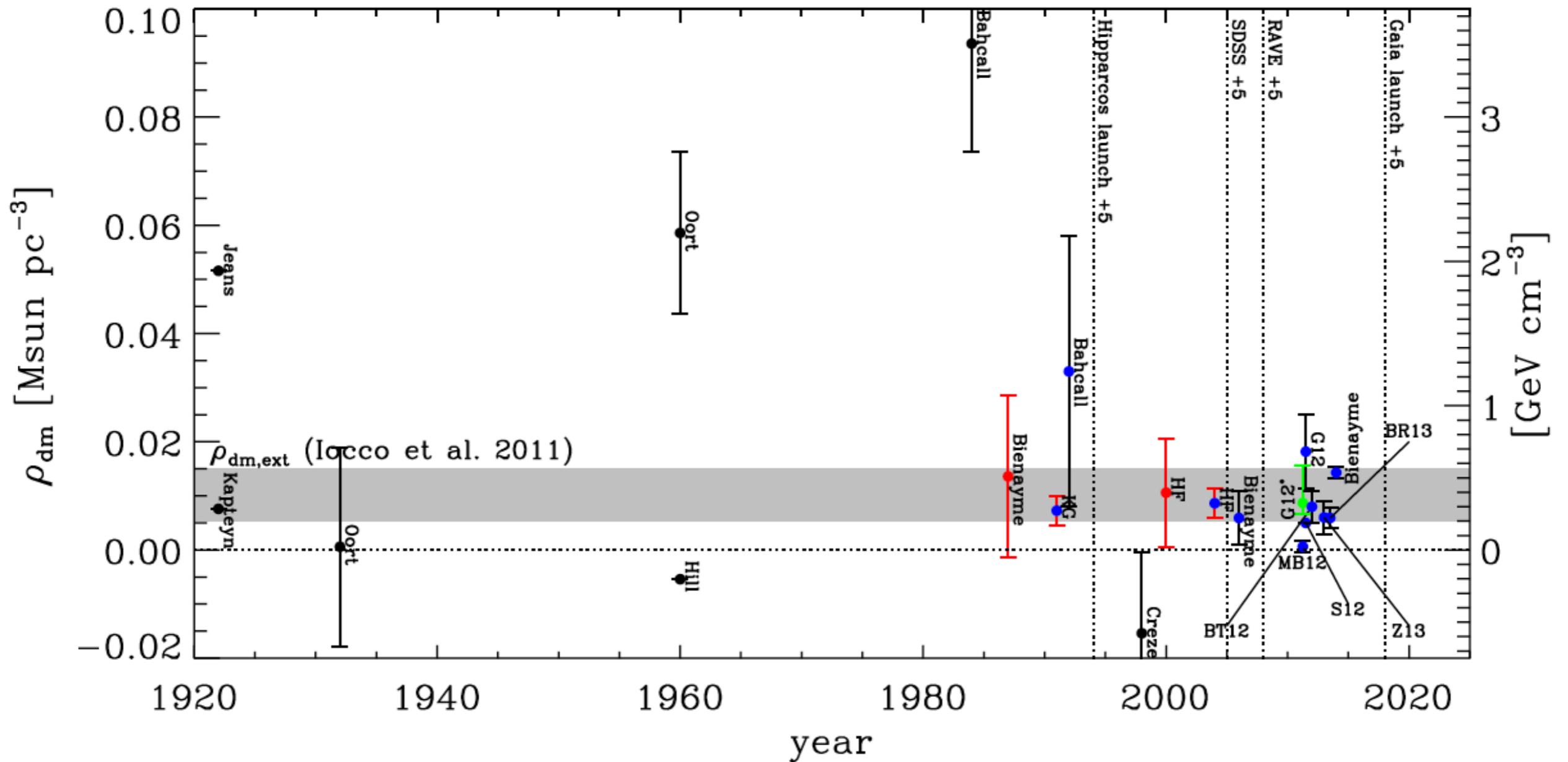
# Local Dark Matter density

## From observations:

- **Local estimates:** use kinematical data from a nearby population of stars.
  - Robust measurements, but need to account for the local contribution of baryons which has significant uncertainties. → *large error bars*
- **Global estimates:** based on mass modeling of the MW, and fits to kinematical data across the Galaxy.
  - Good precision ( $\sim 10\%$ ), but estimates are strongly model dependent. → *systematic uncertainties*



# Local Dark Matter density



Read, 1404.1938

# Local DM velocity distribution

- The velocity distribution depends on the halo model.
- In the **SHM**, a truncated Maxwellian velocity distribution is assumed:

$$f_{\text{gal}}(\mathbf{v}) = \begin{cases} N \exp(-\mathbf{v}^2/v_c^2) & v < v_{\text{esc}} \\ 0 & v \geq v_{\text{esc}} \end{cases}$$

with  $v_c = 220$  km/s and  $v_{\text{esc}} = 550$  km/s.

$\sigma_v = \sqrt{3/2} v_c$  independent of radius.

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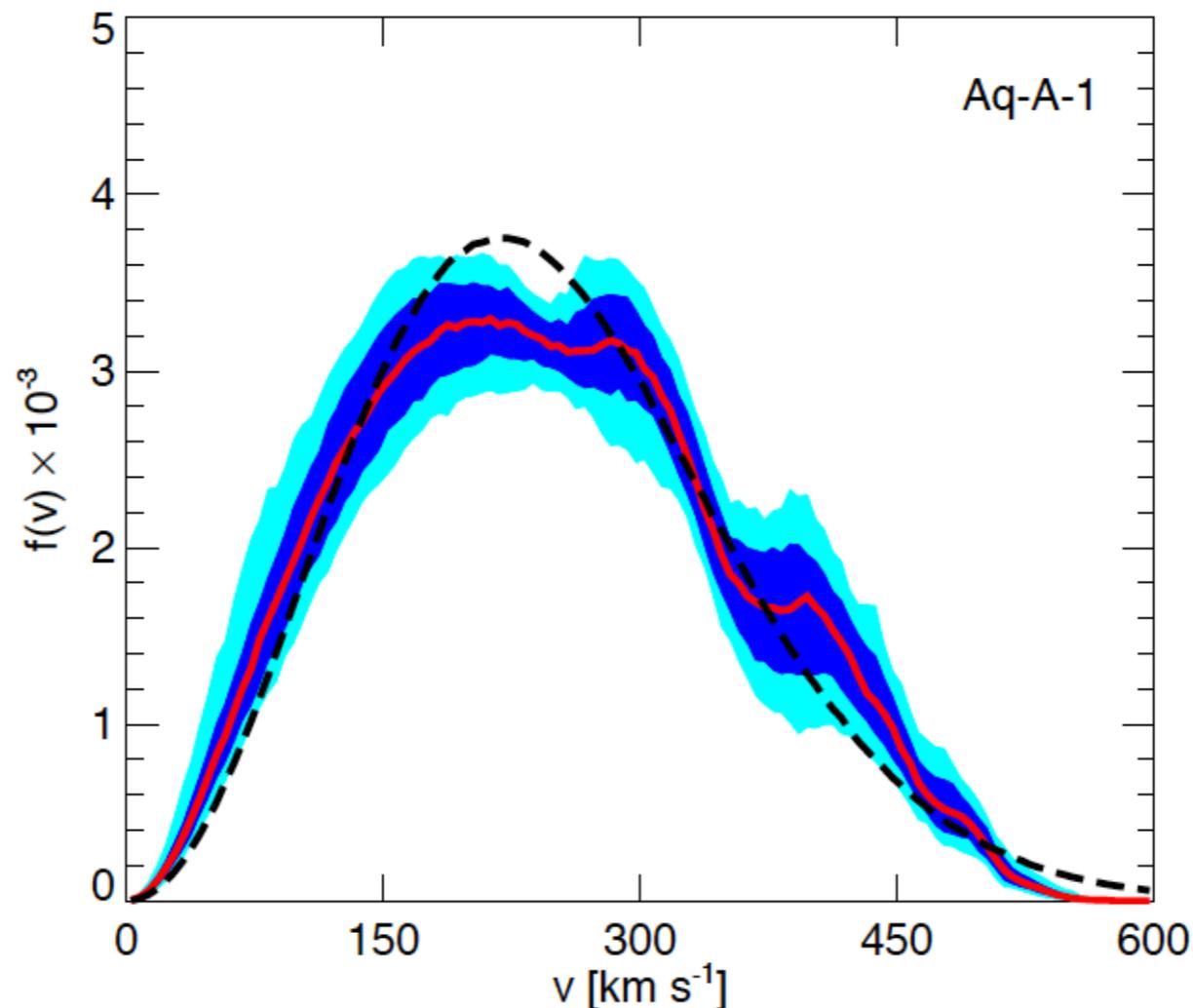
with  $v_c = 220$  km/s and  $v_{\text{esc}} = 550$  km/s.

$\sigma_v = \sqrt{3/2} v_c$  independent of radius.

- What can we learn from numerical simulations of galaxy formation about the local DM velocity distribution?

# Dark Matter only simulations

- DM speed distributions from cosmological N-body simulations **without baryons**, deviate substantially from a Maxwellian.



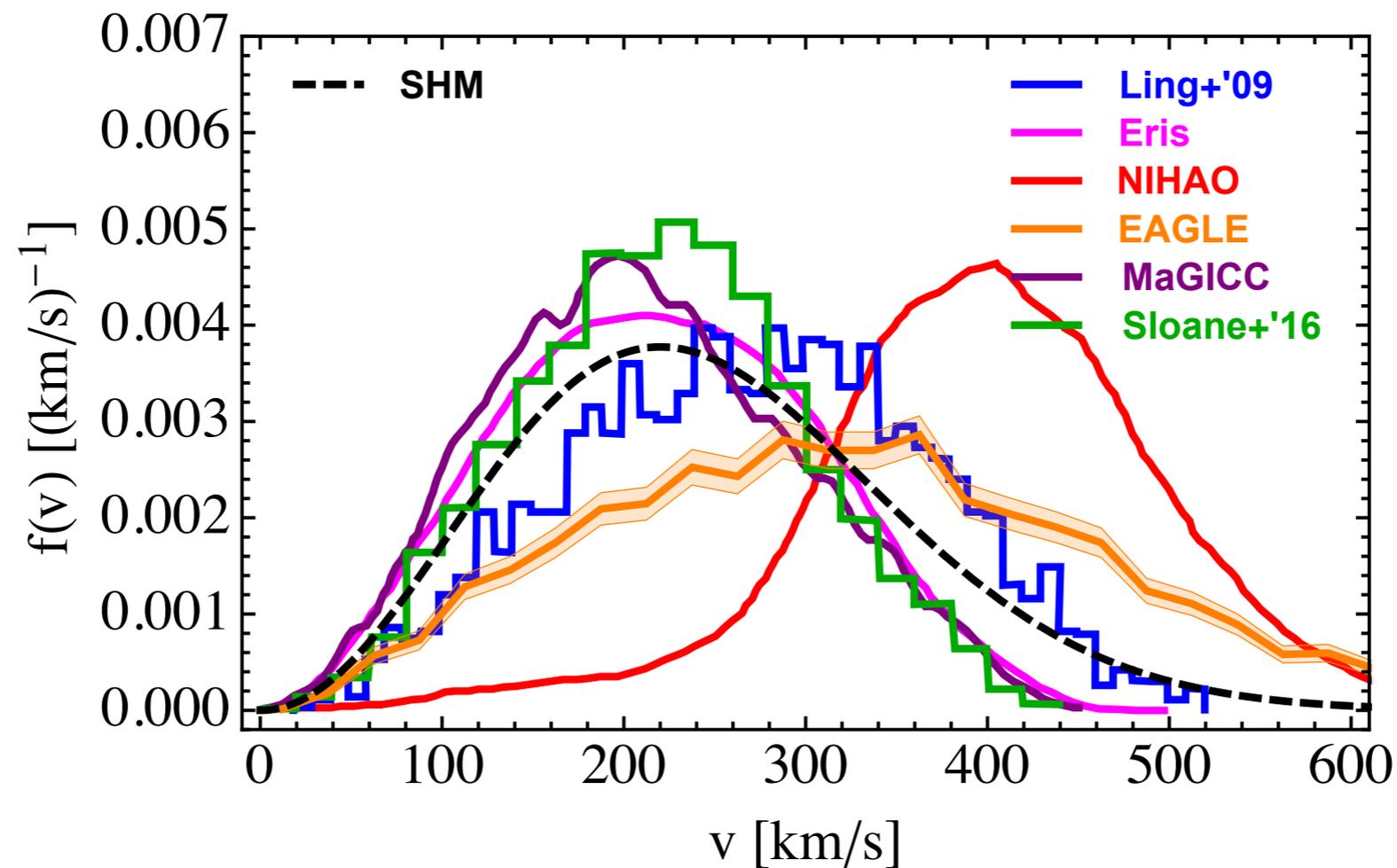
$$f(|\mathbf{v}|) = v^2 \int d\Omega_{\mathbf{v}} f(\mathbf{v})$$

Vogelsberger et al., 0812.0362

- Significant systematic uncertainty since the impact of baryons neglected.*

# Hydrodynamical simulations

- Each hydrodynamical (**DM + baryons**) simulation adopts a different *galaxy formation model, spatial resolution, DM particle mass*.

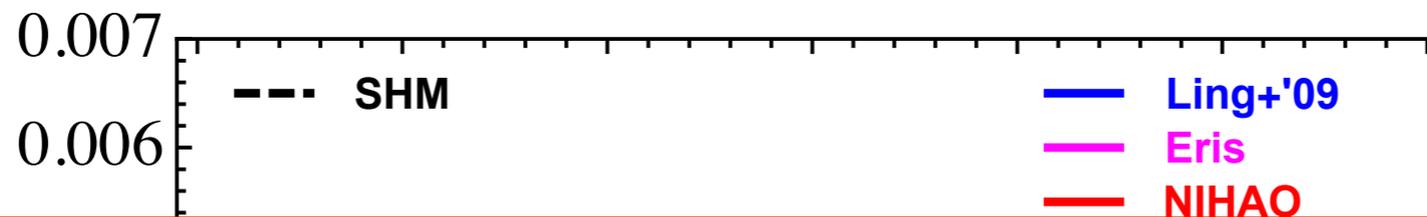


Bozorgnia & Bertone, 1705.05853

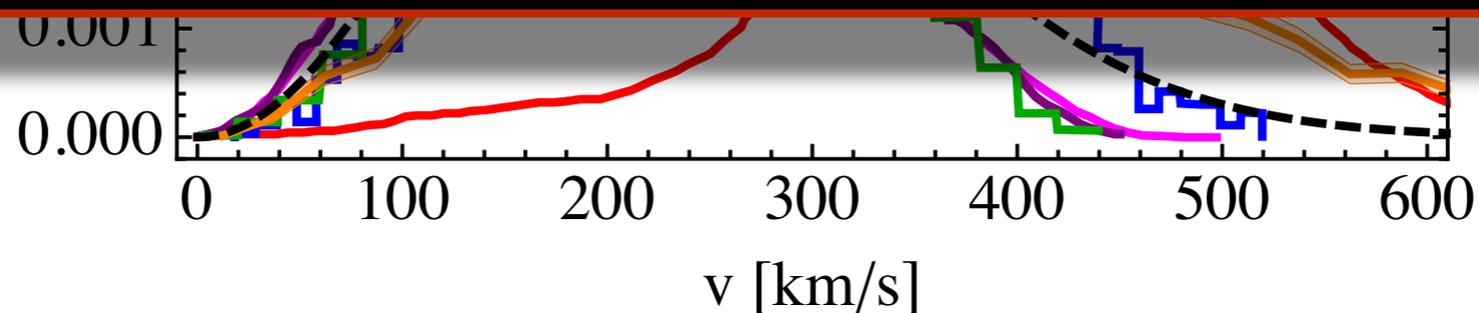
- Large variation in DM speed distributions between the results of different simulations.

# Hydrodynamical simulations

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Different criteria used to identify MW-like galaxies among different groups. The most common criteria is the MW mass constraint, which has a large uncertainty.



Bozorgnia & Bertone, 1705.05853

- Large variation in DM speed distributions between the results of different simulations.

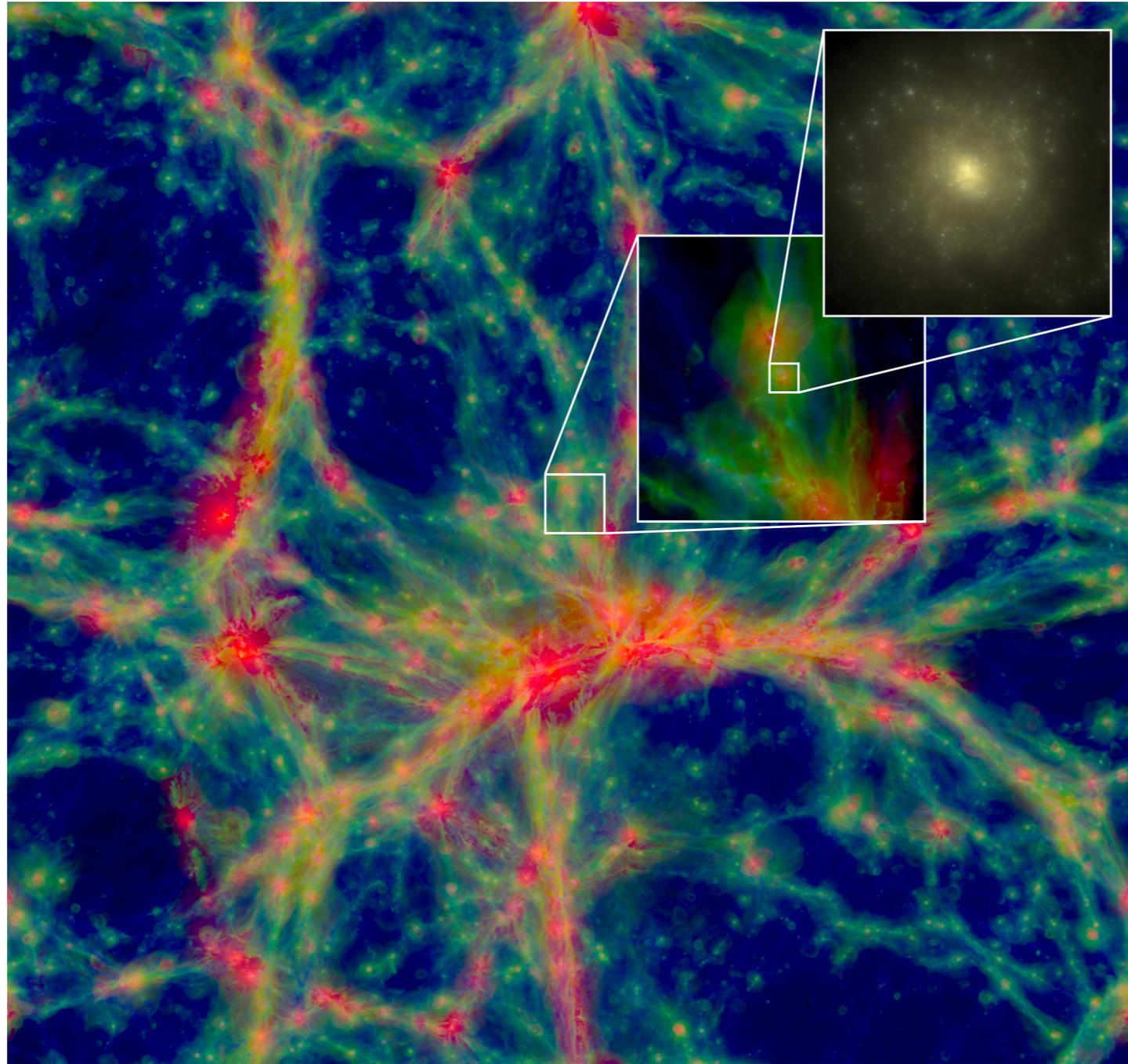
# Hydrodynamical simulations

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  - Model baryonic processes in a way that the main galaxy population properties are broadly reproduced.
  - Identify MW-like galaxies by taking into account observational constraints on the MW.

# Hydrodynamical simulations

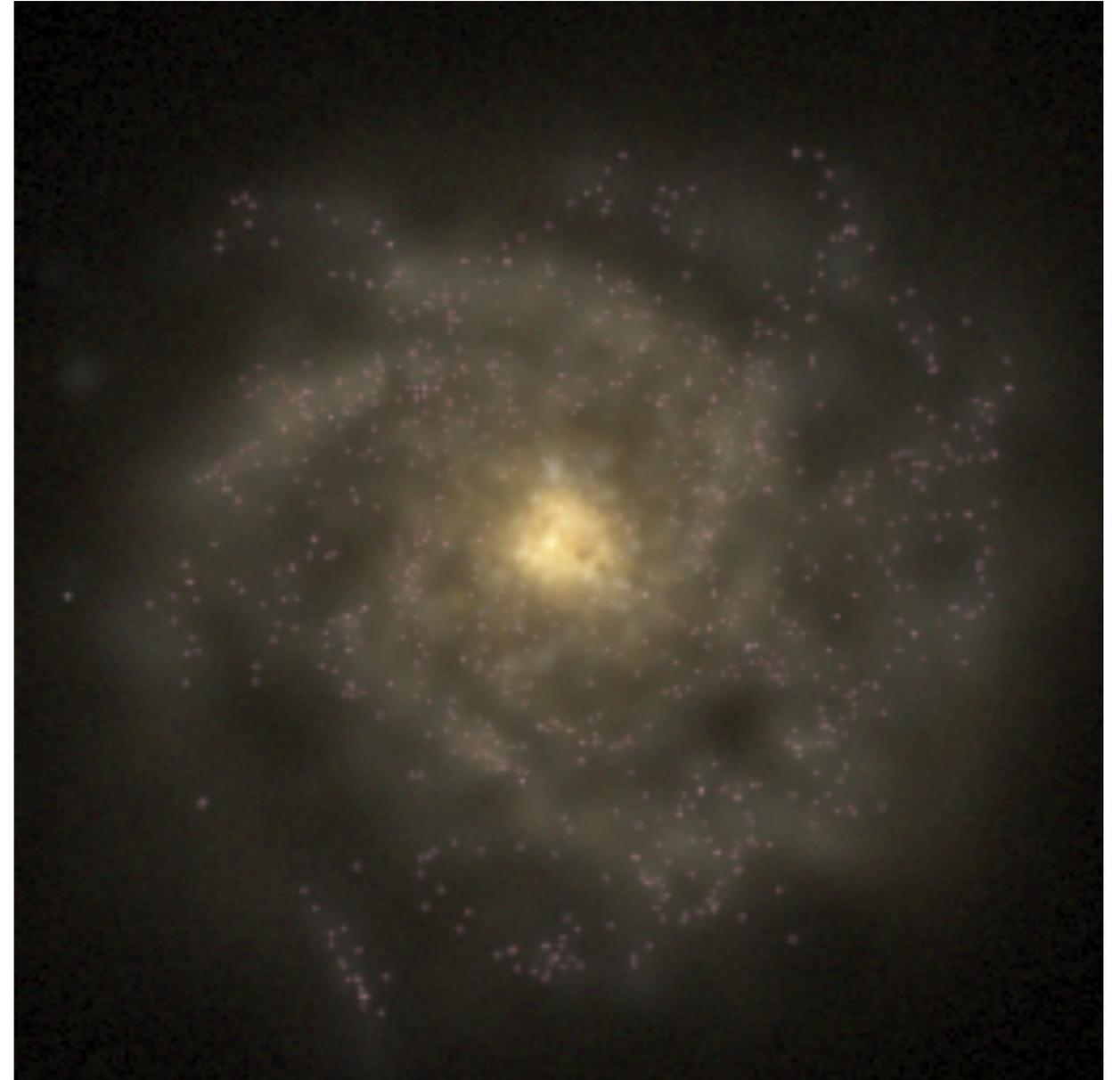
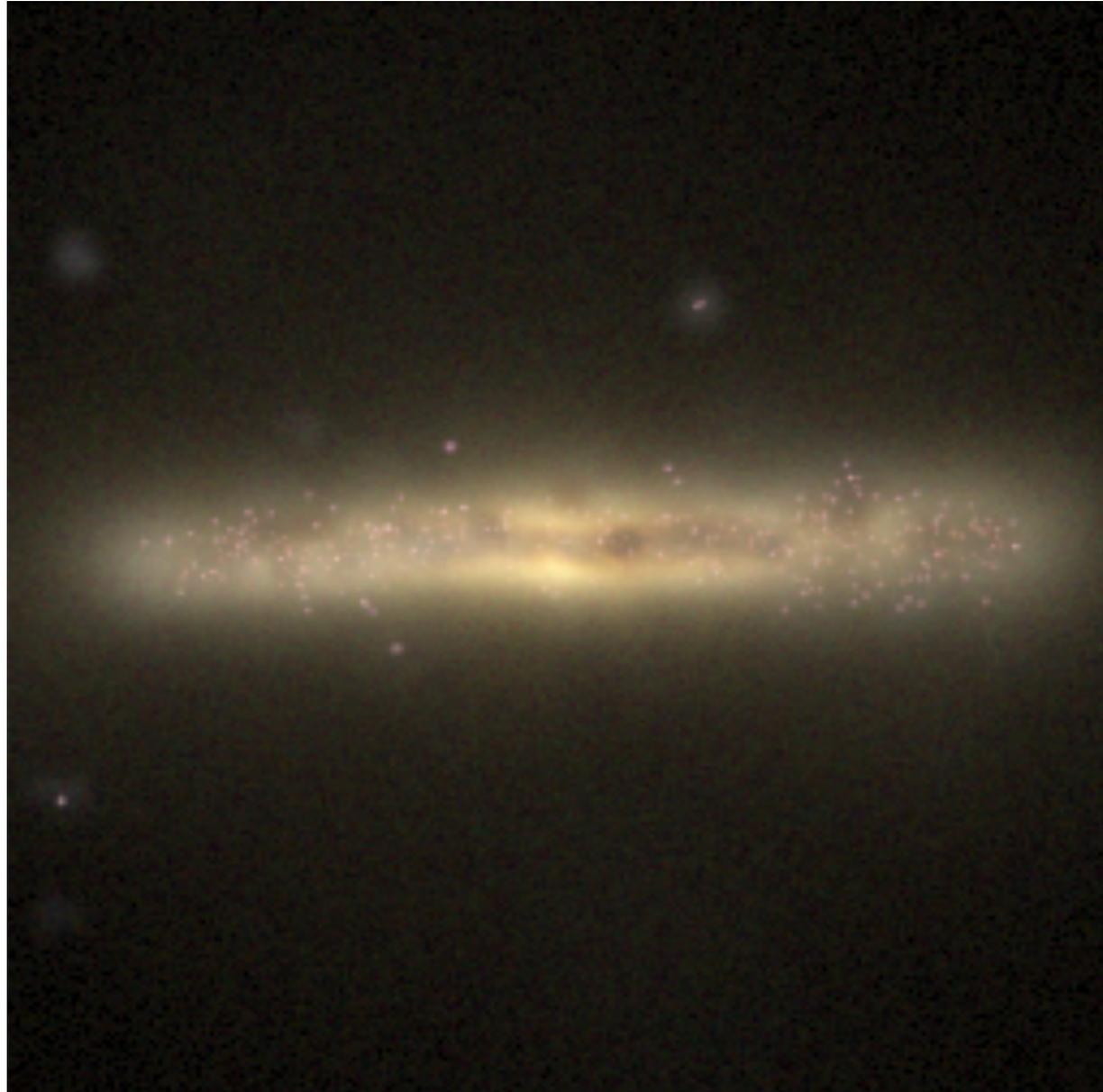
- To make precise quantitative predictions:
  - Model baryonic processes in a way that the main galaxy population properties are broadly reproduced.
  - Identify MW-like galaxies by taking into account observational constraints on the MW.
- We use the **EAGLE** and **APOSTLE** hydrodynamic simulations. *calibrated to reproduce the observed distribution of stellar masses and sizes of low-redshift galaxies.*
- Companion Dark Matter only (DMO) simulations were run assuming all the matter content is collisionless.

# EAGLE Simulations



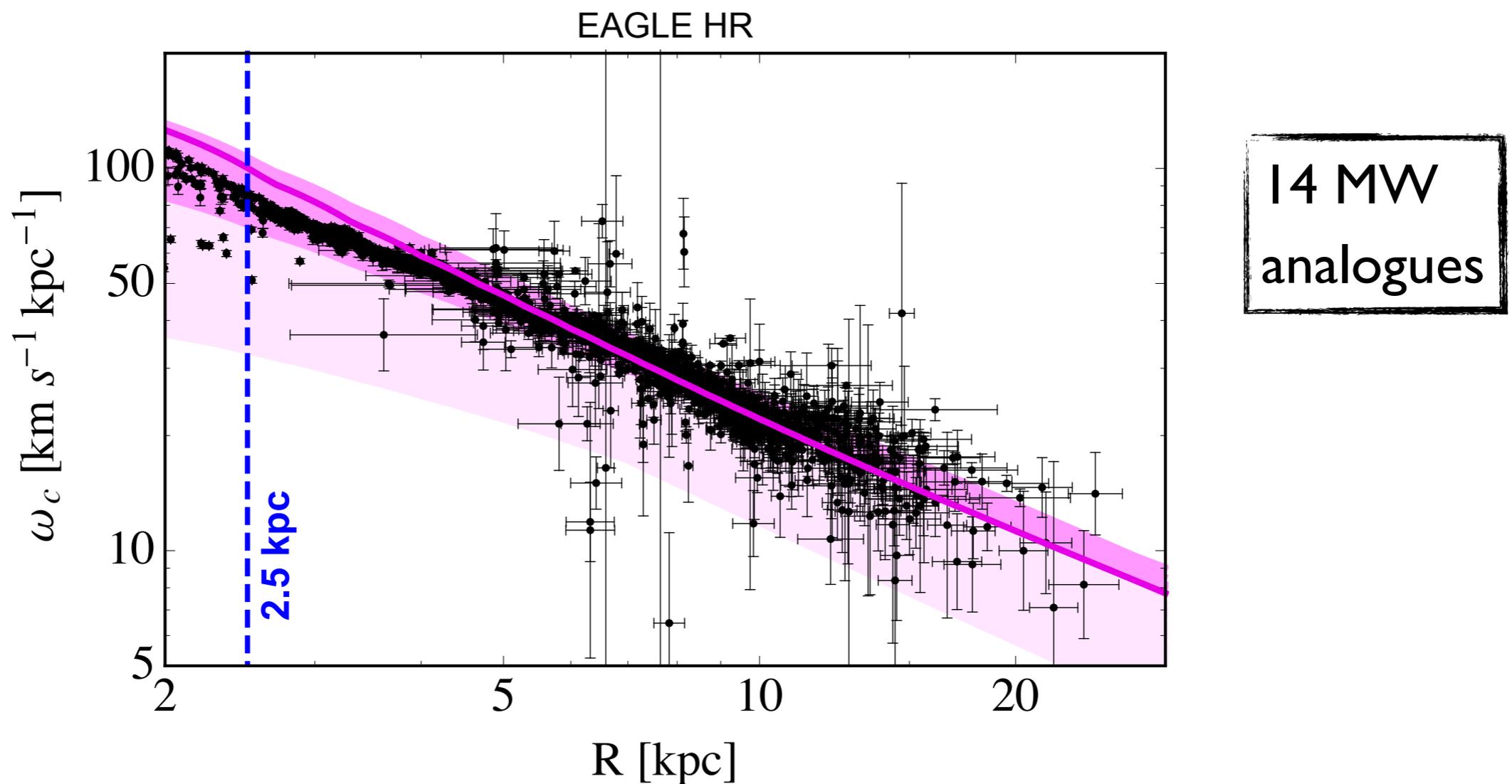
**EAGLE Simulations, I407.7040**

# Milky Way analogues



# Identifying Milky Way analogues

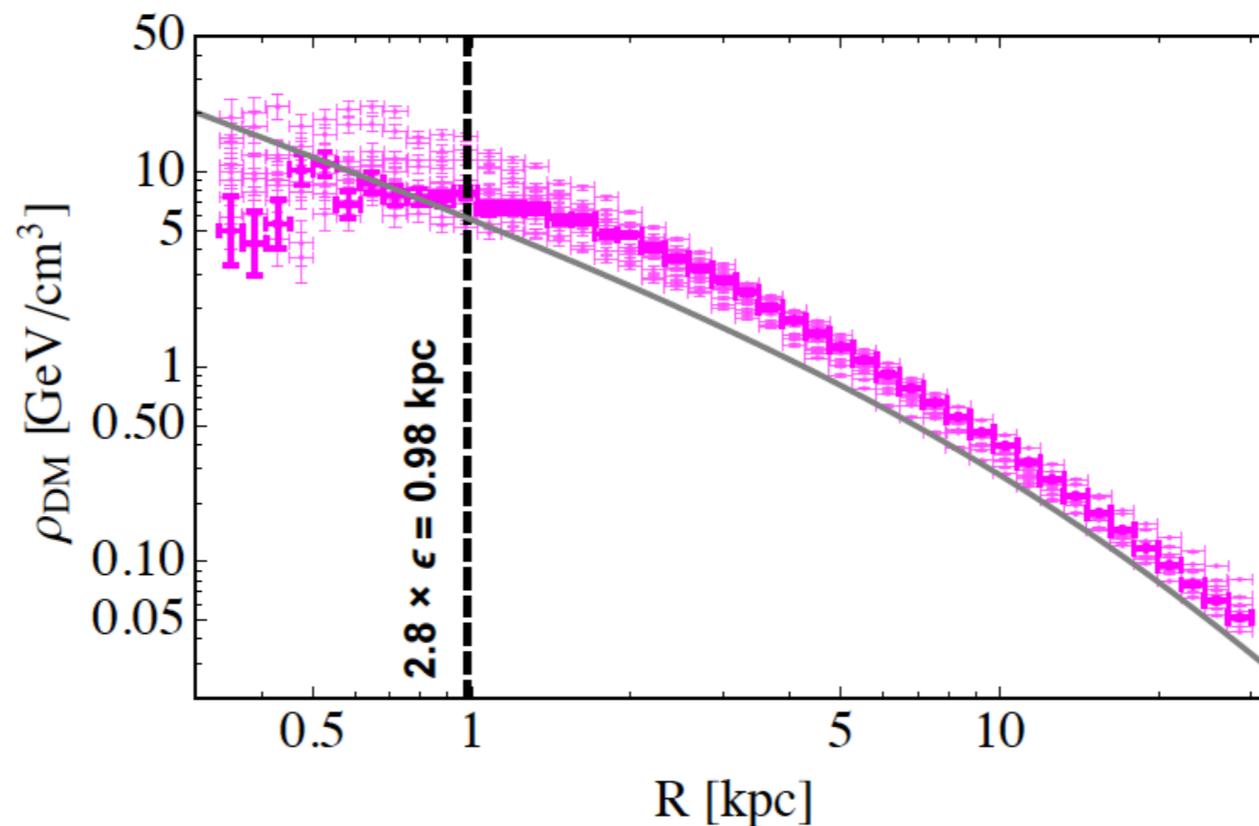
- We introduce new criteria to identify MW analogues using observed MW kinematical data: **rotation curves, total stellar mass.**



Bozorgnia et al., [1601.04707](#)  
Calore, Bozorgnia et al., [1509.02164](#)

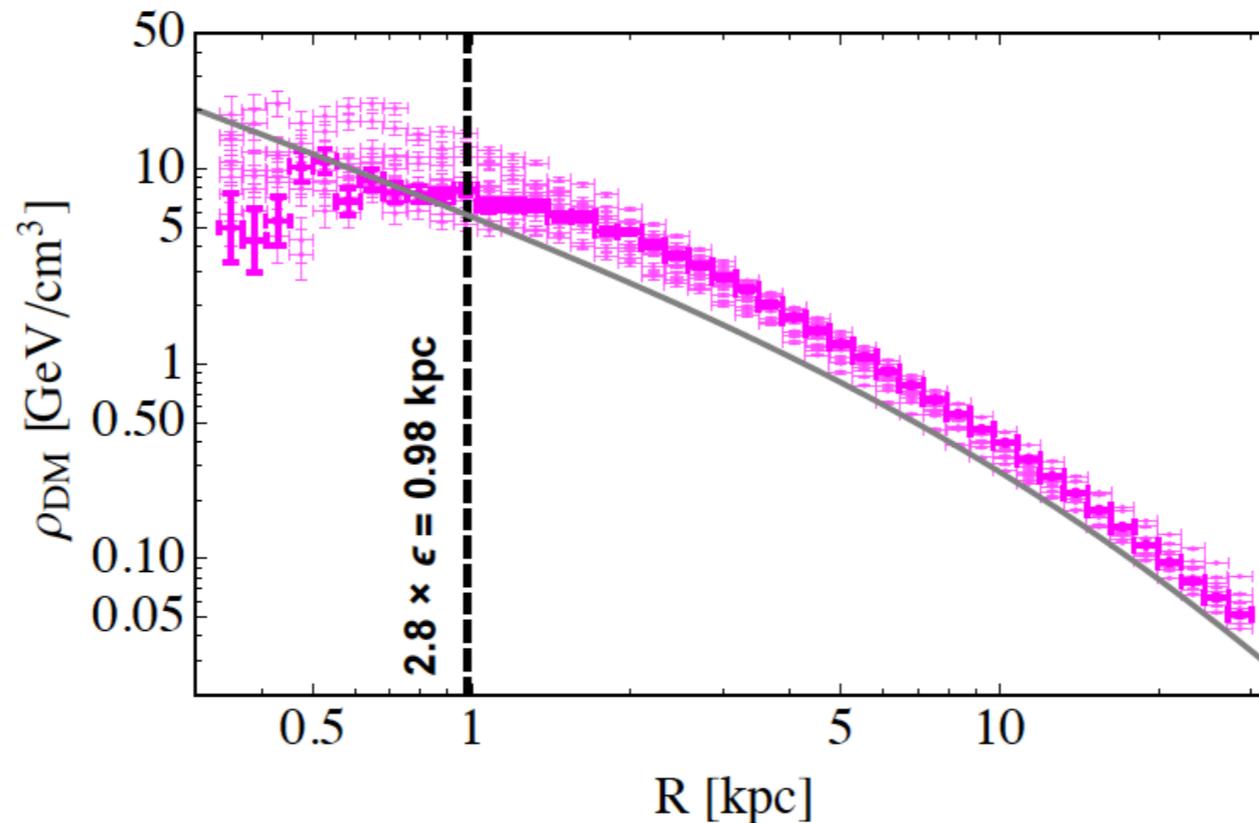
# Dark Matter density profiles

- Spherically averaged DM density profiles of the MW analogues:



# Dark Matter density profiles

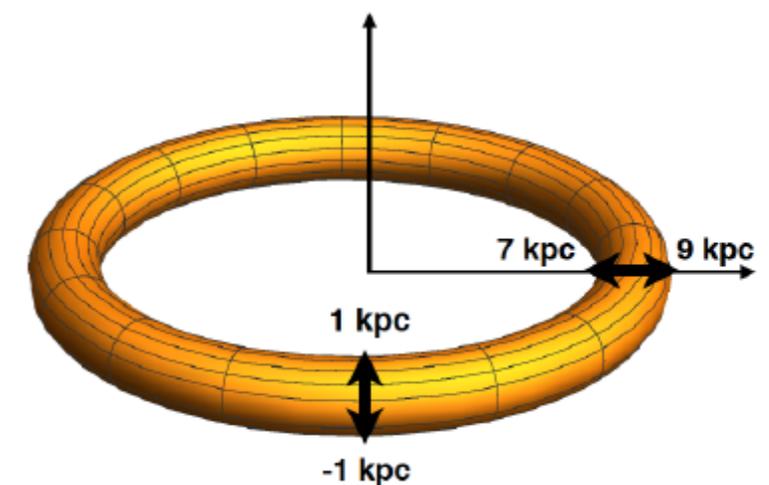
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- To find the DM density at the position of the Sun, consider a torus aligned with the stellar disc.

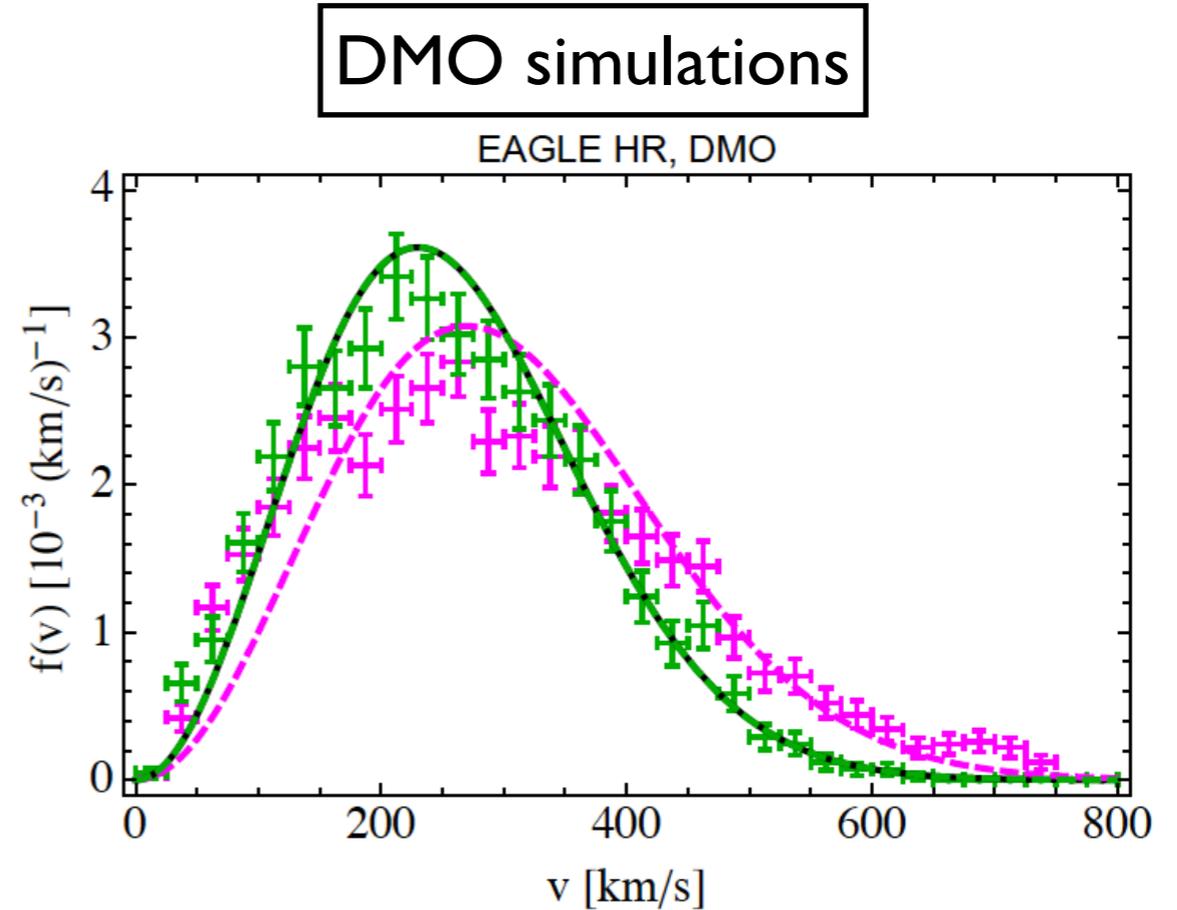
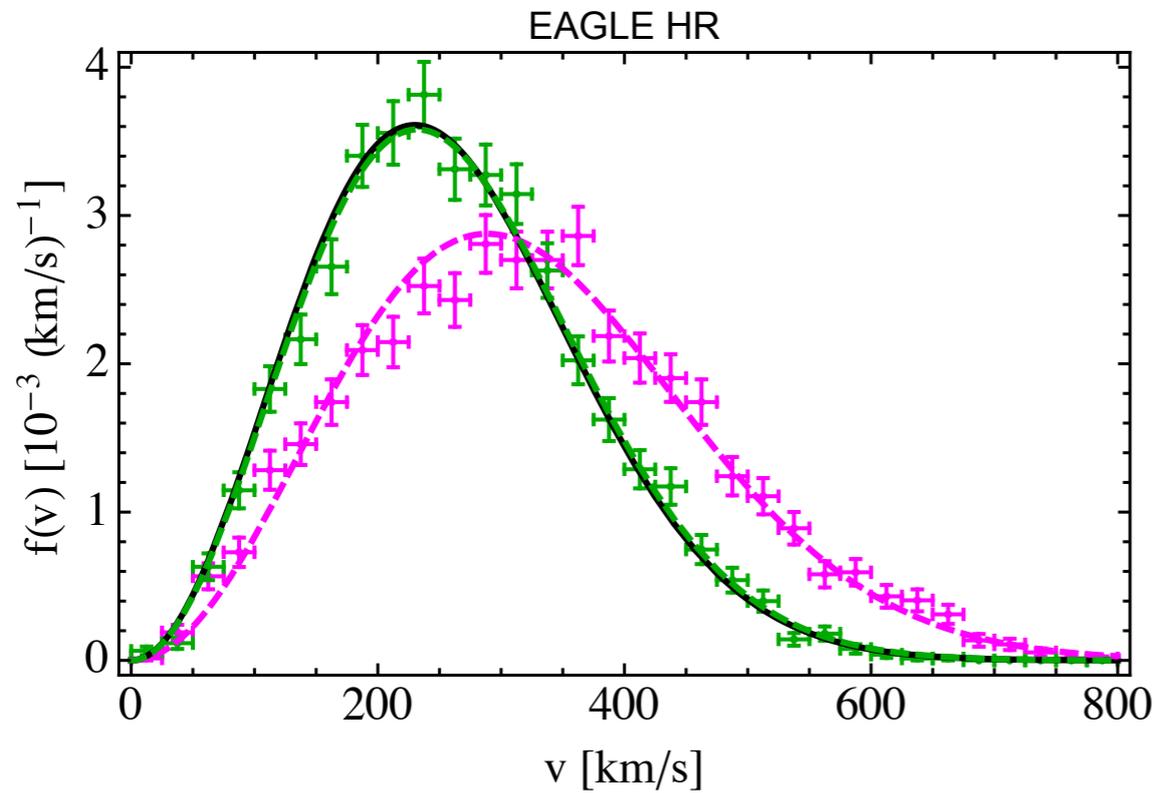
$$\rho_\chi = 0.41 - 0.73 \text{ GeV/cm}^3$$

Bozorgnia et al., 1601.04707



# Local speed distributions

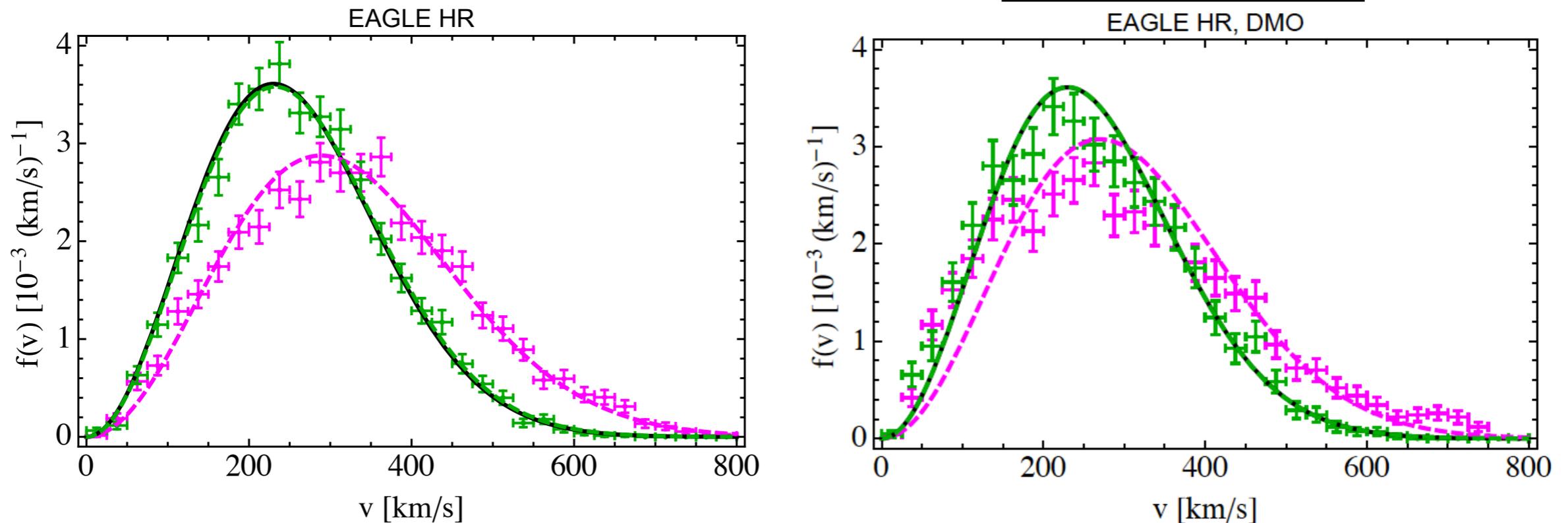
In the galactic rest frame:



Bozorgnia et al., 1601.04707

# Local speed distributions

In the galactic rest frame:



Bozorgnia et al., 1601.04707

- Maxwellian distribution with a free peak provides a better fit to haloes in the hydrodynamical simulations compared to their DMO counterparts.
- Best fit peak speed:  $v_{\text{peak}} = 223 - 289 \text{ km/s}$

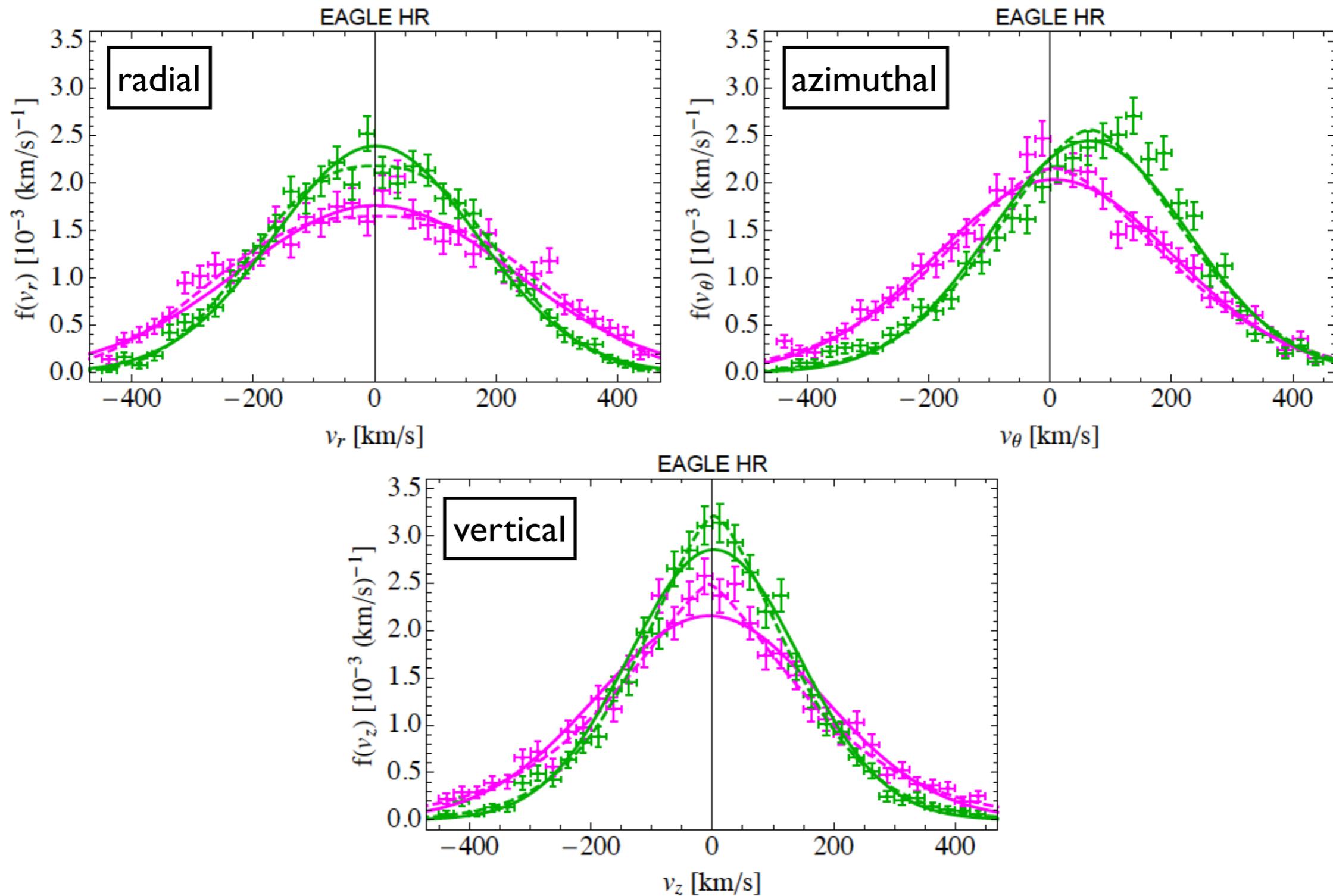
# Local speed distributions

## Common trends in different hydrodynamical simulations:

- Baryons deepen the gravitational potential in the inner halo, shifting the peak of the DM speed distribution to *higher speeds*.
- In most cases, baryons appear to make the local DM speed distribution *more Maxwellian*.

Bozorgnia & Bertone, 1705.05853

# Components of the velocity distribution



Bozorgnia et al., I601.04707

# How common are dark disks?

- Clear velocity anisotropy at the Solar circle.
- Two haloes have a rotating DM component in the disc with mean velocity comparable (within 50 km/s) to that of the stars.

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- Hint for the existence of a co-rotating dark disk in 2 out of 14 MW-like haloes. → Dark disks are relatively rare in our halo sample.

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Schaller et al., 1605.02770

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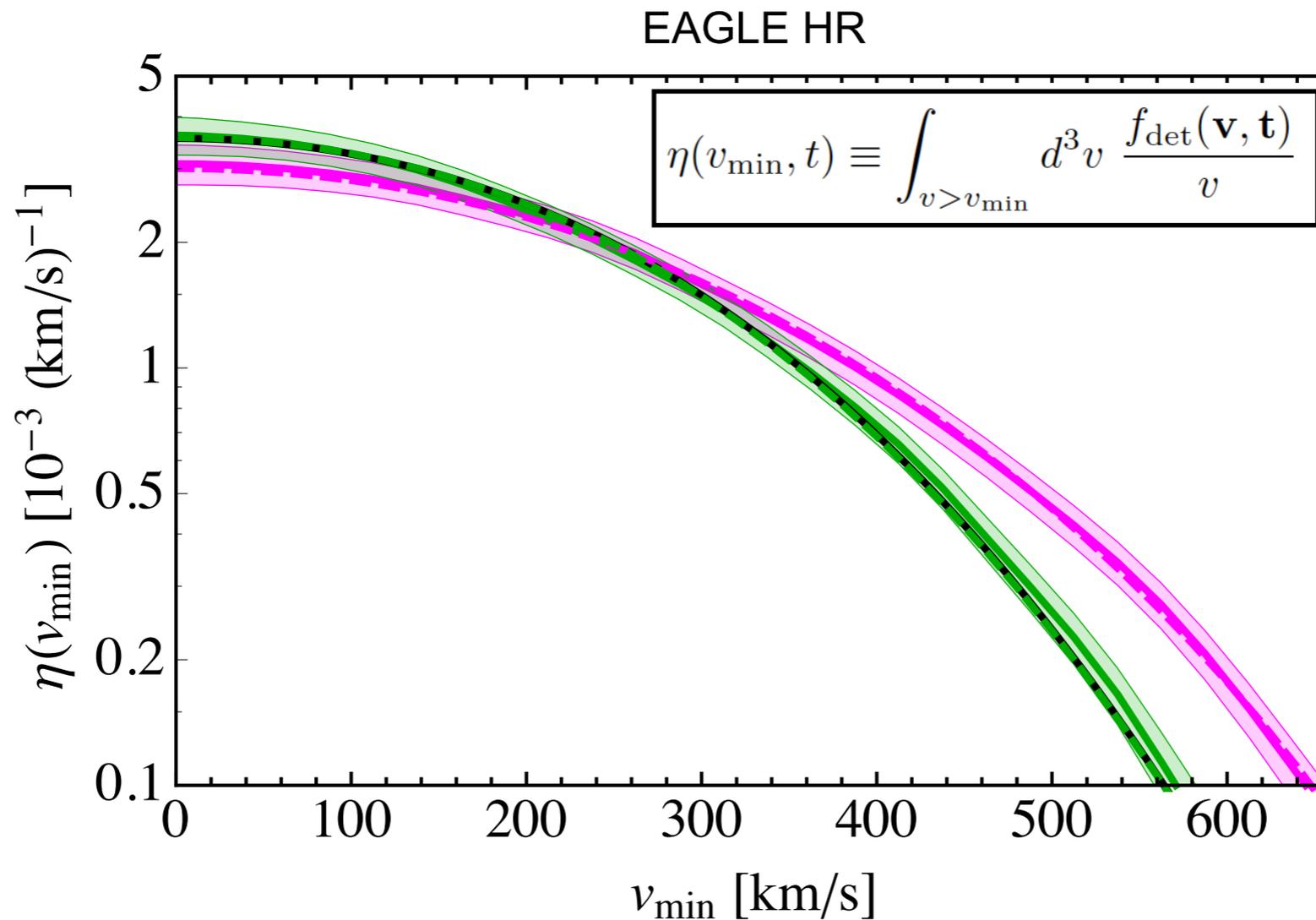
Bozorgnia et al., 1601.04707

Schaller et al., 1605.02770

- *Sizable dark disks also rare in other hydro simulations:*
  - They only appear in simulations where a large satellite merged with the MW in the recent past, which is robustly excluded from MW kinematical data.

Bozorgnia & Bertone, 1705.05853

# The halo integral



- Halo integrals for the best fit Maxwellian velocity distribution (*peak speed 223 - 289 km/s*) fall within the  $1\sigma$  uncertainty band of the halo integrals of the simulated haloes.

Bozorgnia et al., 1601.04707

# The halo integral

Common trend in different hydrodynamical simulations:

- Halo integrals and hence direct detection event rates obtained from a **Maxwellian velocity distribution with a free peak** are similar to those obtained directly from the simulated haloes.

Bozorgnia et al., [1601.04707](#) (EAGLE & APOSTLE)

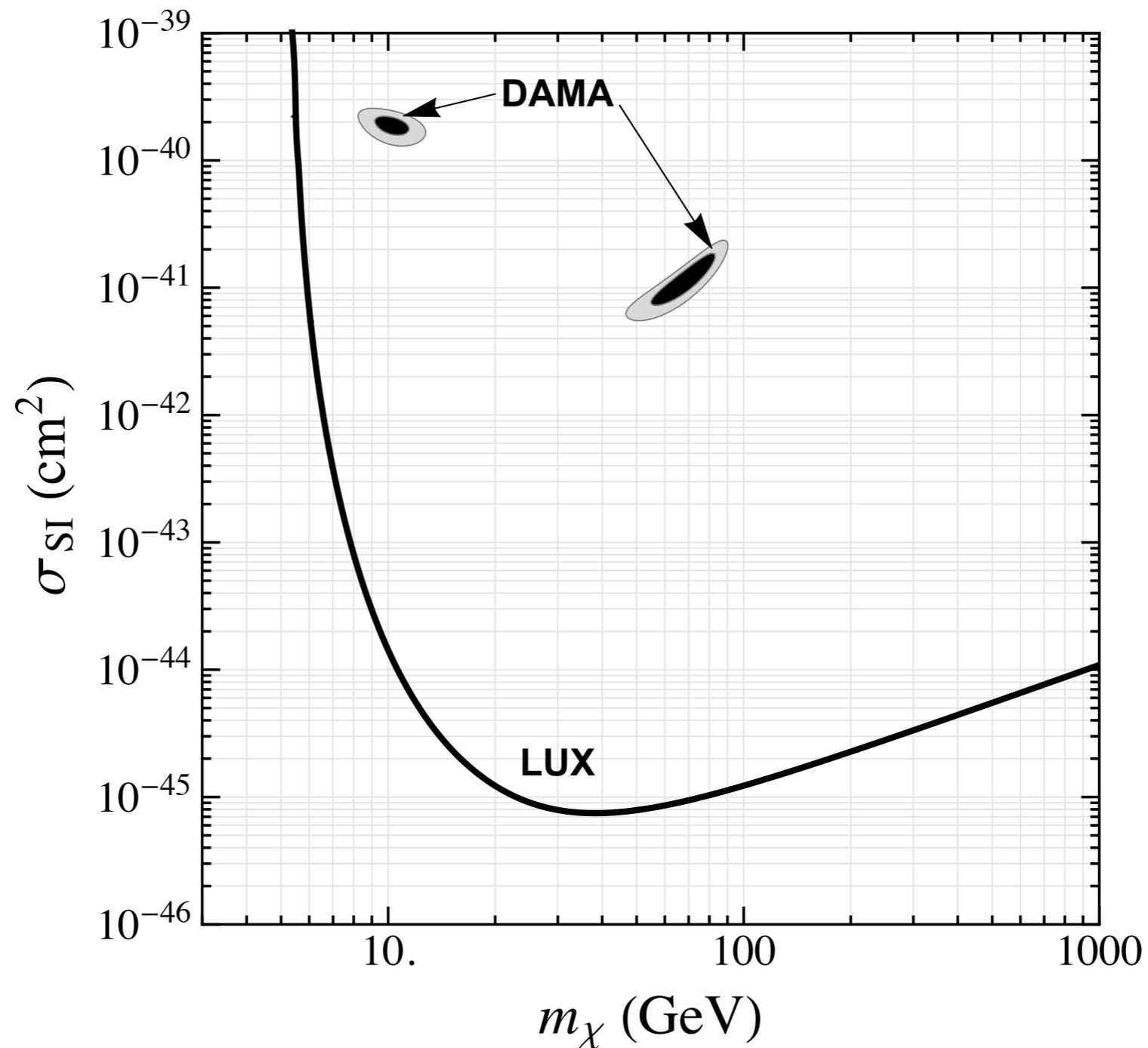
Kelso et al., [1601.04725](#) (MaGICC)

Sloane et al., [1601.05402](#)

Bozorgnia & Bertone, [1705.05853](#)

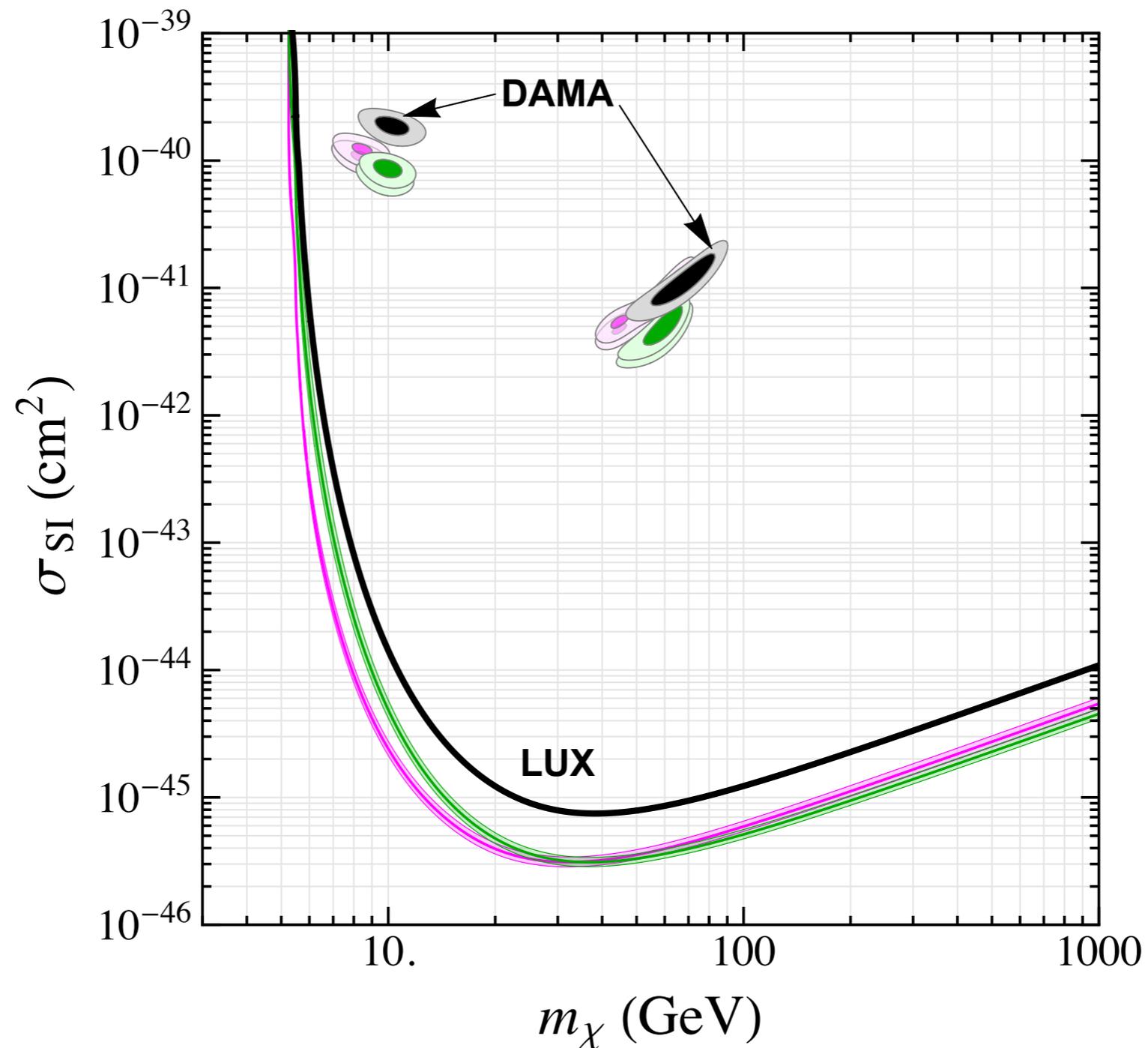
# Implications for direct detection

- Assuming the **Standard Halo Model**:

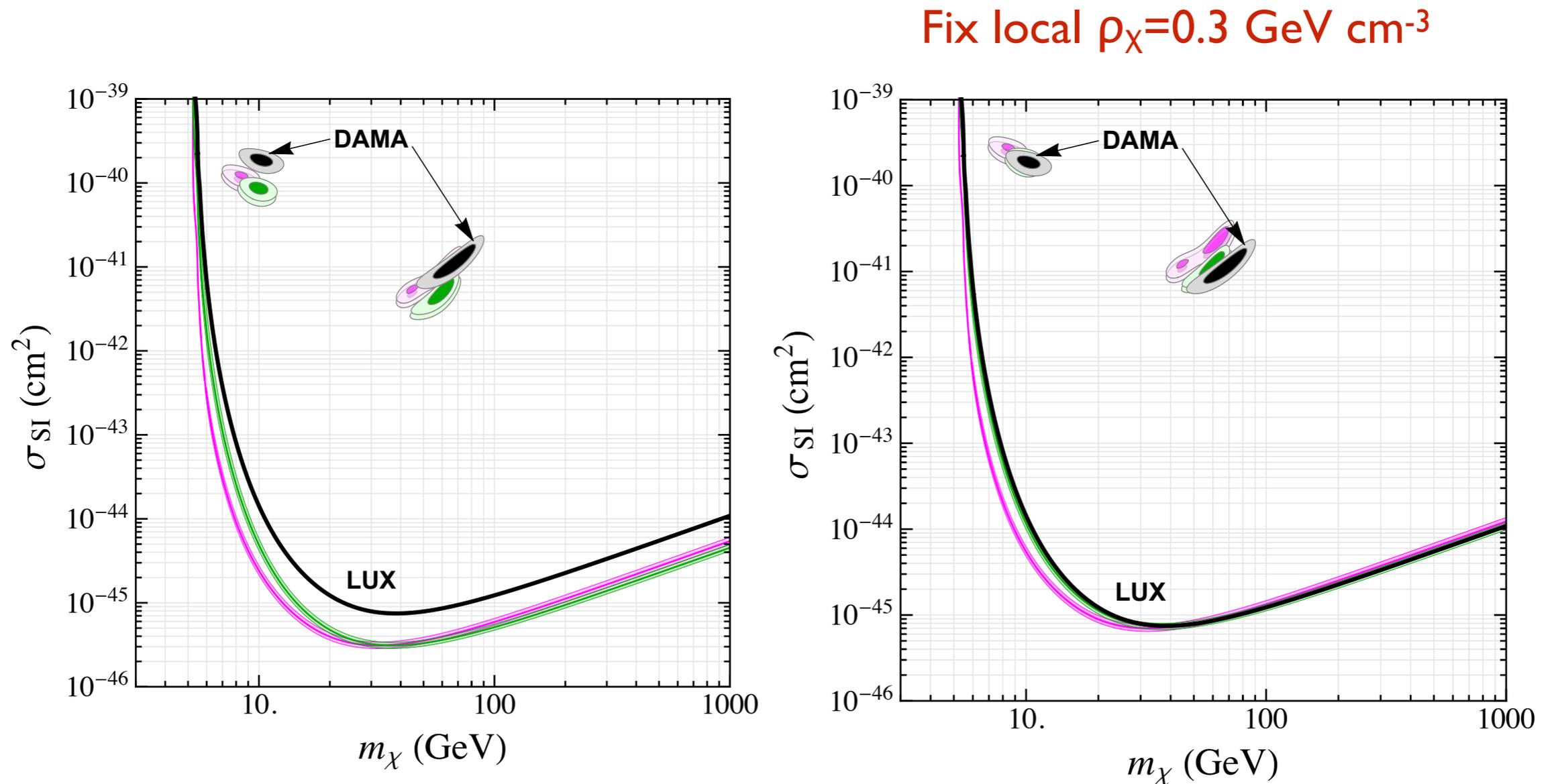


# Implications for direct detection

- Compare with simulated Milky Way-like haloes:



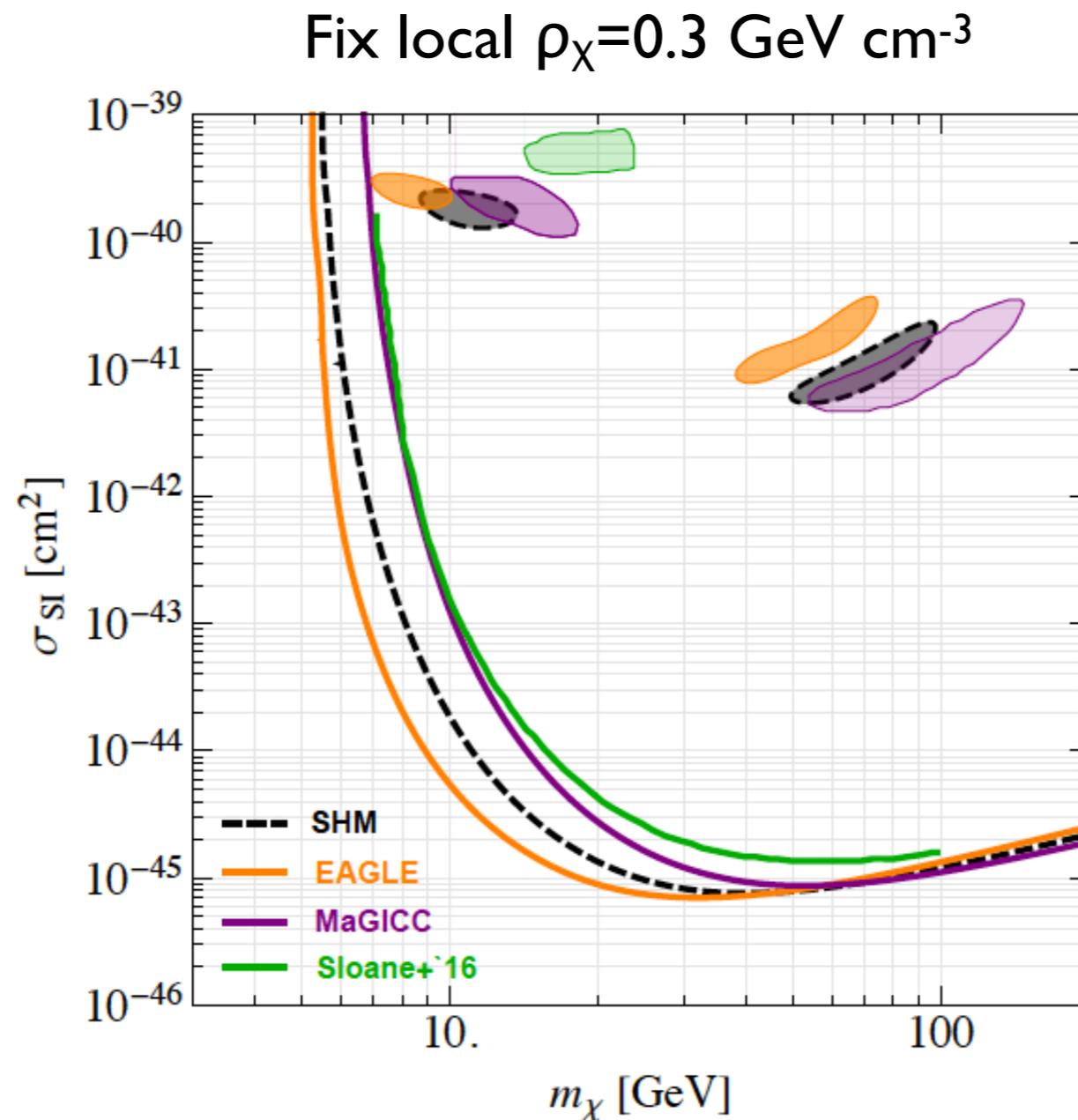
# Implications for direct detection



- Difference in the local DM density  $\rightarrow$  overall difference with the SHM.
- Variation in the peak of the DM speed distribution  $\rightarrow$  shift in the low mass region.

# Implications for direct detection

Comparison to other hydrodynamical simulations:



Bozorgnia & Bertone, 1705.05853

# Non-standard interactions

- For a very general set of non-relativistic effective operators:

Kahlhoefer & Wild, 1607.04418

$$\frac{d\sigma_{\chi N}}{dE_R} = \frac{d\sigma_1}{dE_R} \frac{1}{v^2} + \frac{d\sigma_2}{dE_R}$$

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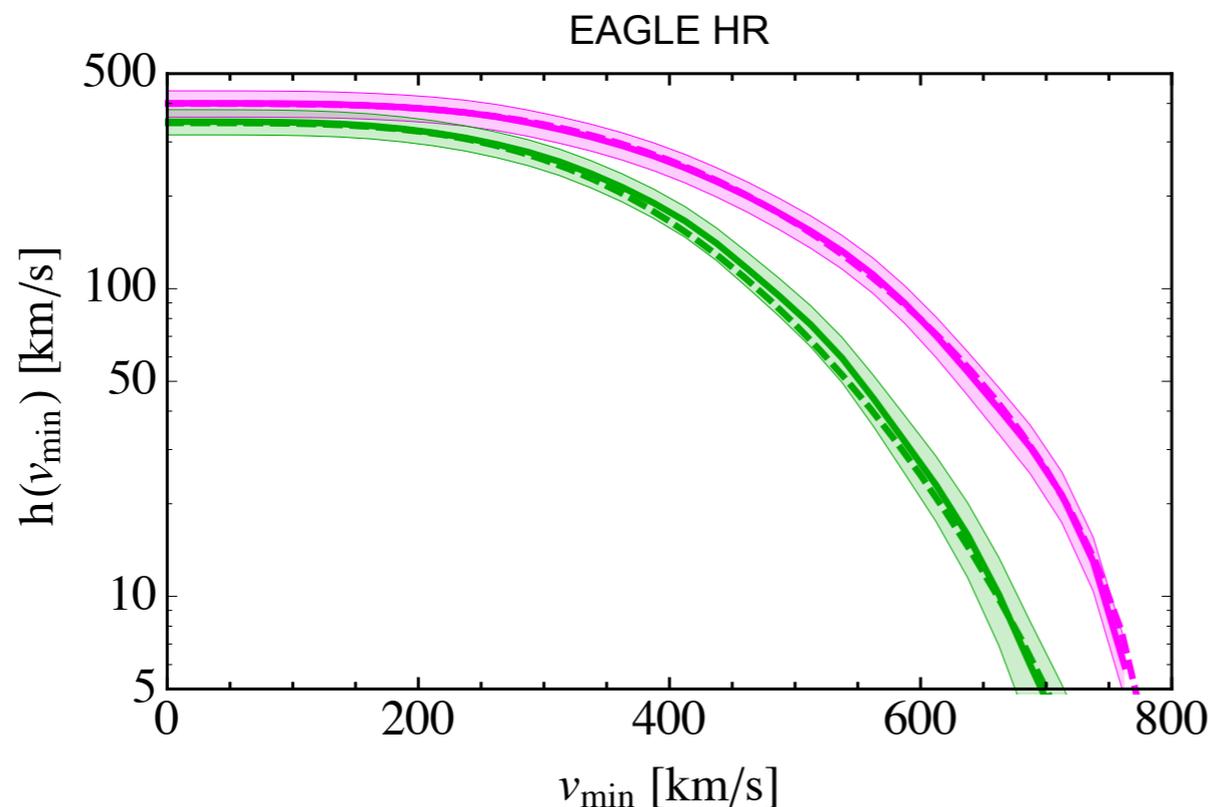
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Bozorgnia & Bertone, 1705.05853

- Best fit Maxwellian  $h(v_{\min})$  falls within the  $1\sigma$  uncertainty band of the  $h(v_{\min})$  of the simulated haloes.

# Prospects for indirect DM searches

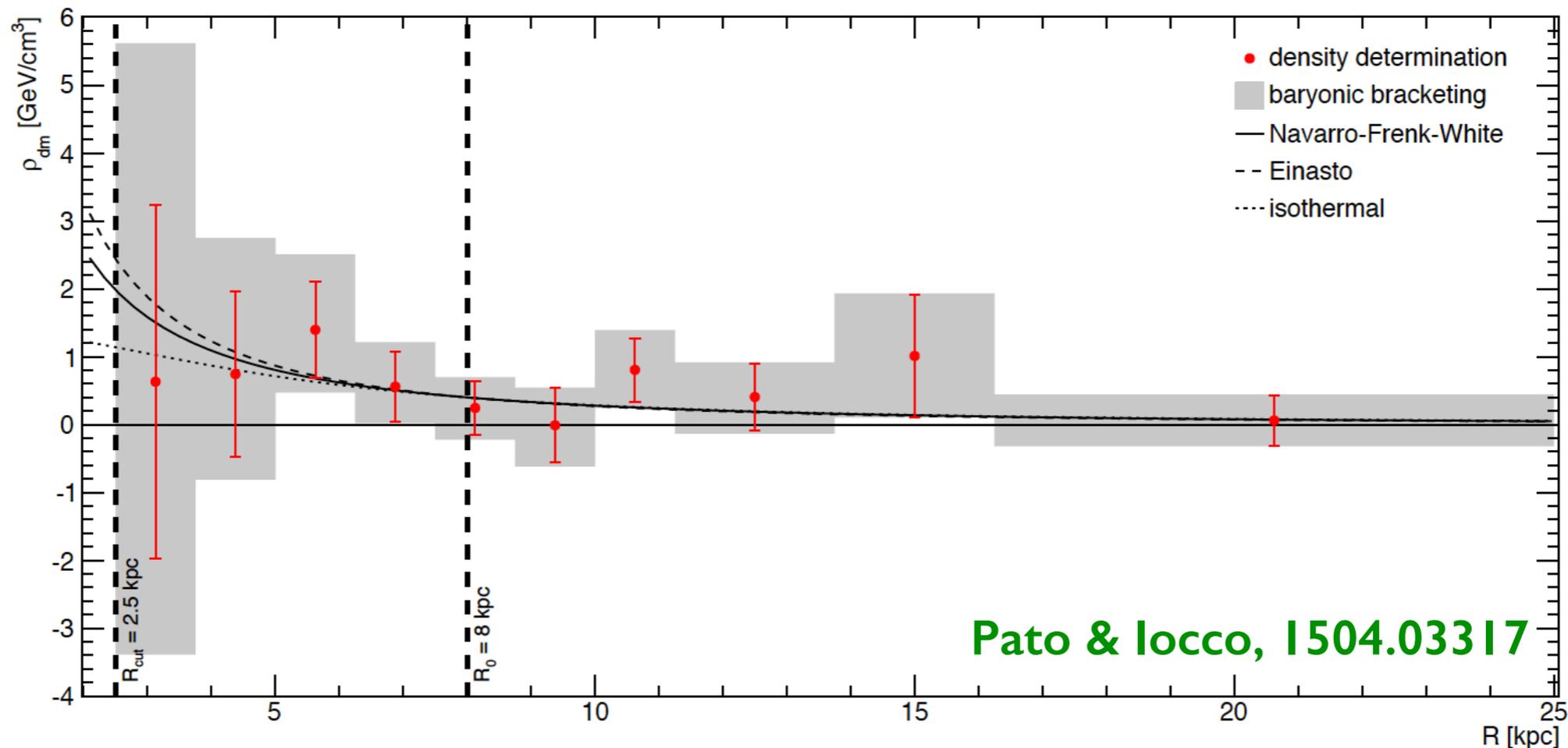
# Indirect DM searches

- Expected gamma-ray flux from DM annihilation:

$$\frac{d\Phi_\gamma}{dE} = \frac{\langle\sigma v\rangle}{8\pi m_\chi^2} \frac{dN_\gamma}{dE} \int_{\text{l.o.s.}} ds \rho^2(r(s, \psi))$$

astrophysics

- Large uncertainties in the DM density profile in the inner few kpc.



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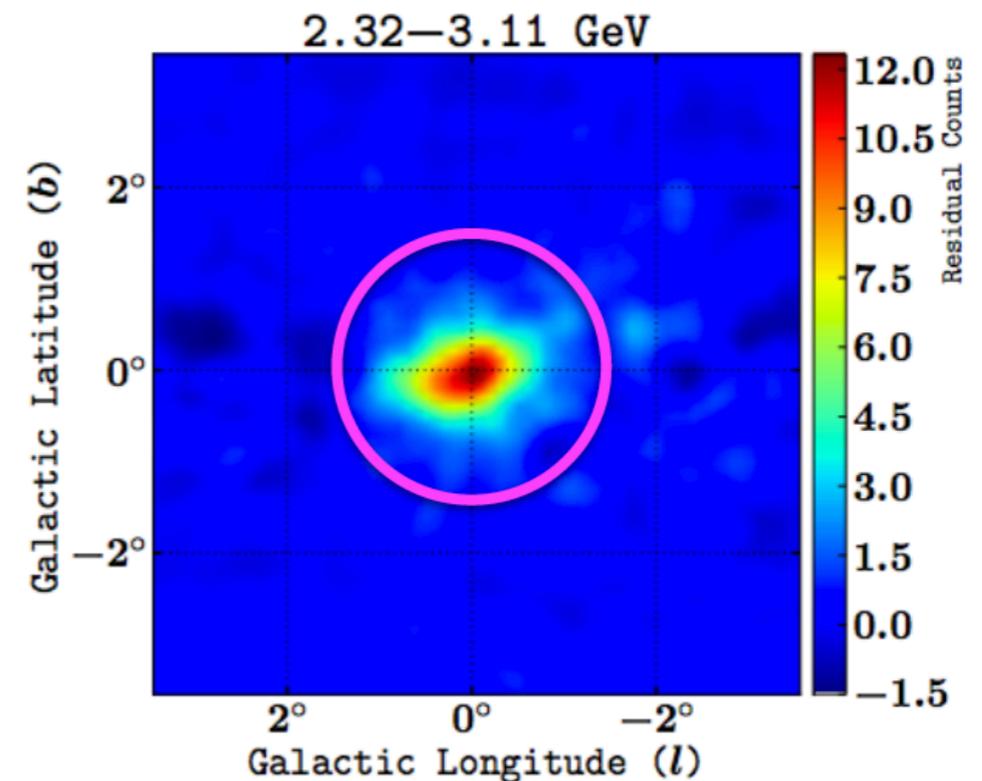
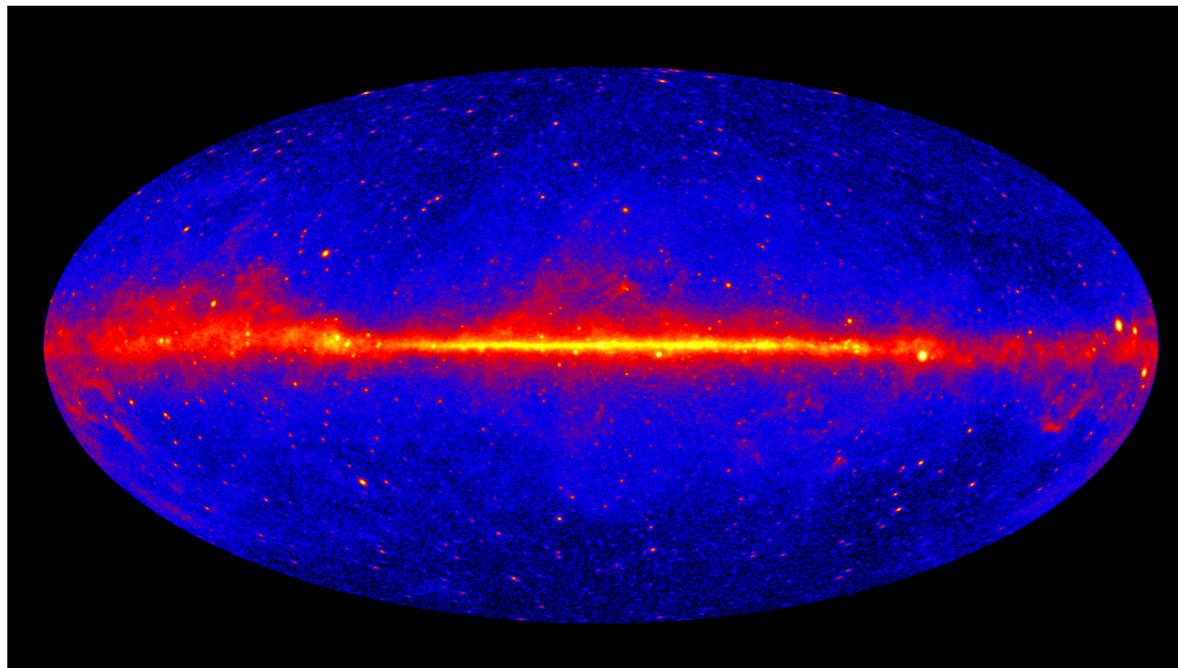
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## Use cosmological simulations:

- DMO simulations predict NFW profile:  $r^{-\gamma}$ , where  $\gamma \approx 1$  in the inner few kpc.
- What is the DM density profile for **MW-like galaxies** in hydrodynamical simulations?

# Galactic centre GeV excess

- Unexplained excess of gamma rays in Fermi-LAT data from the centre of our Galaxy, above the known astrophysical background. **Hooper & Goodenough '09, Vitale & Morselli '09, ...**



**Macias & Gordon, 1312.6671**

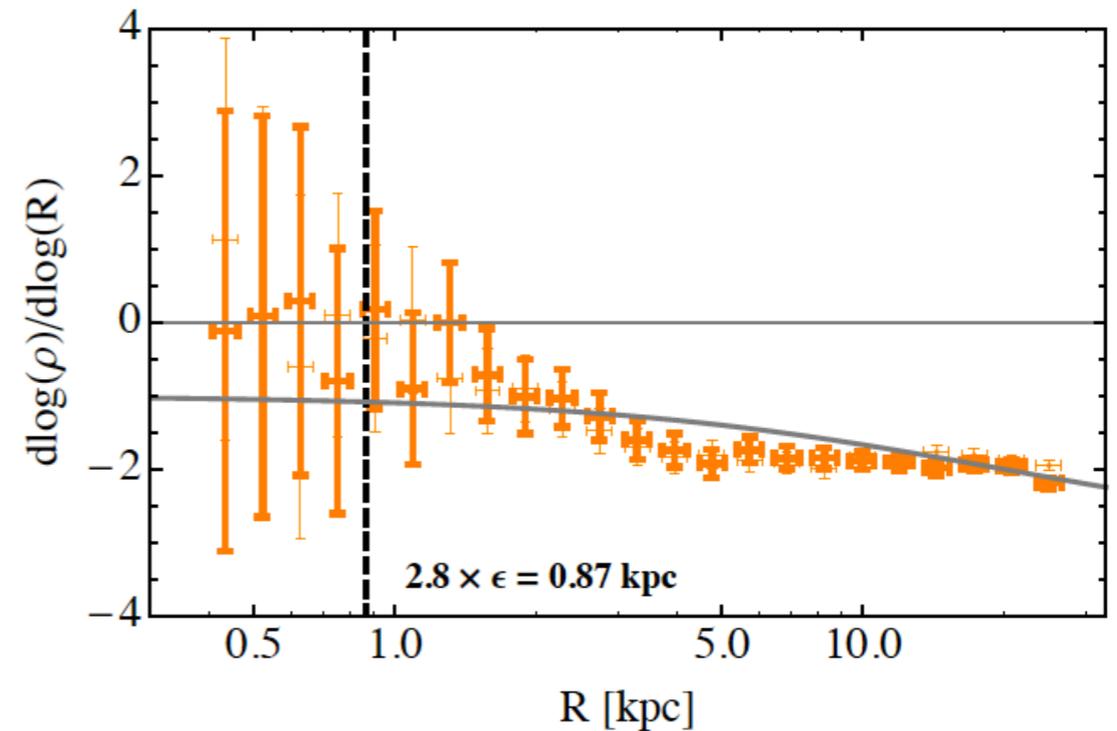
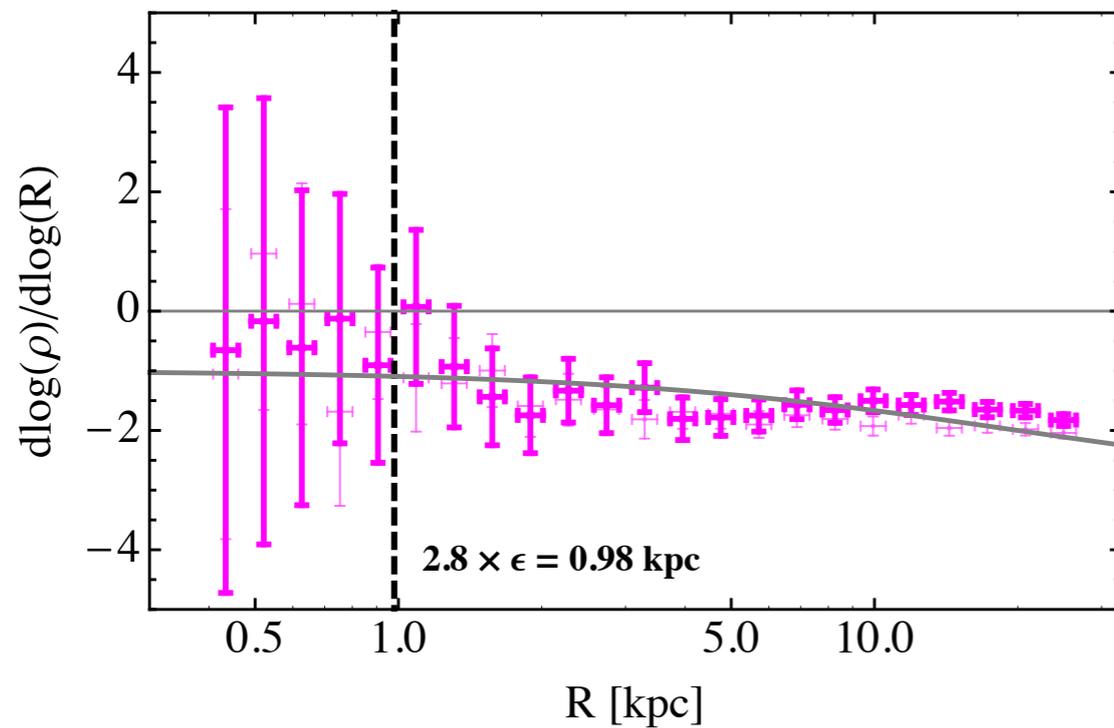
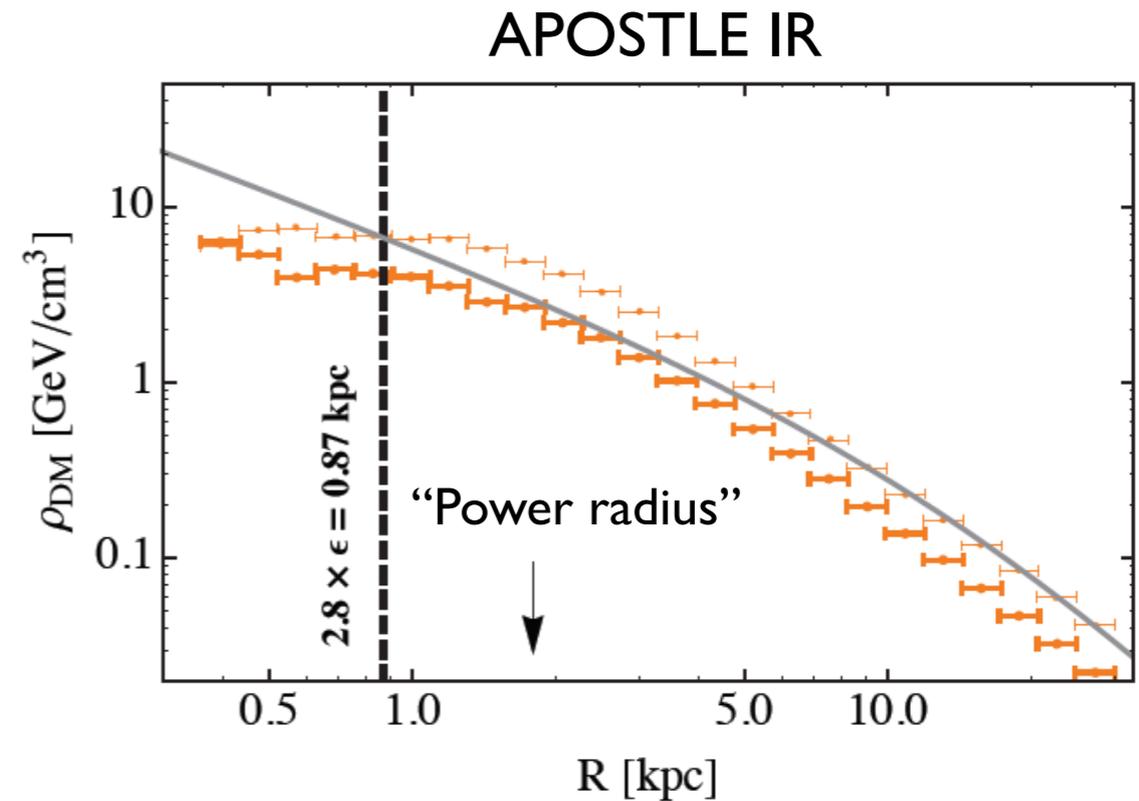
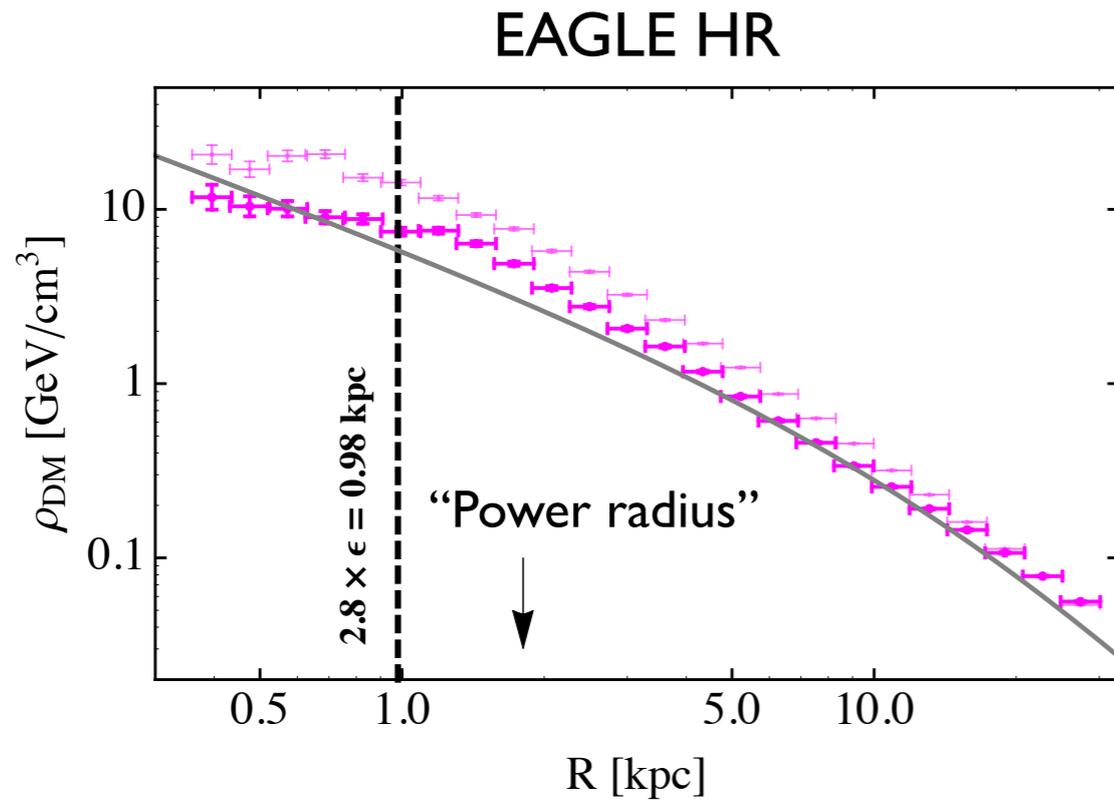
- DM interpretation:**  
Best fit value for the inner slope:  $\gamma = 1.26 \pm 0.15$
- Other interpretations:** *unresolved millisecond pulsars, diffuse photons from cosmic rays, stellar source population in the Galactic bulge, ...*

# Galactic centre GeV excess

- Test the DM density profile predicted by hydrodynamical simulations against the GeV excess data.
- **Additional selection criterion of MW-like galaxies:** substantial stellar disk component.

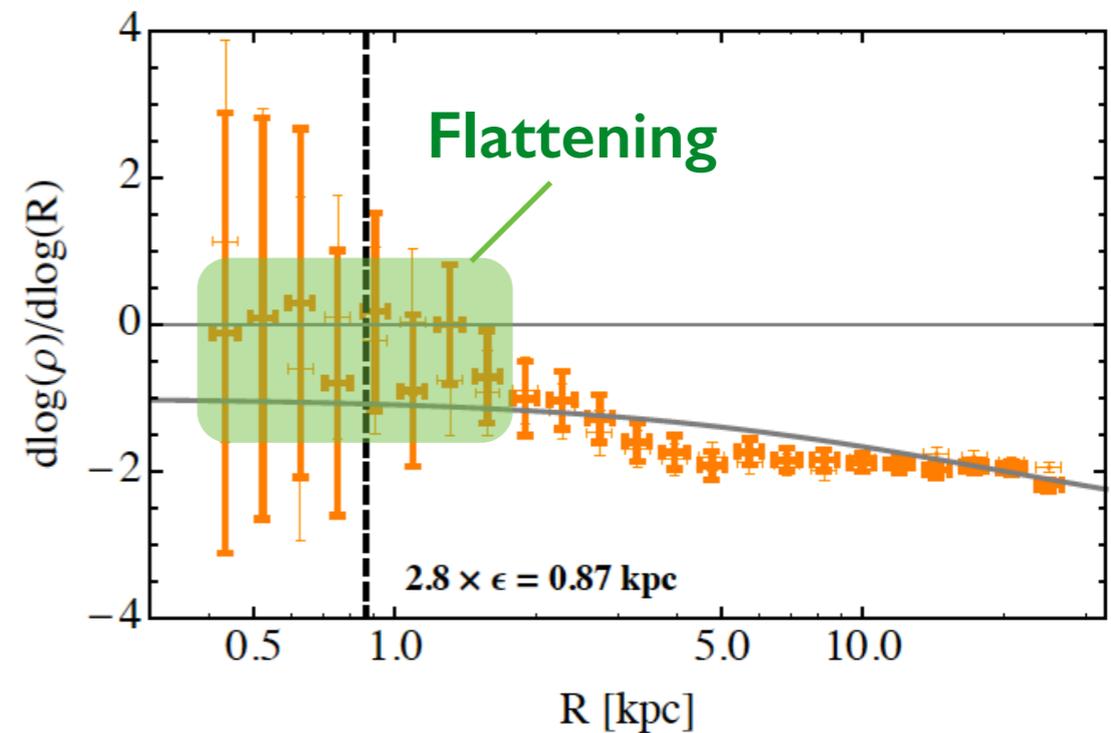
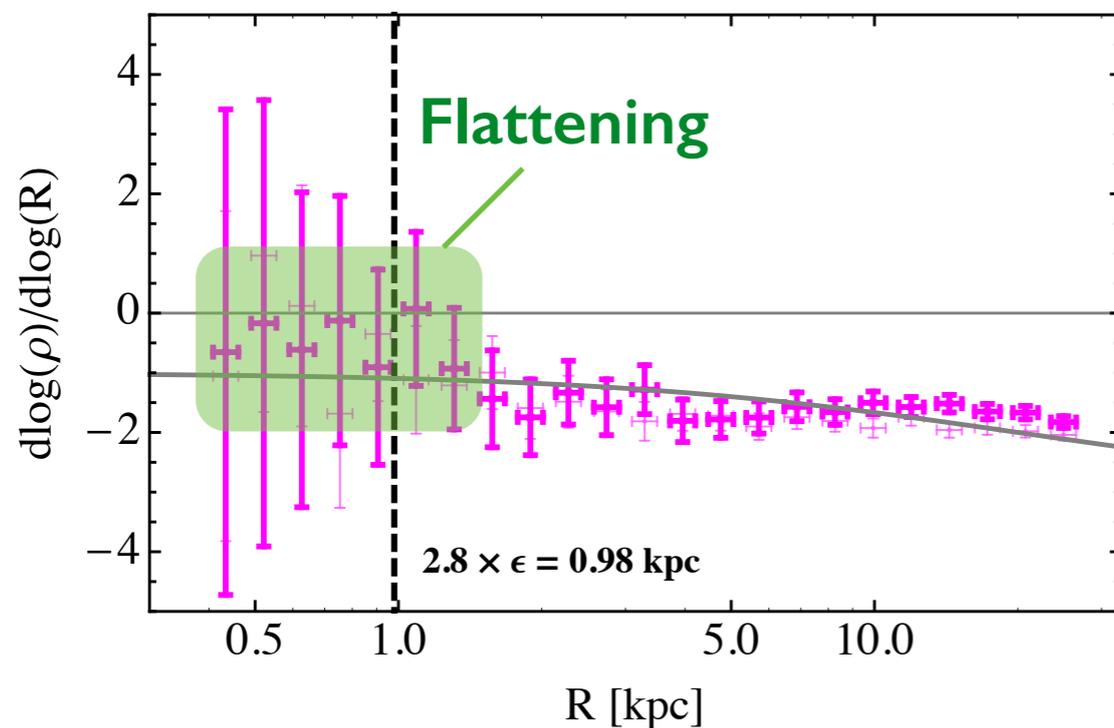
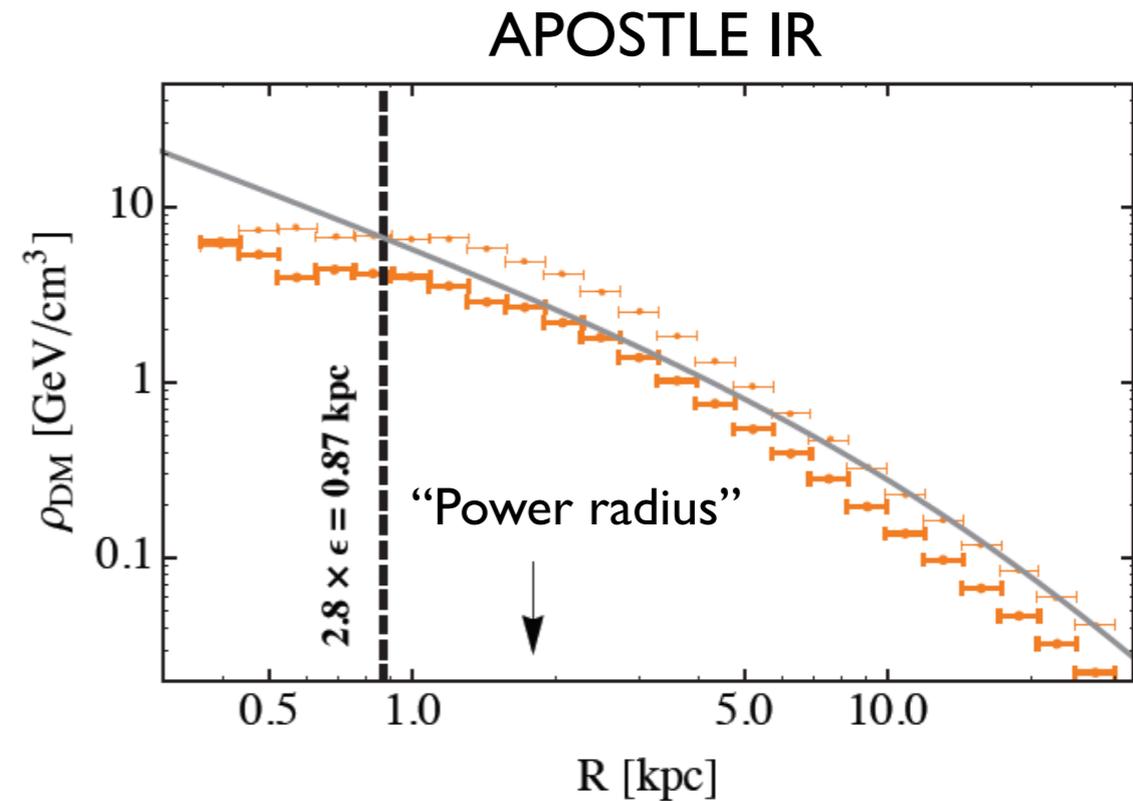
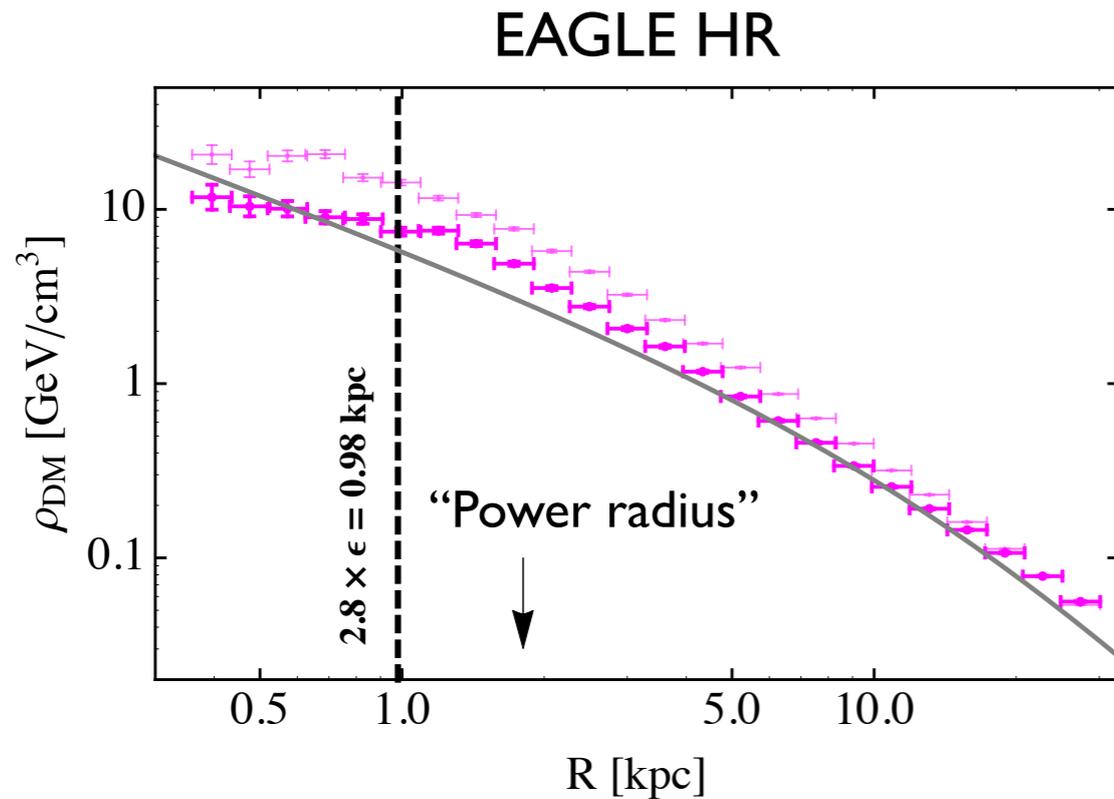
4 MW analogues:  
2 EAGLE + 2 APOSTLE

# DM density profiles



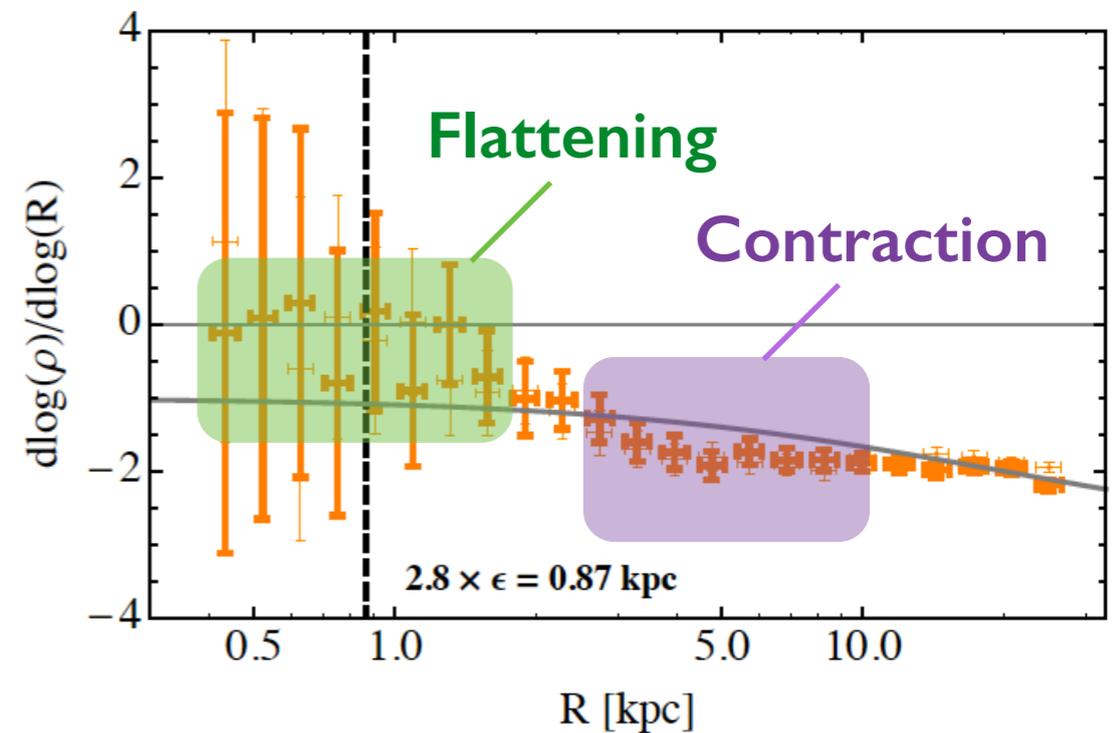
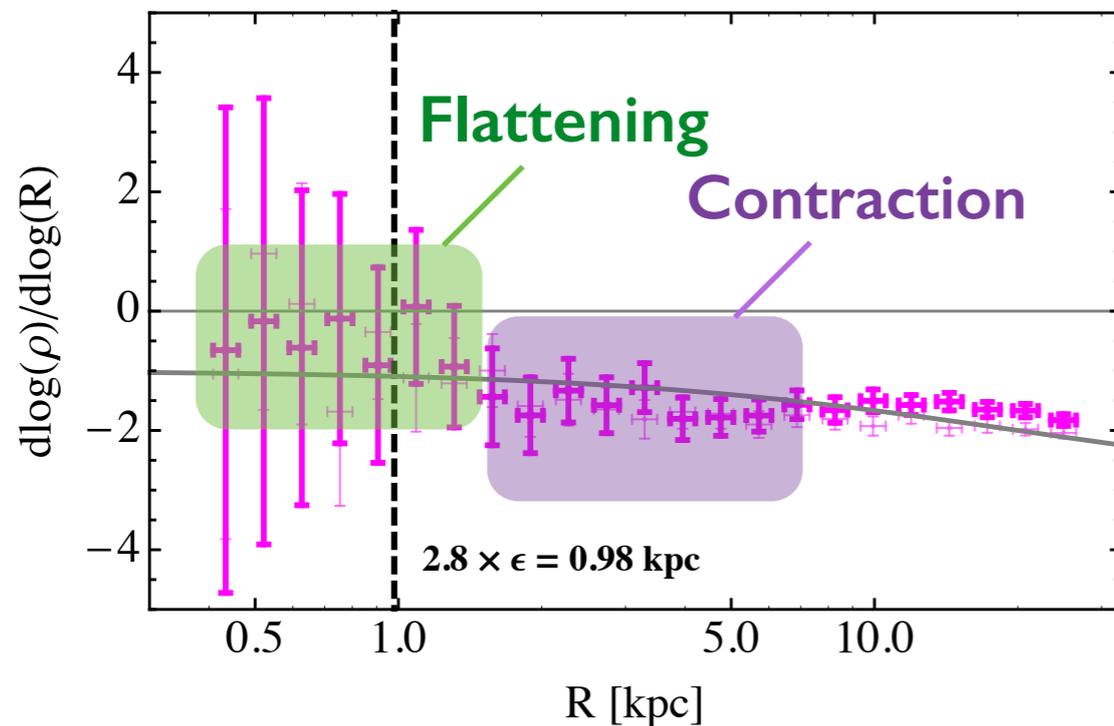
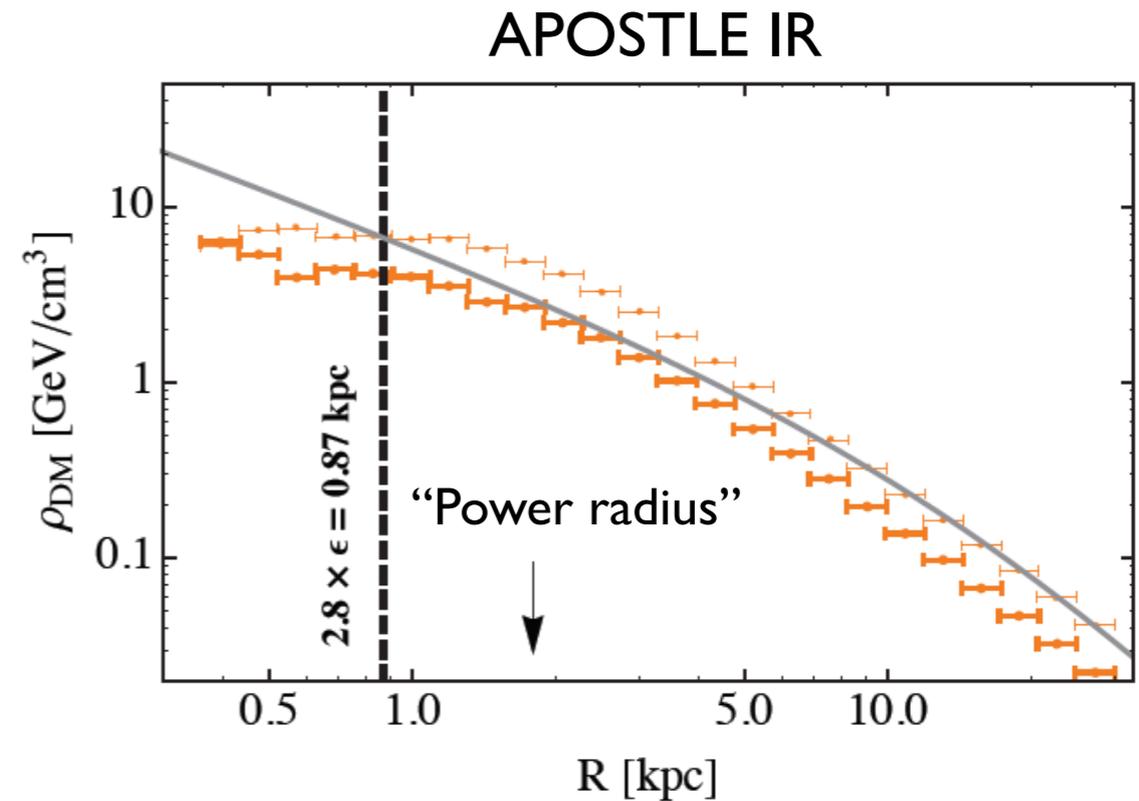
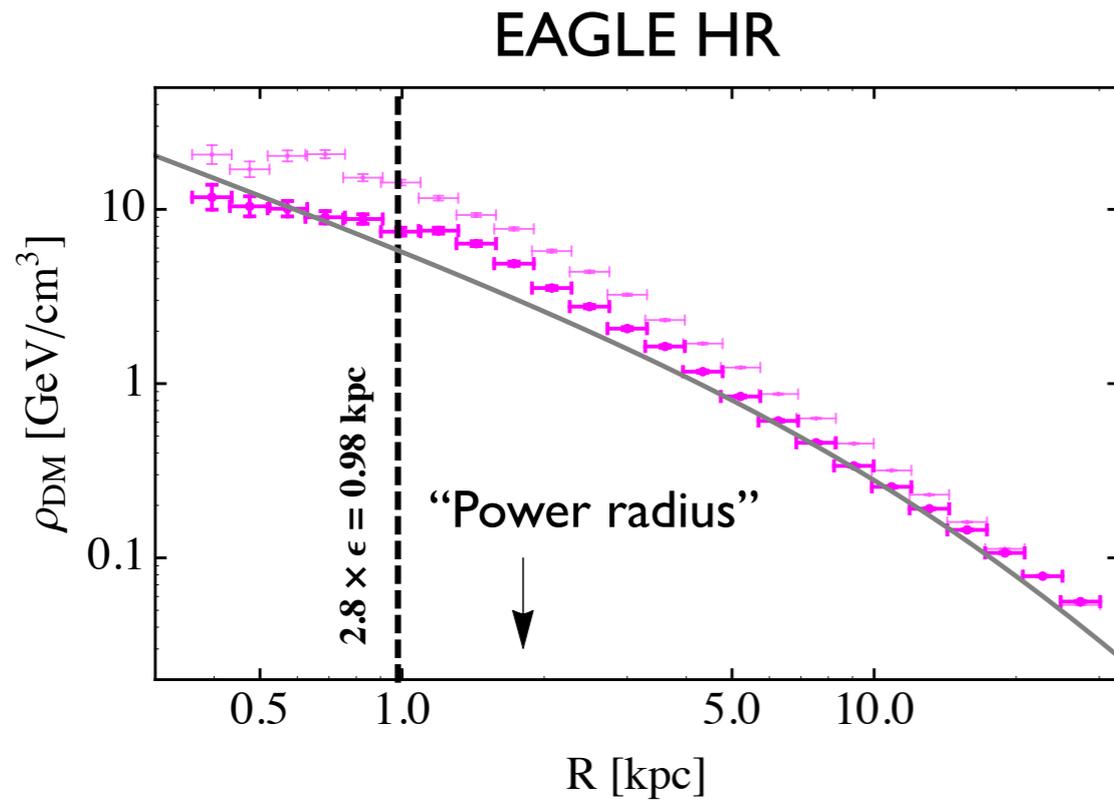
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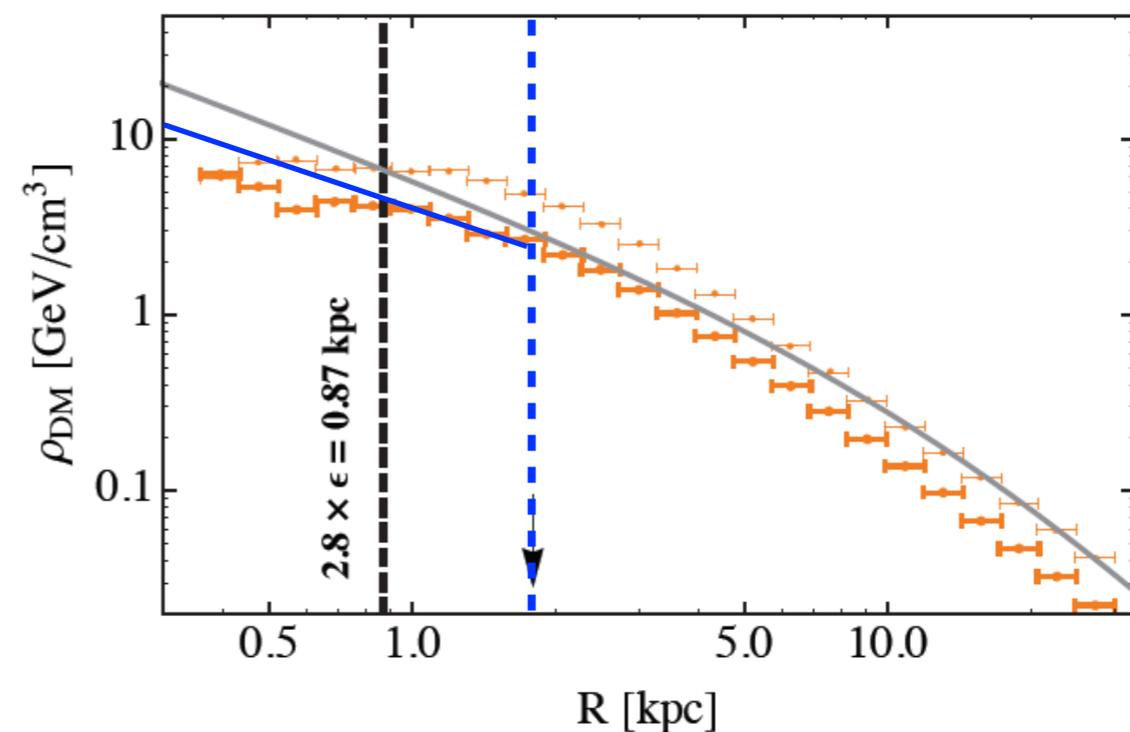
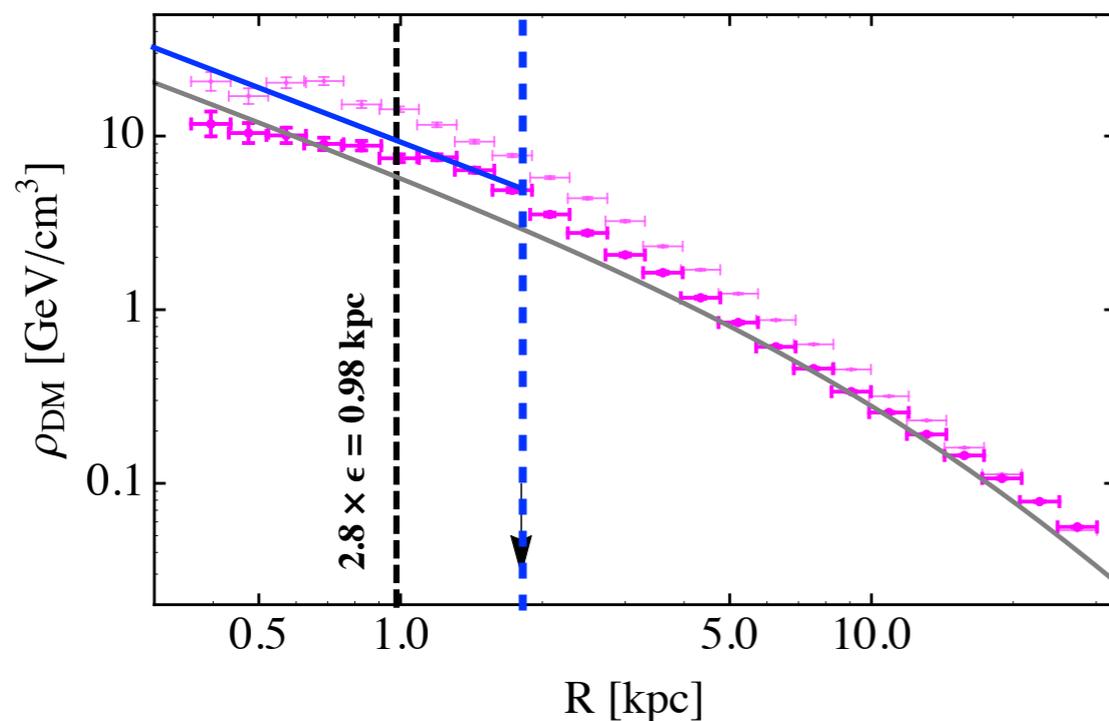
# DM density profiles

- GeV excess data analyzed in the region:

$$2^\circ \leq |b| \leq 20^\circ \ \& \ |l| \leq 20^\circ$$

radial scale: 0.3 - 3 kpc

- **A very conservative approach:** power-law extrapolation with maximal asymptotic slope at the Power radius.

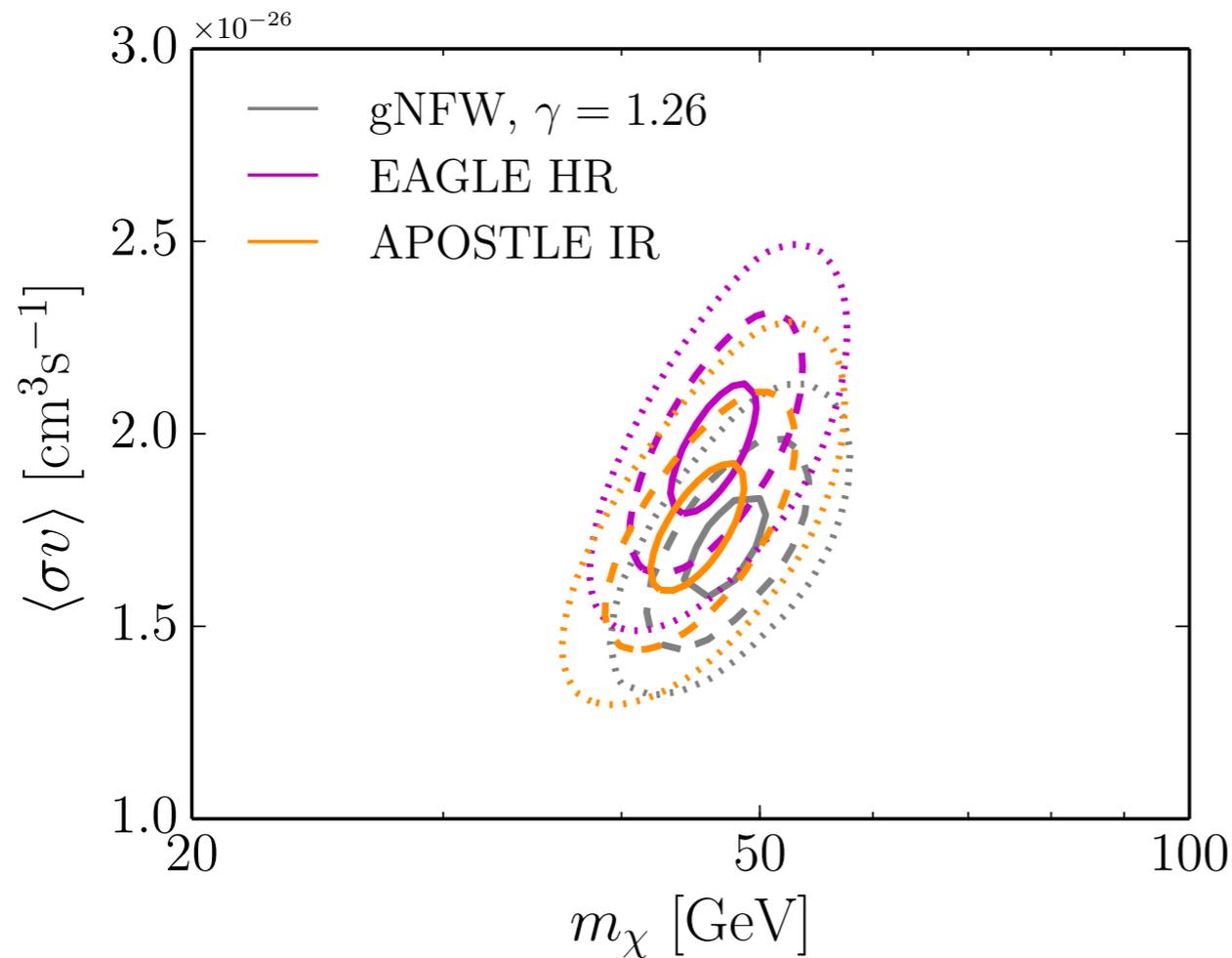


EAGLE HR (2 haloes):  $0.94 < \gamma_{\max} < 0.98$  at  $R_{P03} = 1.8$  kpc

APOSTLE IR (2 haloes):  $0.50 < \gamma_{\max} < 0.62$  at  $R_{P03} = 1.8$  kpc.

# Fitting the GeV excess

- Assuming 100% annihilation into b-quarks:



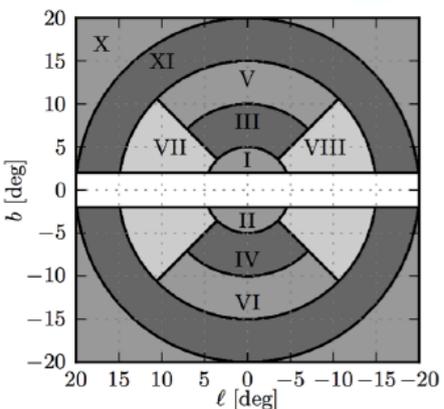
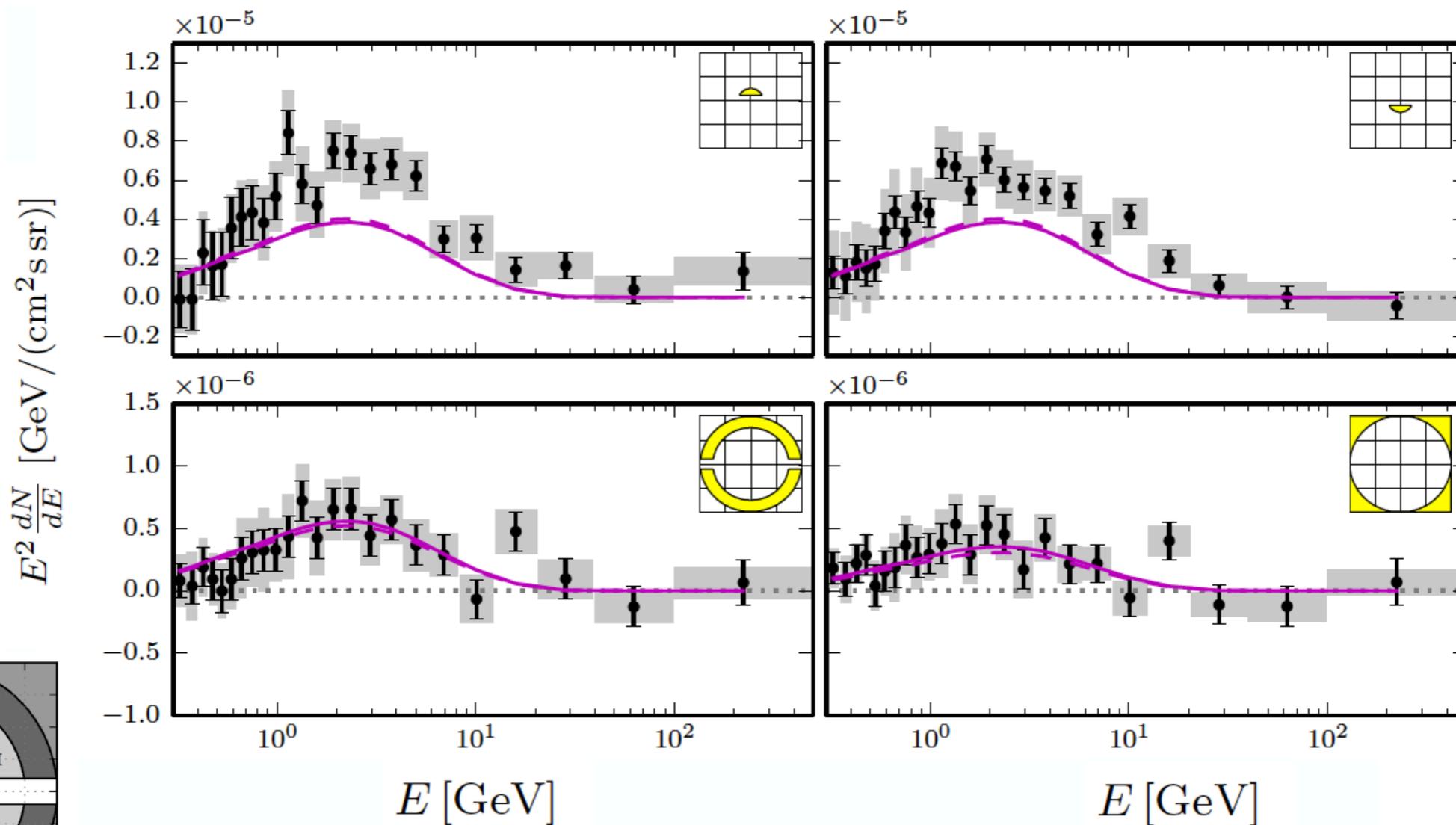
- Similar constraints on DM mass and annihilation cross section, but significantly worse fit.

(238 dof)

| Profile                | $\langle\sigma v\rangle [\times 10^{-26} \text{ cm}^3/\text{s}]$ | $m_\chi$ [GeV]   | $\chi^2$ | $p$ -value |
|------------------------|--|------------------|----------|------------|
| gNFW ( $\gamma=1.26$ ) | $1.71 \pm 0.11$  | $47.32 \pm 1.07$ | 223.9    | 0.73       |
| EAGLE HR               | $1.96 \pm 0.14$  | $46.37 \pm 1.37$ | 246.3    | 0.34       |
| APOSTLE IR             | $1.76 \pm 0.16$  | $45.36 \pm 2.96$ | 283.9    | 0.02       |

# Fitting the GeV excess

- Even under our very conservative assumption, DM density profiles of our MW-like galaxies do not reproduce the correct morphology of the GeV excess in the inner most regions.



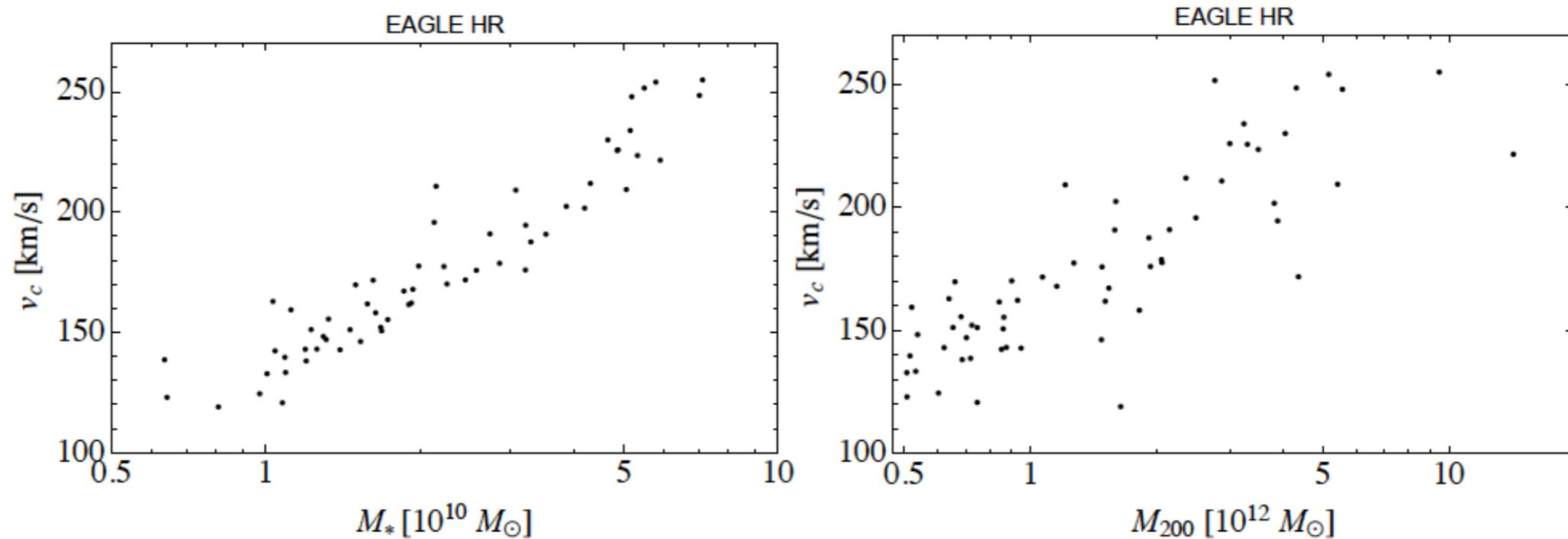
Calore, Bozorgnia et al., 1509.02164

# Summary

- Need a precise determination of the DM distribution in the MW.  
→ *Identify MW analogues* in simulations by taking into account observational constraints on the MW.
- **Local DM density** agrees with local and global estimates.
- **DM density profiles** show flattening in the inner few kpc and contraction up to 10 kpc.
- **Halo integrals** match well those obtained from best fit Maxwellian velocity distributions.
- A **Maxwellian velocity distribution** with *peak speed* constrained by hydro simulations, and independent from the *local circular speed*, could be used for the analysis of direct detection data.
- DM density profiles of MW-like galaxies fail to reproduce the GeV excess.

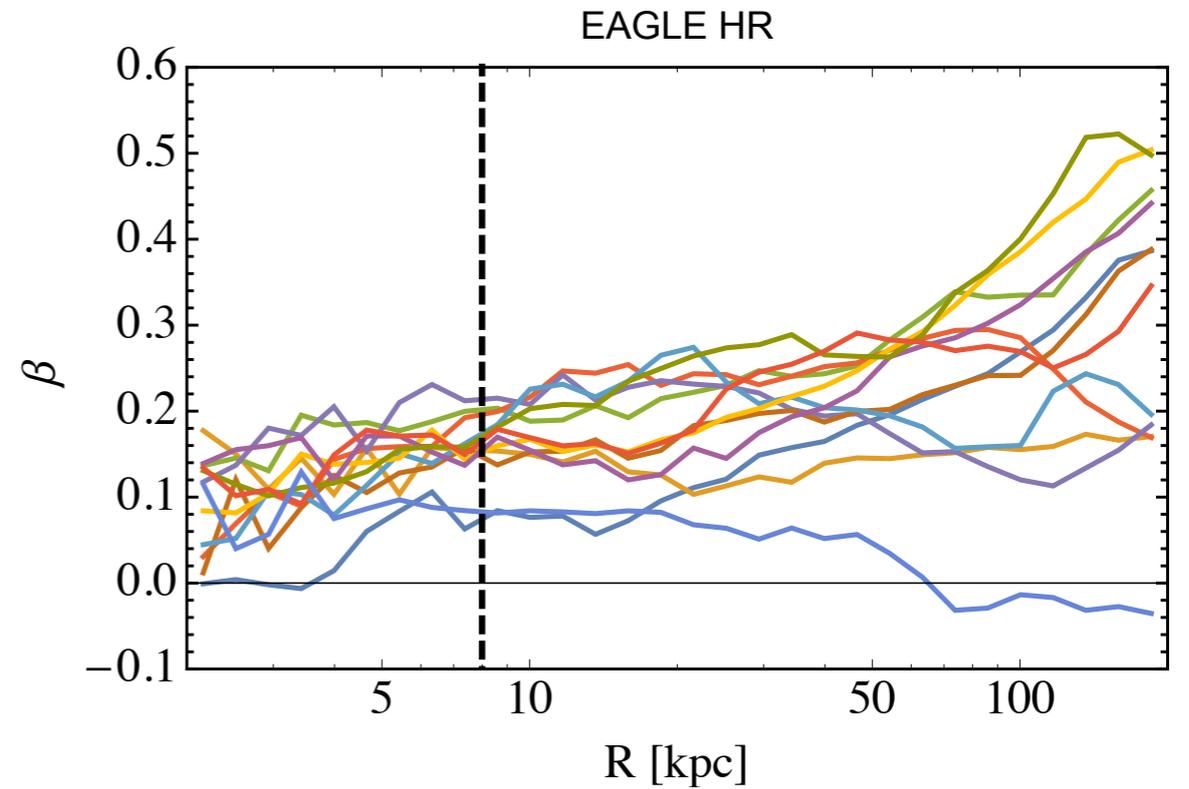
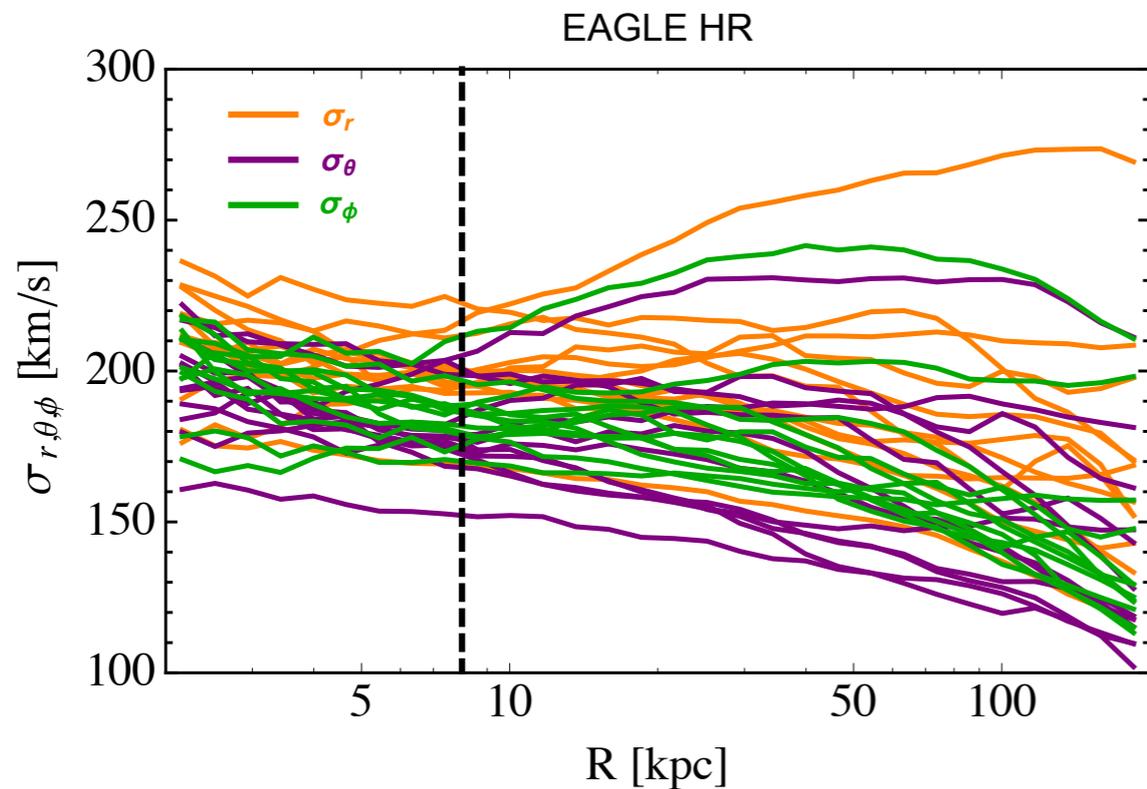
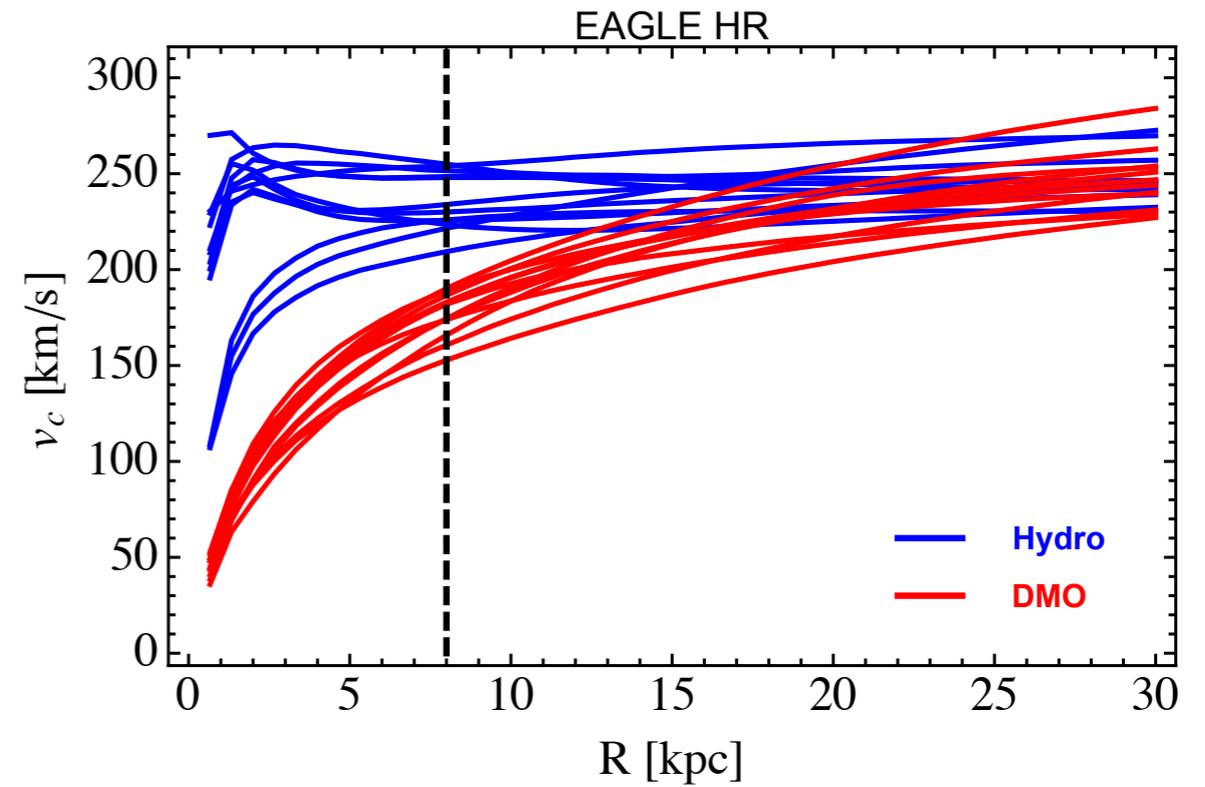
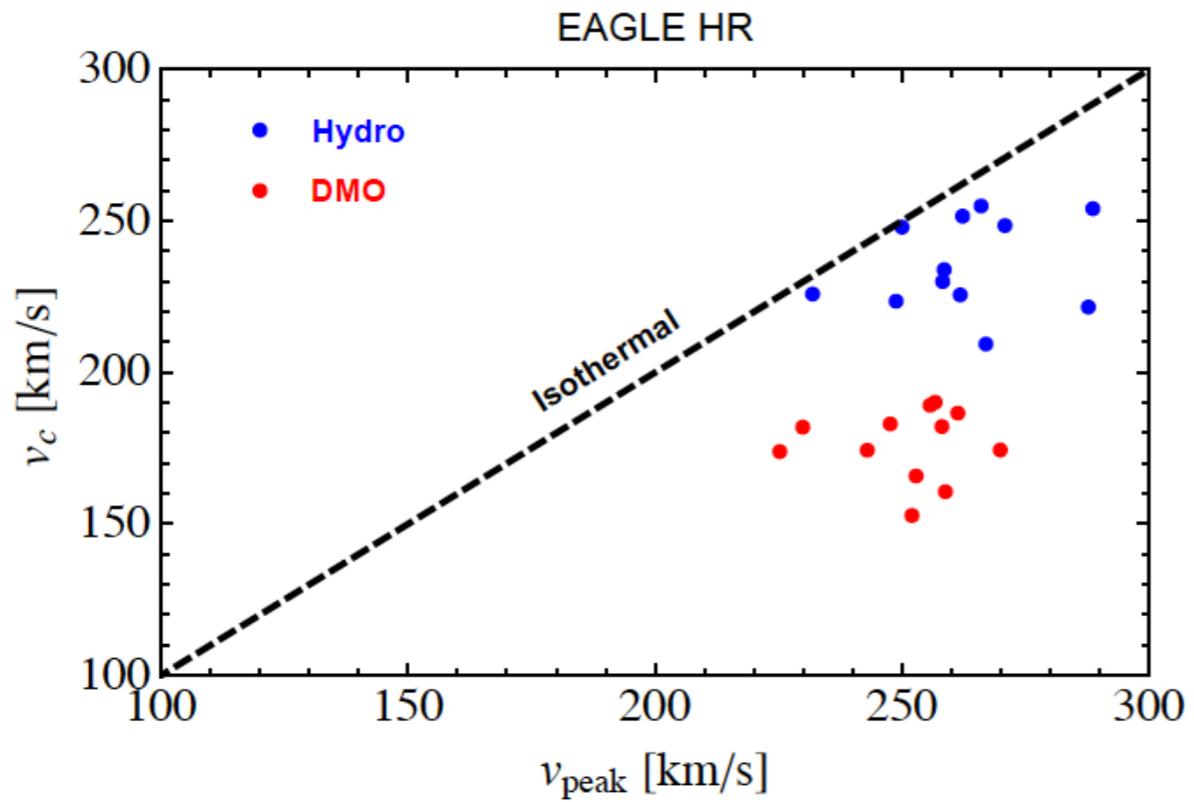
# Backup Slides

# Selection criteria for MW analogues



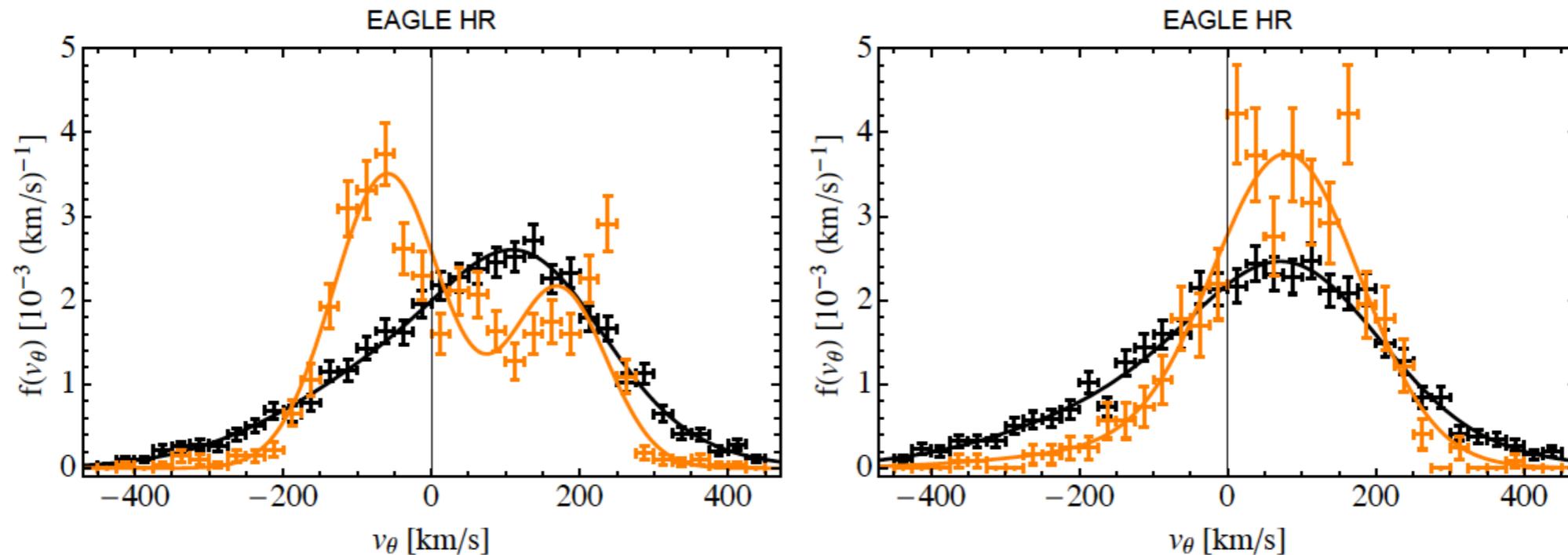
- ▶  $M_*$  strongly correlated with  $v_c$  at 8 kpc, while the correlation of  $M_{200}$  with  $v_c$  is weaker.
- ▶  $M_*(R < 8 \text{ kpc}) = (0.5 - 0.9)M_*$ .
- ▶  $M_{\text{tot}}(R < 8 \text{ kpc}) = (0.01 - 0.1)M_{200}$ .
- ▶ Over the small halo mass range probed, little correlation between  $M_{\text{DM}}(R < 8 \text{ kpc})$  and  $M_{200}$ .

# Departure from isothermal



# Searching for dark disks

DM and stellar velocity distributions:



- ▶ Fit with a double Gaussian. Difference in the mean speed of second Gaussian between DM and stars is 35 km/s in the left, and 7 km/s in the right panel.
- ▶ Fraction of second Gaussian is 32% in the left panel and 43% in the right panel.

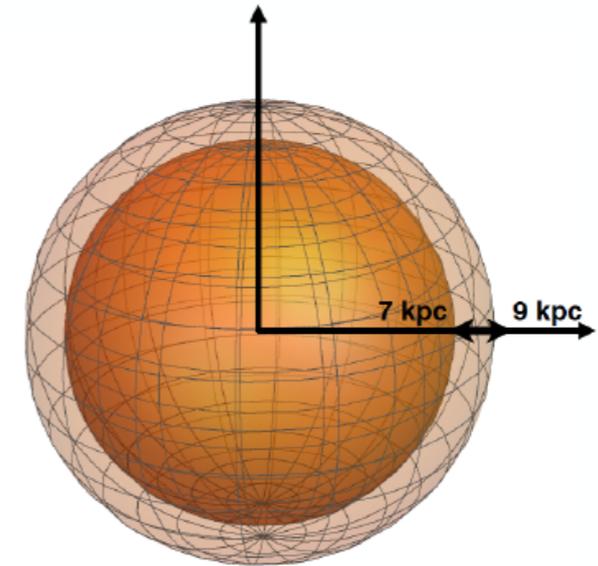
# Searching for dark disks

Is there an enhancement of the local DM density in the **Galactic disc** compared to the **halo**?

- ▶ Compare the the average  $\rho_{\text{DM}}$  in the torus with the value in a spherical shell at  $7 < R < 9$  kpc.

$\rho_{\text{DM}}^{\text{torus}}$  is larger than  $\rho_{\text{DM}}^{\text{shell}}$  by:

2 – 27% for 10 haloes,  
greater than 10% for 5 haloes, and  
greater than 20% for only two haloes.



- ▶ The increase in the DM density in the disc could be due to the DM halo contraction as a result of dissipational baryonic processes.

# Halo shapes

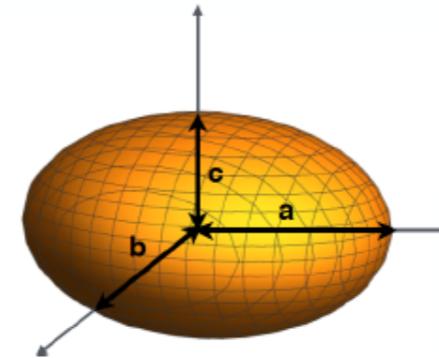
- ▶ To study the shape of the inner ( $R < 8$  kpc) DM haloes, we calculate the inertia tensor of DM particles within 5 and 8 kpc.  
⇒ ellipsoid with three axes of length  $a \geq b \geq c$ .
- ▶ Calculate the **sphericity**:  $s = c/a$ .
  - ▶  $s = 1$ : perfect sphere.  $s < 1$ : increasing deviation from sphericity.
  - ▶ At 5 kpc,  $s = [0.85, 0.95]$ . At 8 kpc,  $s$  lower by less than 10%.
  - ▶ Due to dissipational baryonic processes, DM sphericity systematically higher in the hydrodynamic simulations compared to DMO haloes in which  $s = [0.75, 0.85]$ .

# Halo shapes

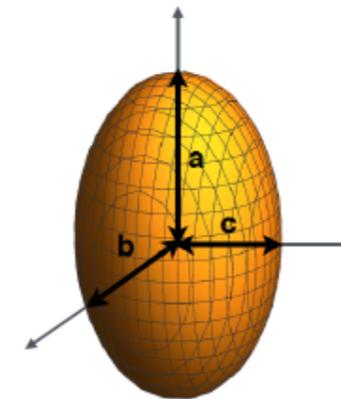
- ▶ Describe a deviation from sphericity by the triaxiality parameter:

$$T = \frac{a^2 - b^2}{a^2 - c^2}$$

- ▶ Oblate systems,  $a \approx b \gg c \Rightarrow T \approx 0$ .



- ▶ Prolate systems,  $a \gg b \approx c \Rightarrow T \approx 1$ .



- ▶ In the hydro case, since inner haloes are very close to spherical, deviation towards either oblate or prolate is small. **DMO counterparts** have a preference for *prolate* inner haloes.

# Parameters of the simulations

| Simulation           | code               | $N_{\text{DM}}$ | $m_g [M_{\odot}]$  | $m_{\text{DM}} [M_{\odot}]$ | $\epsilon$ [pc] |
|----------------------|--------------------|-----------------|--------------------|-----------------------------|-----------------|
| Ling <i>et al.</i>   | RAMSES             | 2662            | –                  | $7.46 \times 10^5$          | 200             |
| Eris                 | GASOLINE           | 81213           | $2 \times 10^4$    | $9.80 \times 10^4$          | 124             |
| NIHAO                | EFS-GASOLINE2      | –               | $3.16 \times 10^5$ | $1.74 \times 10^6$          | 931             |
| EAGLE (HR)           | P-GADGET (ANARCHY) | 1821–3201       | $2.26 \times 10^5$ | $1.21 \times 10^6$          | 350             |
| APOSTLE (IR)         | P-GADGET (ANARCHY) | 2160, 3024      | $1.3 \times 10^5$  | $5.9 \times 10^5$           | 308             |
| MaGICC               | GASOLINE           | 4849, 6541      | $2.2 \times 10^5$  | $1.11 \times 10^6$          | 310             |
| Sloane <i>et al.</i> | GASLOINE           | 5847–7460       | $2.7 \times 10^4$  | $1.5 \times 10^5$           | 174             |

## Properties of the selected MW analogues

| Simulation           | Count | $M_{\text{star}} [\times 10^{10} M_{\odot}]$ | $M_{\text{halo}} [\times 10^{12} M_{\odot}]$ | $\rho_{\chi} [\text{GeV}/\text{cm}^3]$ | $v_{\text{peak}} [\text{km/s}]$ |
|----------------------|-------|--|--|--|---------------------------------|
| Ling <i>et al.</i>   | 1     | $\sim 8$                                     | 0.63   | 0.37–0.39                              | 239                             |
| Eris                 | 1     | 3.9  | 0.78   | 0.42                                   | 239                             |
| NIHAO                | 5     | 15.9   | $\sim 1$                                     | 0.42                                   | 192–363                         |
| EAGLE (HR)           | 12    | 4.65–7.12                                    | 2.76–14.26                                   | 0.42–0.73                              | 232–289                         |
| APOSTLE (IR)         | 2     | 4.48, 4.88                                   | 1.64–2.15                                    | 0.41–0.54                              | 223–234                         |
| MaGICC               | 2     | 2.4–8.3                                      | 0.584, 1.5                                   | 0.346, 0.493                           | 187, 273                        |
| Sloane <i>et al.</i> | 4     | 2.24–4.56                                    | 0.68–0.91                                    | 0.3–0.4                                | 185–204                         |

# Morphology of simulated haloes

- ▶ Select simulated galaxies whose stellar kinematics show a disc component, rather than ellipticals or undergoing mergers.
- ▶ Characterize the morphology of each simulated galaxy by looking for evidence of coherent rotation.
- ▶ Use the distribution of angular momentum vectors of individual particles relative to the net angular momentum of the galaxy to discriminate between discs (coherent rotation) and spheroids (no coherent rotation).
- ▶ Derive the distribution of the stellar orbital circularity parameter,

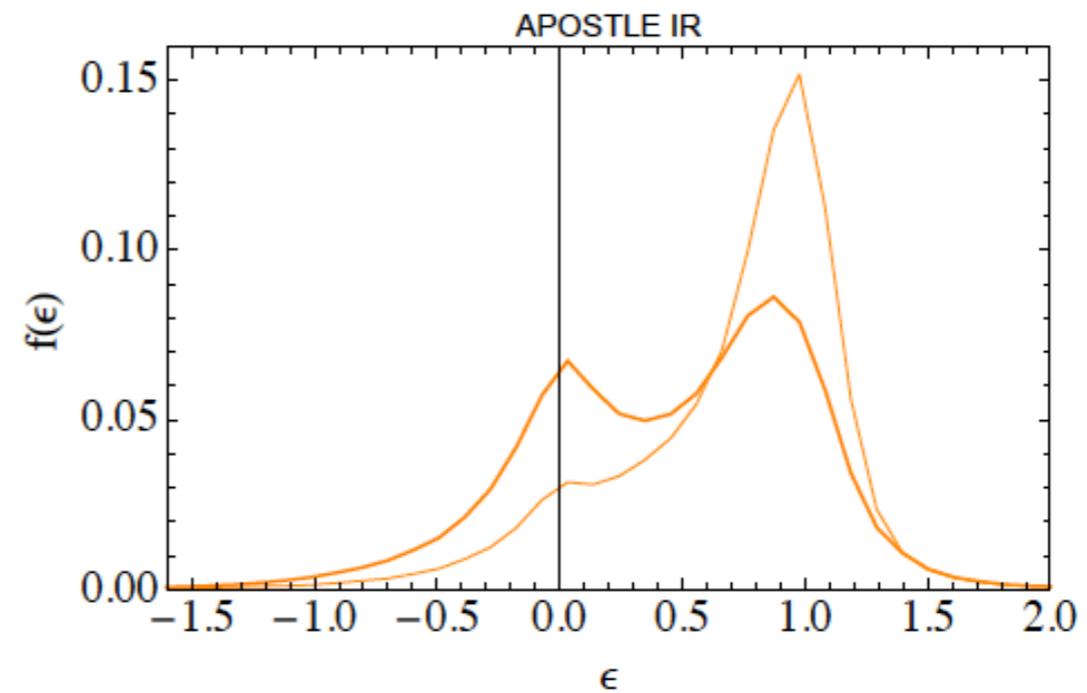
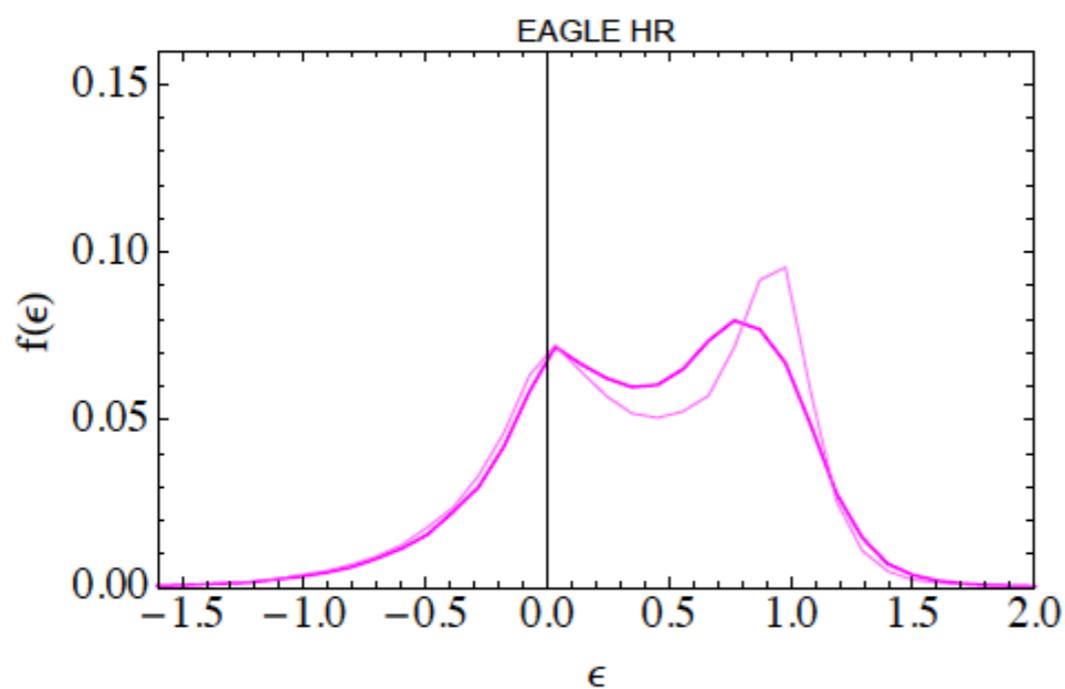
$$\epsilon(r) = \frac{j_z}{j_c(r)}$$

A distribution peaked at  $\epsilon = 1 \Rightarrow$  disc

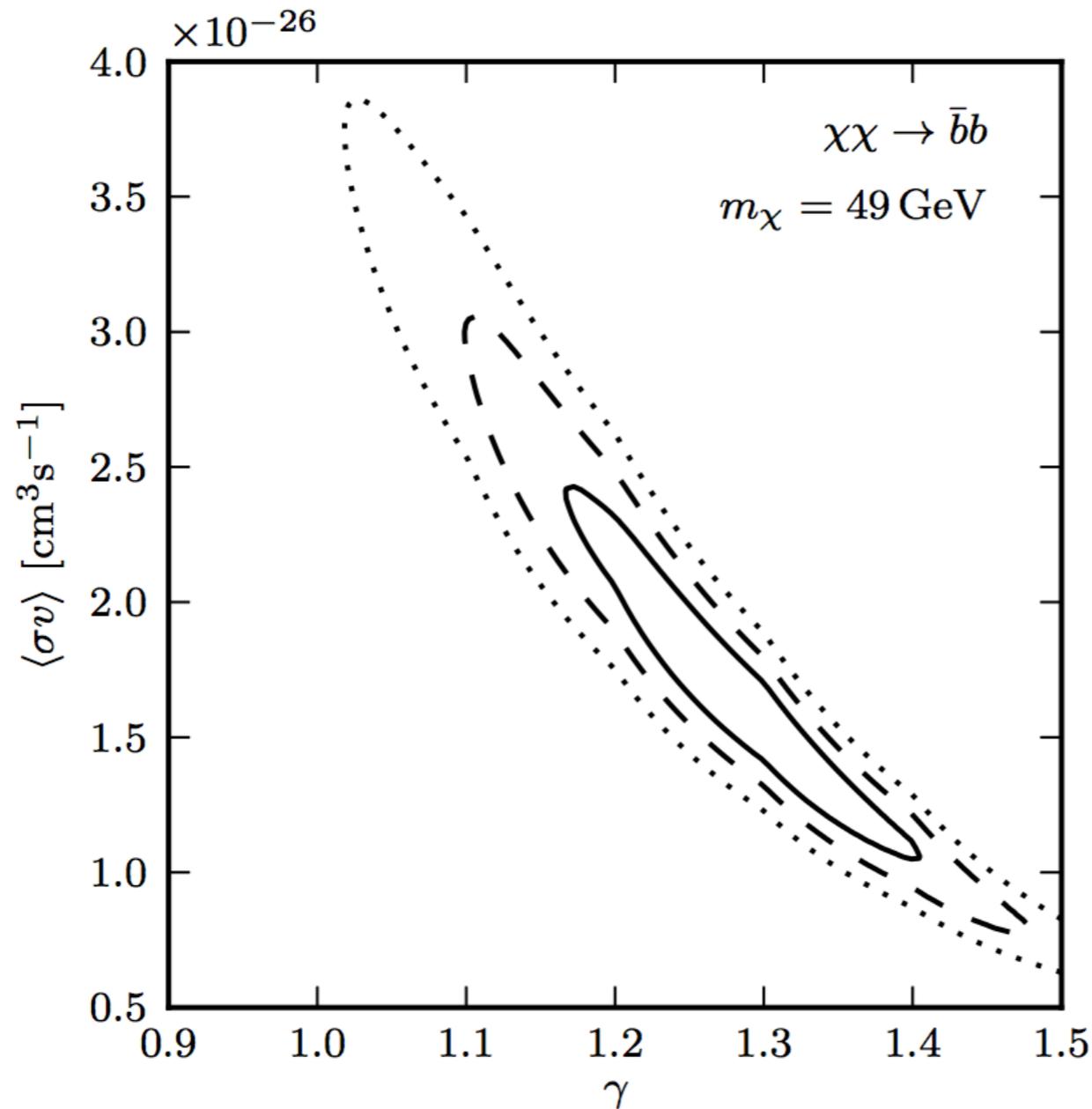
An almost symmetric distribution around  $\epsilon = 0 \Rightarrow$  spheroidal system

# Morphology of simulated haloes

- ▶ We retain a galaxy if the stellar fraction in the range  $\epsilon > 0.45$  is larger than 50%.
- ▶ With this criterion we can identify galaxies that have a dominant disc, and remove galaxies that show an almost symmetric distribution around  $\epsilon = 0$ .



# GeV excess spatial profile

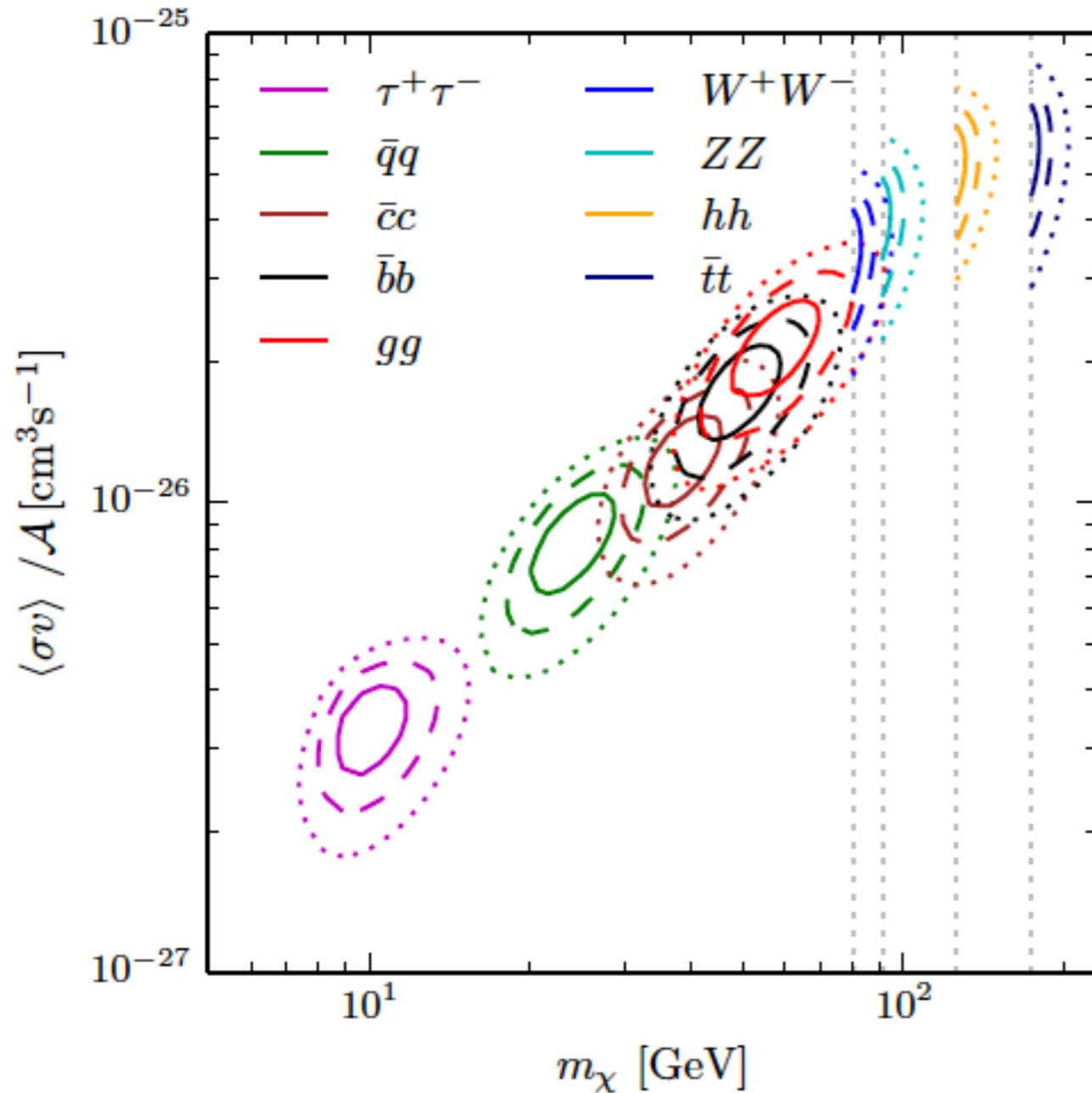


Calore et al., I409.0042

Generalized NFW:

$$\rho(r) = \rho_s \frac{r_s^3}{r^\gamma (r + r_s)^{3-\gamma}}$$

# GeV excess DM interpretation



| Channel        | $\langle\sigma v\rangle$<br>( $10^{-26} \text{ cm}^3 \text{ s}^{-1}$ ) | $m_\chi$<br>(GeV)      | $\chi^2_{\min}$ | $p$ -value             |
|----------------|--|------------------------|-----------------|------------------------|
| $\bar{q}q$     | $0.83^{+0.15}_{-0.13}$   | $23.8^{+3.2}_{-2.6}$   | 26.7            | 0.22                   |
| $\bar{c}c$     | $1.24^{+0.15}_{-0.15}$   | $38.2^{+4.7}_{-3.9}$   | 23.6            | 0.37                   |
| $\bar{b}b$     | $1.75^{+0.28}_{-0.26}$   | $48.7^{+6.4}_{-5.2}$   | 23.9            | 0.35                   |
| $\bar{t}t$     | $5.8^{+0.8}_{-0.8}$  | $173.3^{+2.8}_{-0}$    | 43.9            | 0.003                  |
| $gg$           | $2.16^{+0.35}_{-0.32}$   | $57.5^{+7.5}_{-6.3}$   | 24.5            | 0.32                   |
| $W^+W^-$       | $3.52^{+0.48}_{-0.48}$   | $80.4^{+1.3}_{-0}$     | 36.7            | 0.026                  |
| $ZZ$           | $4.12^{+0.55}_{-0.55}$   | $91.2^{+1.53}_{-0}$    | 35.3            | 0.036                  |
| $hh$           | $5.33^{+0.68}_{-0.68}$   | $125.7^{+3.1}_{-0}$    | 29.5            | 0.13                   |
| $\tau^+\tau^-$ | $0.337^{+0.047}_{-0.048}$  | $9.96^{+1.05}_{-0.91}$ | 33.5            | 0.055                  |
| $[\mu^+\mu^-]$ | $1.57^{+0.23}_{-0.23}$   | $5.23^{+0.22}_{-0.27}$ | 43.9            | $0.0036]_{\text{JES}}$ |

Calore et al., [1411.4647](#)